

# “Generalized” bag models for the QCD pressure

Models for the pressure in QCD in hydrodynamics:

$T > T_c$ : (semi) ideal gas plus MIT bag constant.

$T < T_c$ : hadron resonance gas

2007: Lattice results with dynamical quarks *not* near continuum limit.

Today: try to extract *general* properties with *simple* models

Suggest:

$T > 1.1 T_c$ : “generalized bag model”

$T$ :  $0.9 T_c$  to  $1.1 T_c$ : *nearly* critical region

$T < 0.9 T_c$ : hadron resonance gas

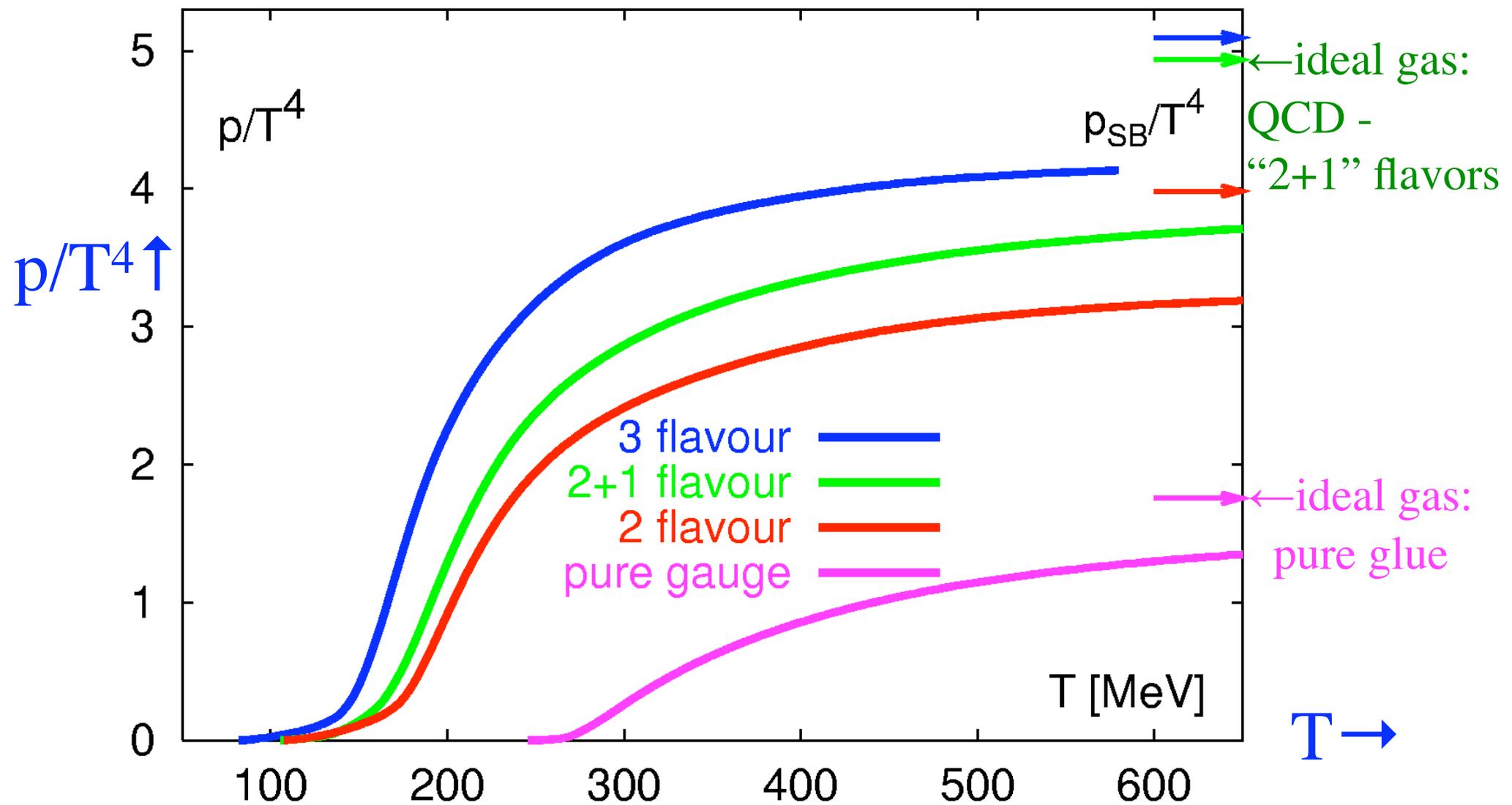
# Lattice: pressure with, and without, quarks

Adding quarks, ideal gas value increases by  $\sim 3$ .

“ $T_c$ ” decreases. Today:  $\sim 190 \pm ?$  MeV Not 1st order, *crossover*.

Pressure nonzero  $< T_c$ : from pions, kaons..., masses  $\sim T_c$

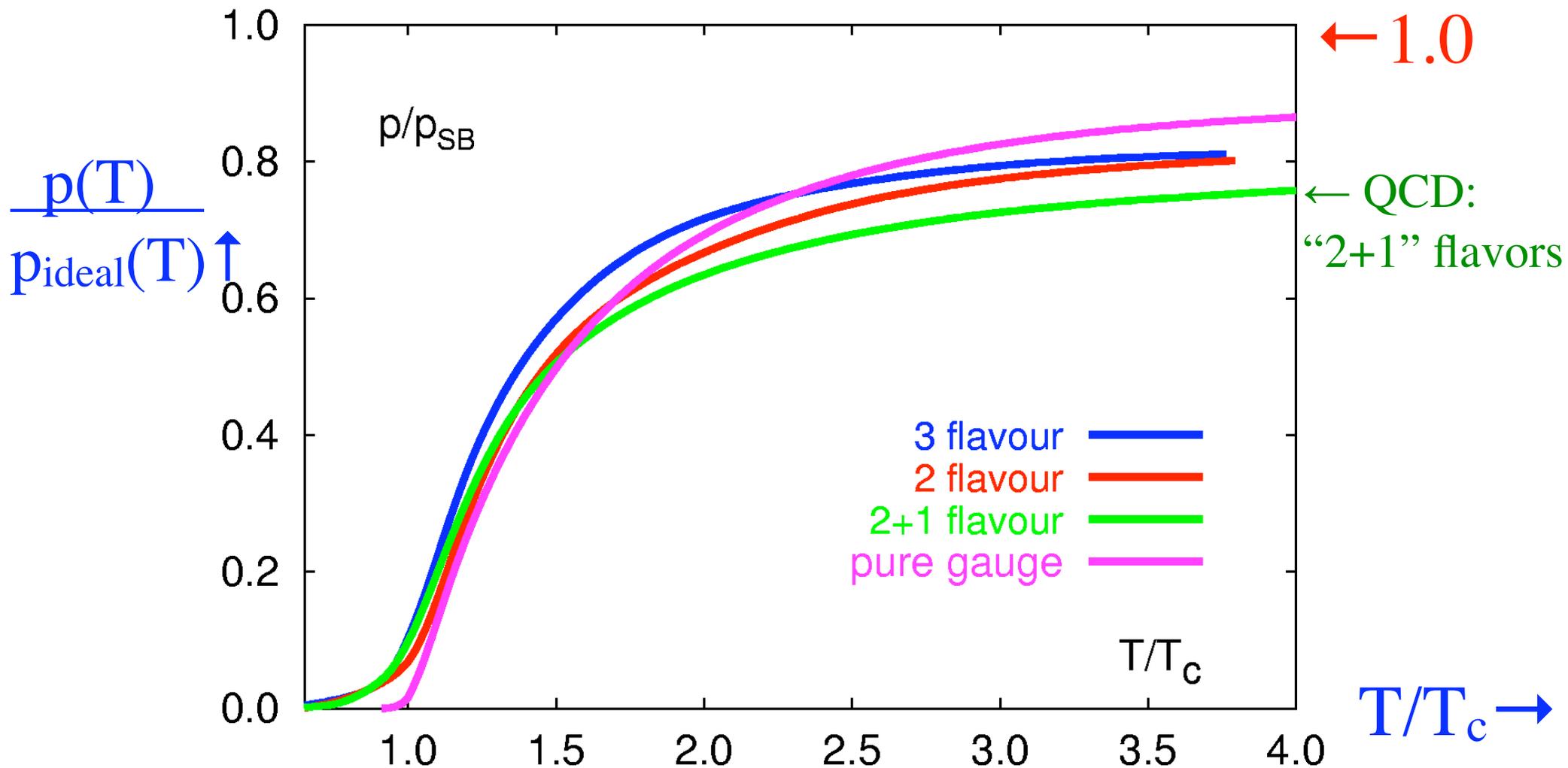
Pressure within 20% of ideal by  $\sim 3 \times T_c$ . Curves *seem* very different



# “Flavor independence”

Bielefeld: scale  $p$  by  $p_{\text{ideal}}$ ,  $T$  by  $T_c$ : *nearly* universal form for pressure

Not exact, approximate. Now concentrate on the differences.



# Some elementary thermodynamics

Let  $e(T)$  = energy,  $p(T)$  = pressure

$$\frac{e - 3p}{T^4} = T \frac{\partial}{\partial T} \frac{p(T)}{T^4}$$

$e - 3p$  = trace of energy momentum tensor.

Measures deviation from conformally symmetric limit.

More prosaically: if

$$p(T) = c_4 T^4 + c_2 T^2 + c_0 + \dots$$

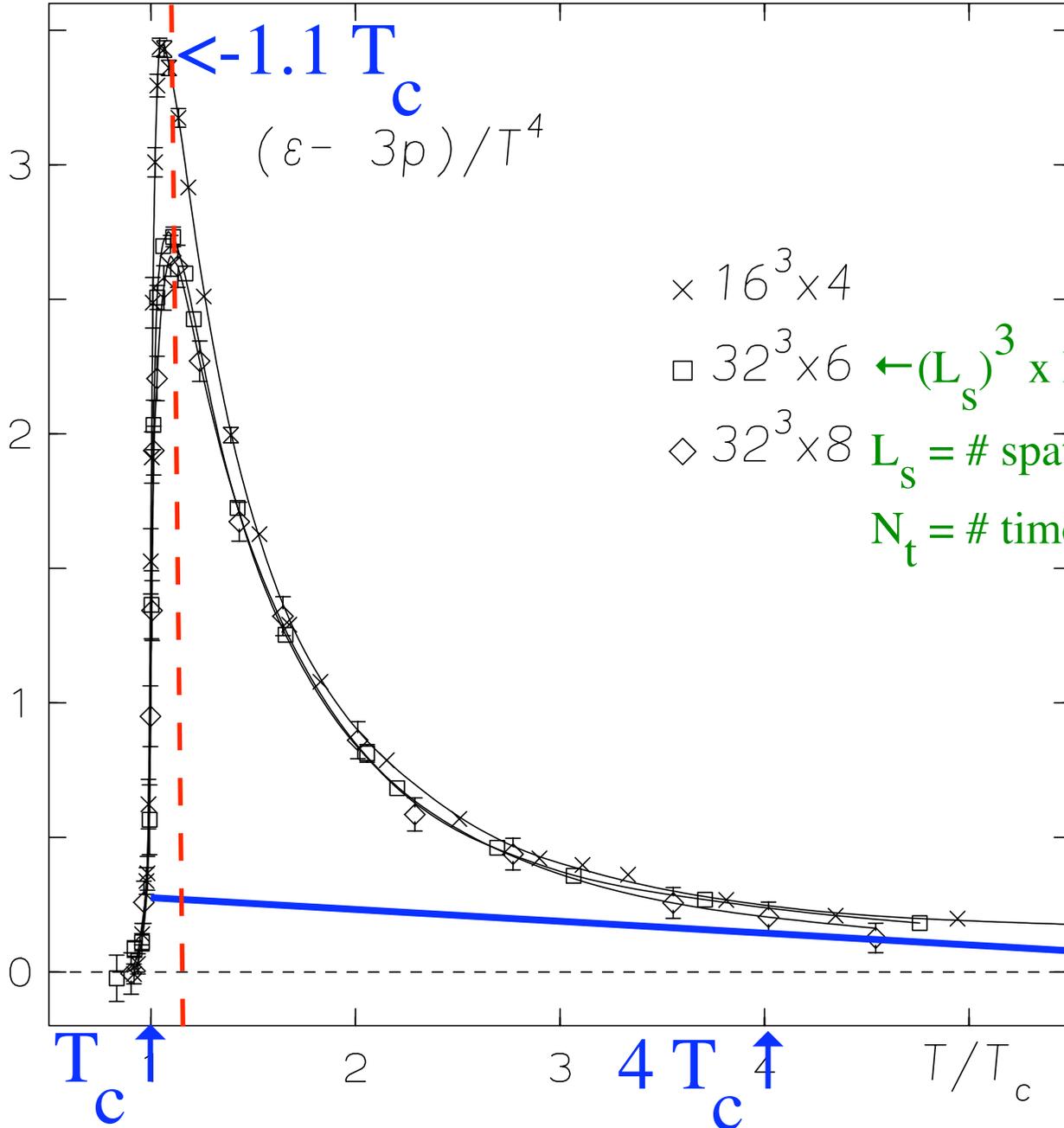
for constant  $c_4, c_2, c_0$ ,  $(e-3p)/T^4$  “strips” off  $T^4$  term, leaving *non-ideal* terms:

$$\frac{e - 3p}{T^4} \times T^2 = -2 c_2 - 4 \frac{c_0}{T^2} + \dots$$

# Lattice: *pure SU(3) glue, no quarks*

Pure glue:  $N_t = \#$  time steps. Claim:  $N_t = 6, 8$  close to continuum limit

$$\frac{e - 3p}{T^4} \uparrow$$



$\times 16^3 \times 4$

$\square 32^3 \times 6 \leftarrow (L_s)^3 \times N_t$

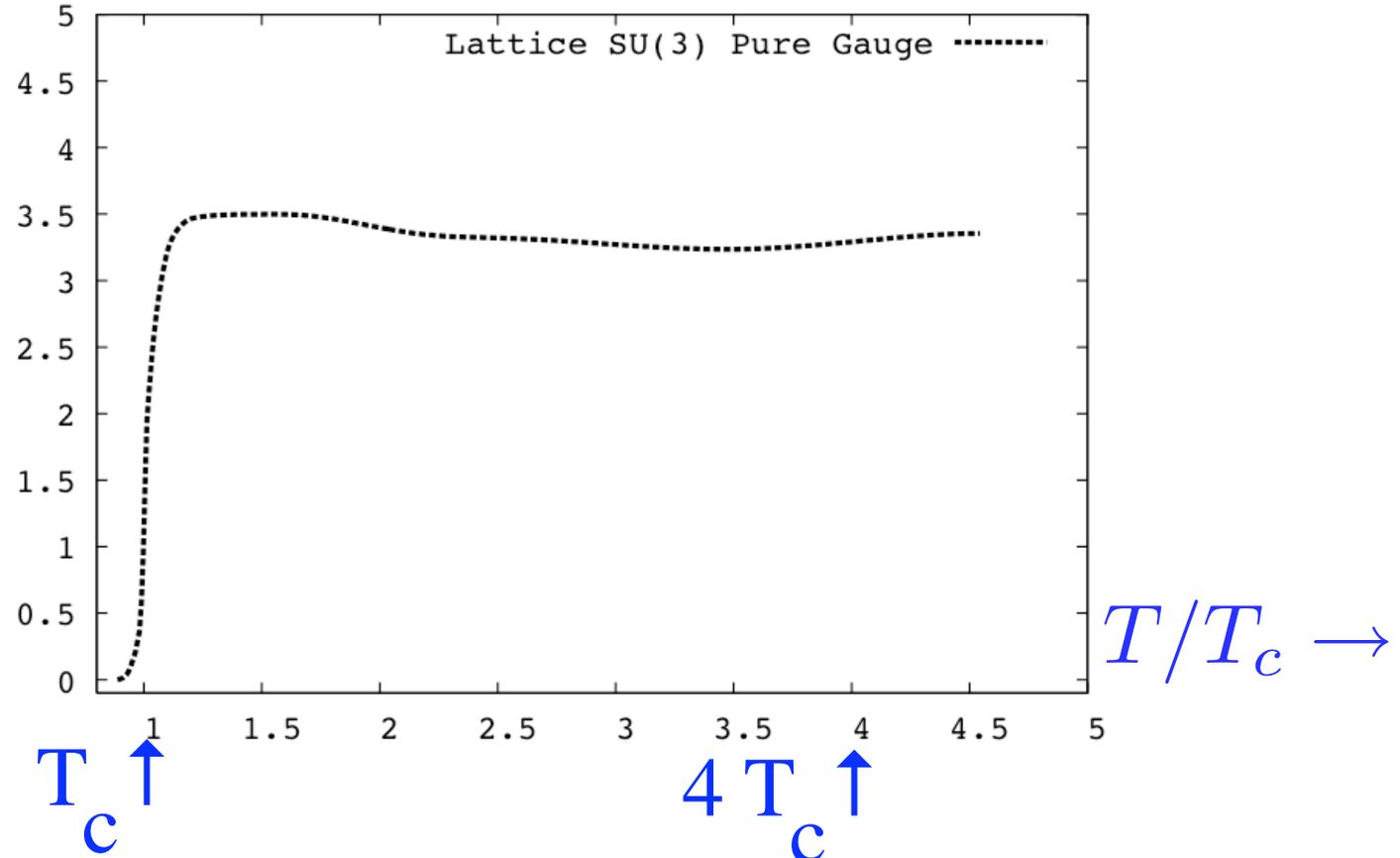
$\diamond 32^3 \times 8$   $L_s = \#$  spatial lattice spacings  
 $N_t = \#$  time steps

Old story:  
 Bielefeld,  
 lat/9602007

$\leftarrow$  Perturbative  
 cont.  $\sim \alpha_s^2$

# “Fuzzy” bag picture

$$\frac{e - 3p}{T^4} \times T^2 \uparrow$$



For  $T$ :  $1.1 T_c$  to  $4.0 T_c$ , pressure sum of just  $T^4$  and  $T^2$  terms!

Since pressure vanishes at  $T_c$ , know constant in front of  $T^2$  term:

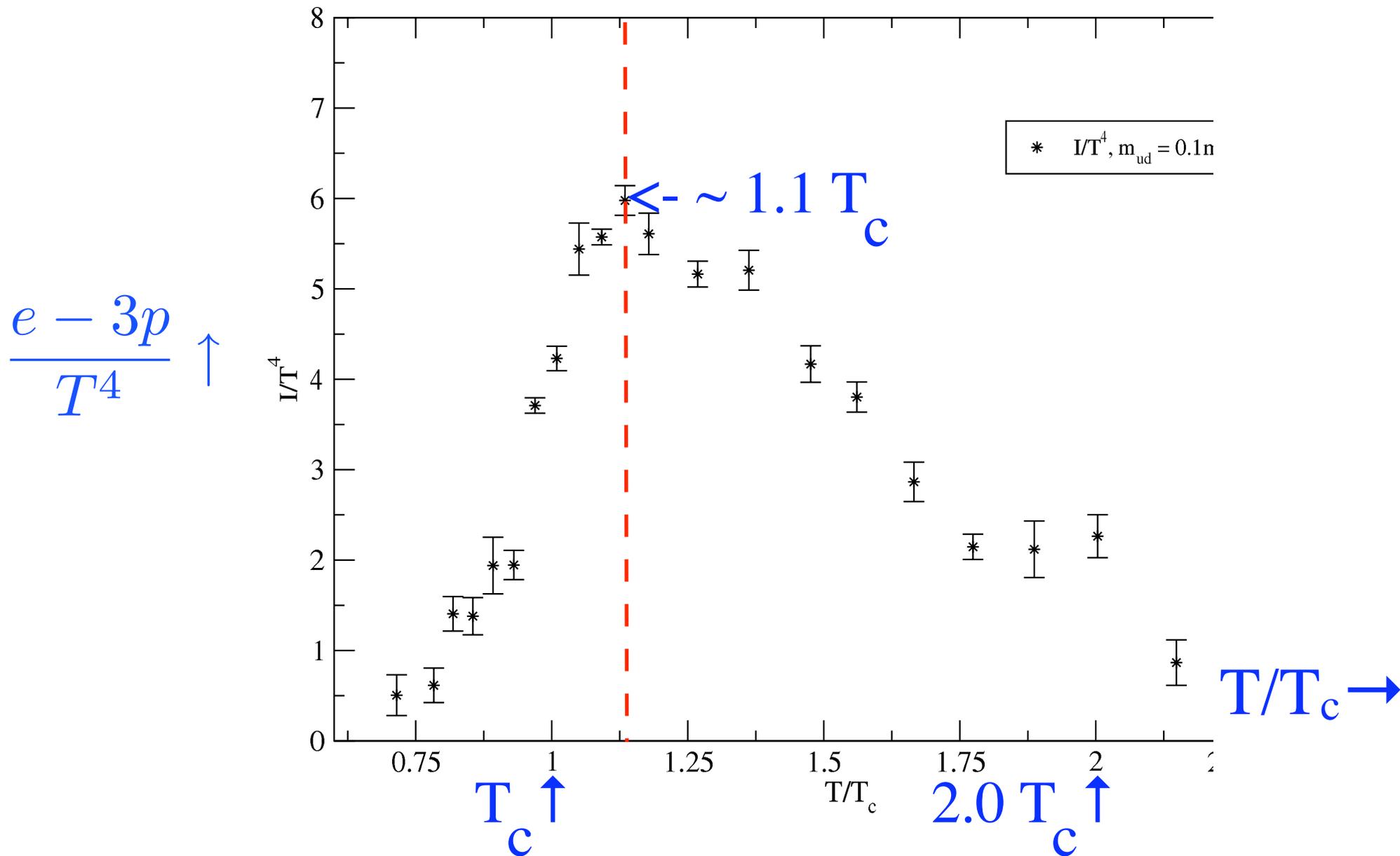
$$p(T) \approx f_{pert}(T^4 - T_c^2 T^2)$$

$f_{pert} \sim \text{constant}$ ,  $T$ :  $1.1 T_c$  to  $4.0 T_c$ . For  $f_{pert}$ , weak coupling resummations of perturbation theory work down to  $T_c$

# Lattice: SU(3) glue *with* quarks

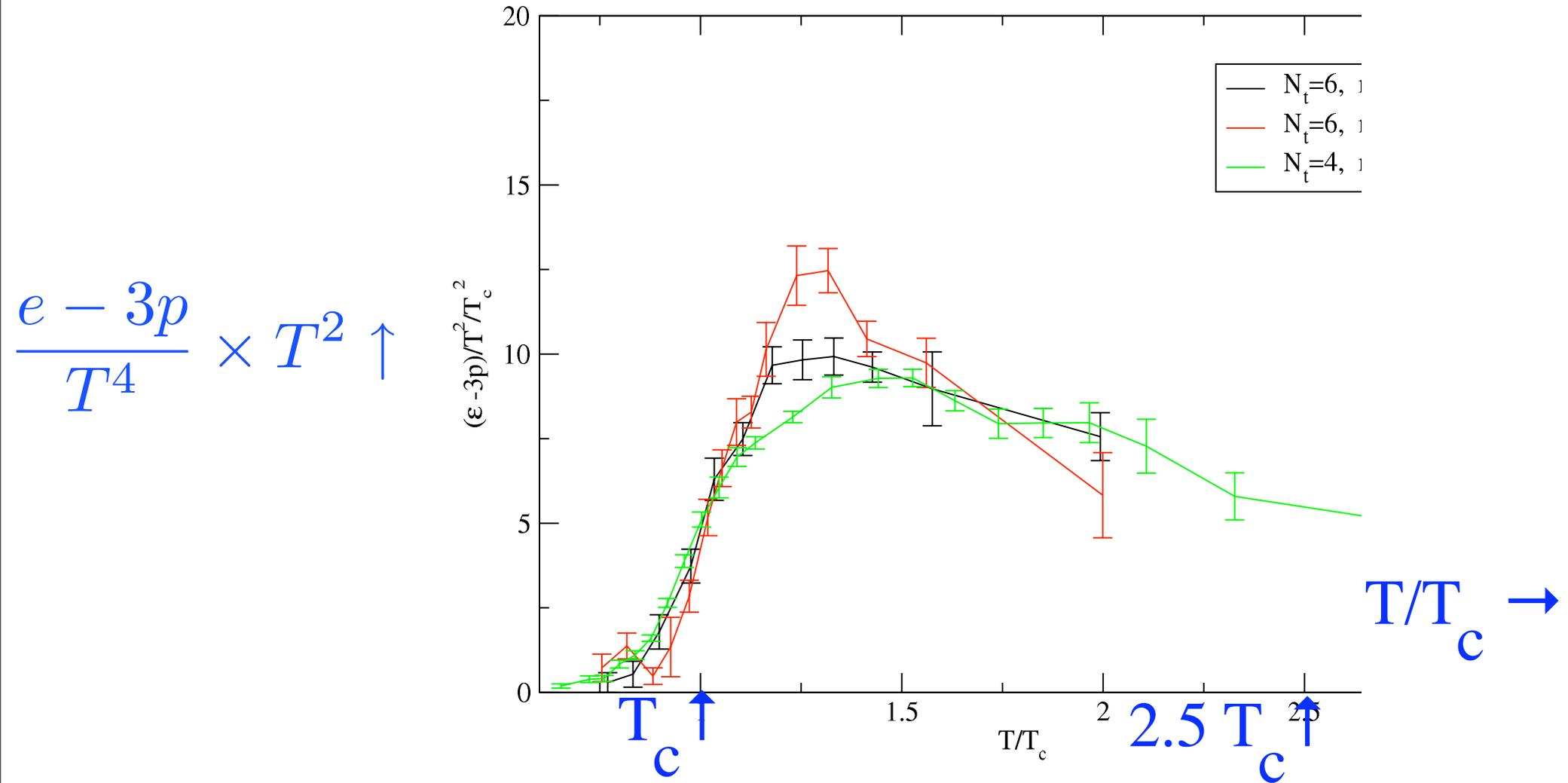
From MILC collaboration: L. Lefkova, private comm.

Peak in  $(e-3p)/T^4$  at  $\sim 1.1 T_c$ . *Much* broader peak, both above and below  $T_c$ .



# Lattice: SU(3) glue *with* quarks

2+1 dynamical flavors



Not just  $T^2$  term in pressure: also “MIT” bag term.

# “Generalized” bag model

Above maximum in  $(e-3p)/T^4$ :  $T > \sim 1.1 T_c$

pressure =  $T^4$  times powers series in  $1/T^2$ :

$$p(T) = f_{pert} T^4 - B_{fuzzy} T^2 - B_{MIT} + \dots$$

$B_{fuzzy}$  “fuzzy” bag constant: dominates MIT bag constant,  $B_{MIT}$ , away from  $T_c$

Resummations of perturbation theory contribute only to  $f_{pert}(g^2)$ :  
should work down to  $T_c$  without difficulty.

Perturbation theory fails because of *non*-perturbative terms, powers in  $1/T^2$ .

For hydro.: corrections to ideality are not  $1/T^4$ , as in MIT bag model, but  $1/T^2$ .

Does it matter much? Maybe not. But it *is* what the lattice shows.

## “Critical” region: $T: 0.9 T_c \rightarrow 1.1 T_c$

What about *below* peak in  $(e-3p)/T^4$ ? Say,  $T: 0.9 T_c \rightarrow 1.1 T_c$ .

While  $(e-3p)/T^4$  is big, in absolute terms, the pressure  $(p/T^4)$  is *small*.

For pure glue, *one* light mode about  $T_c$ : Polyakov loop.

For QCD, at least *four* light modes:

- Polyakov loop

- 4 modes for chiral symmetry restoration: sigma, + 3 pi's.

- Light eta? (if axial U(1) restored)

Suggestion: model with one light mode (+ heavy modes to avoid transitions)

Vary how light mass gets: will change minimum in speed of sound at  $T_c$ .

Will change hydro. results dramatically.