

Exploring Quark Gluon Plasma with Realistic Hydrodynamic Evolution

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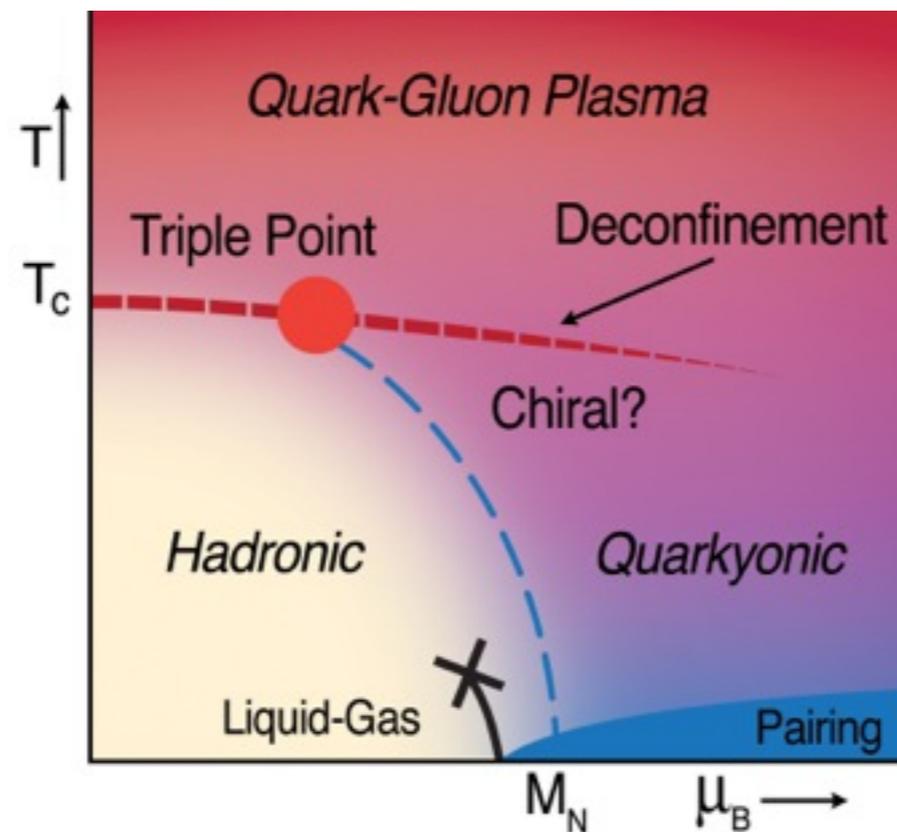


- ◆ Ho-ung Yee and YY, arXiv:1311.2574
- ◆ YY, arXiv:1312.4434

Nuclear theory seminar, BNL, Feb. 28th .

Quark-Gluon Plasma

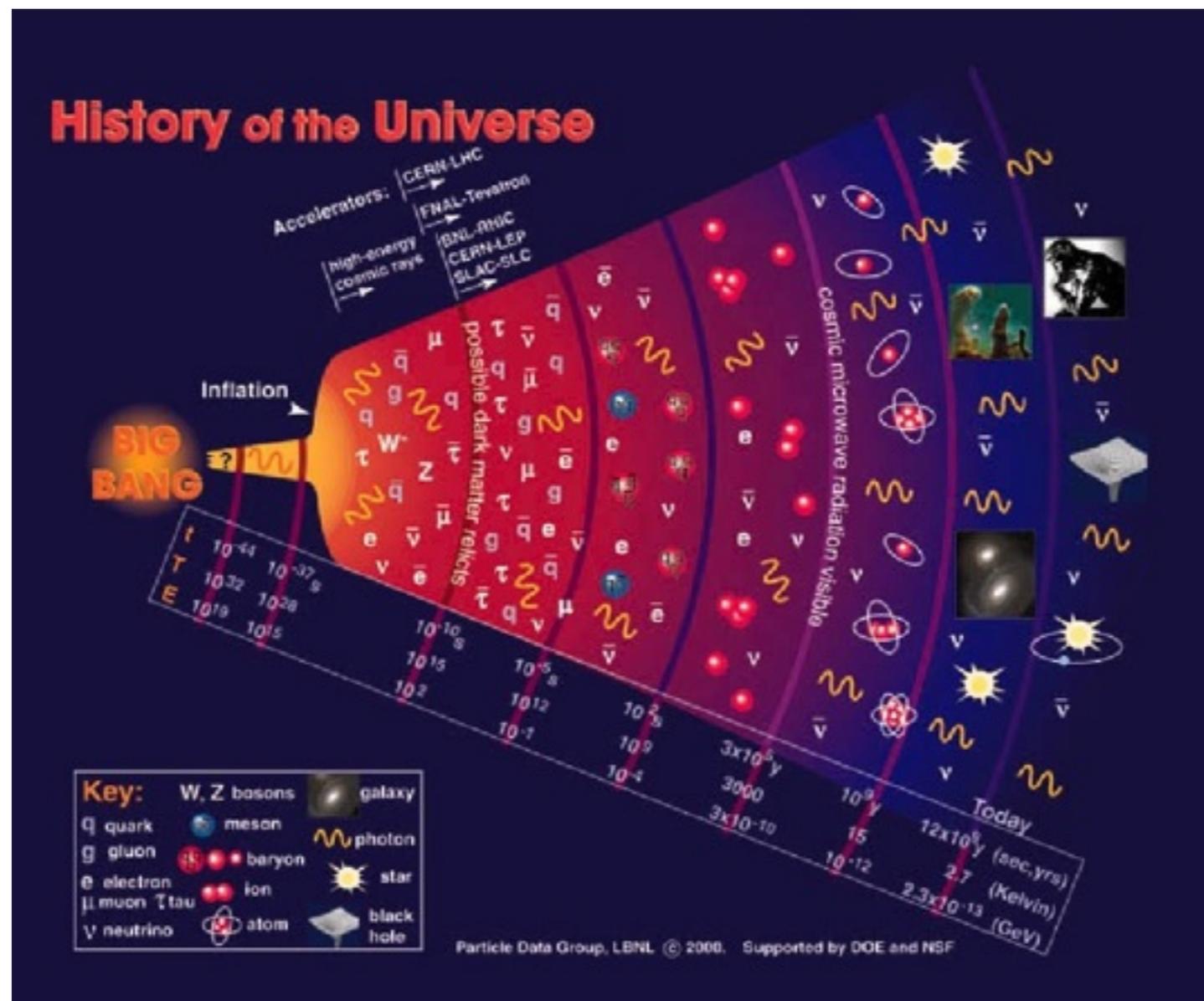
- Quark-gluon plasma(QGP) is a phase of QCD matter with temperature above transition temperature T_c .



- Where to find quark-gluon plasma?

Where to find Quark-Gluon Plasma(QGP)? Option I:

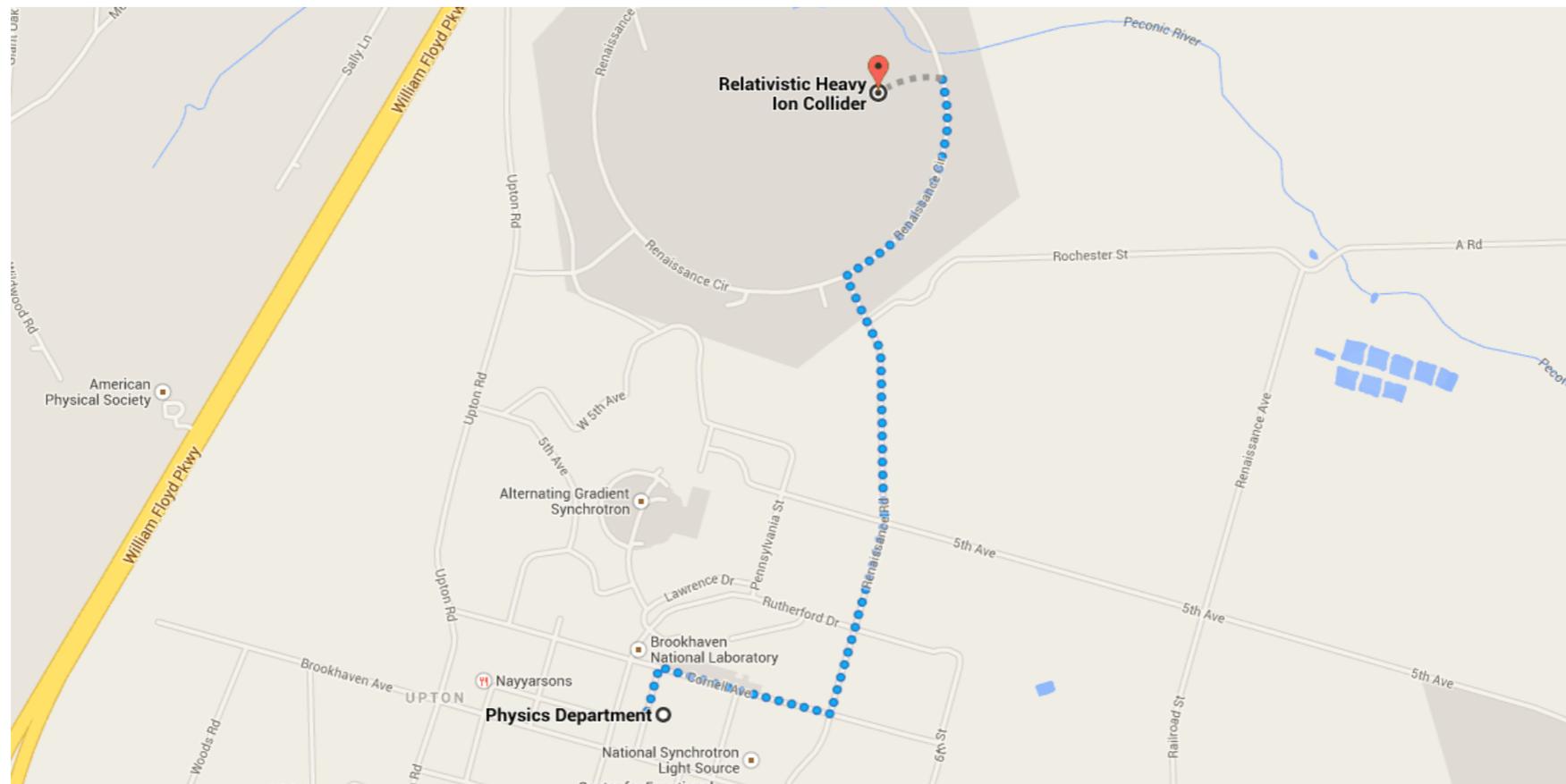
- Traveling back to the early universe.



Where to find Quark-Gluon Plasma(QGP)? Option II:

Directions from Physics Department to Relativistic Heavy Ion Collider

Walk 1.6 mi, 30 min



○ Physics Department
Pennsylvania St, Upton, NY 11973

Use caution - may involve errors or sections not suited for walking

 This route has restricted usage or private roads.



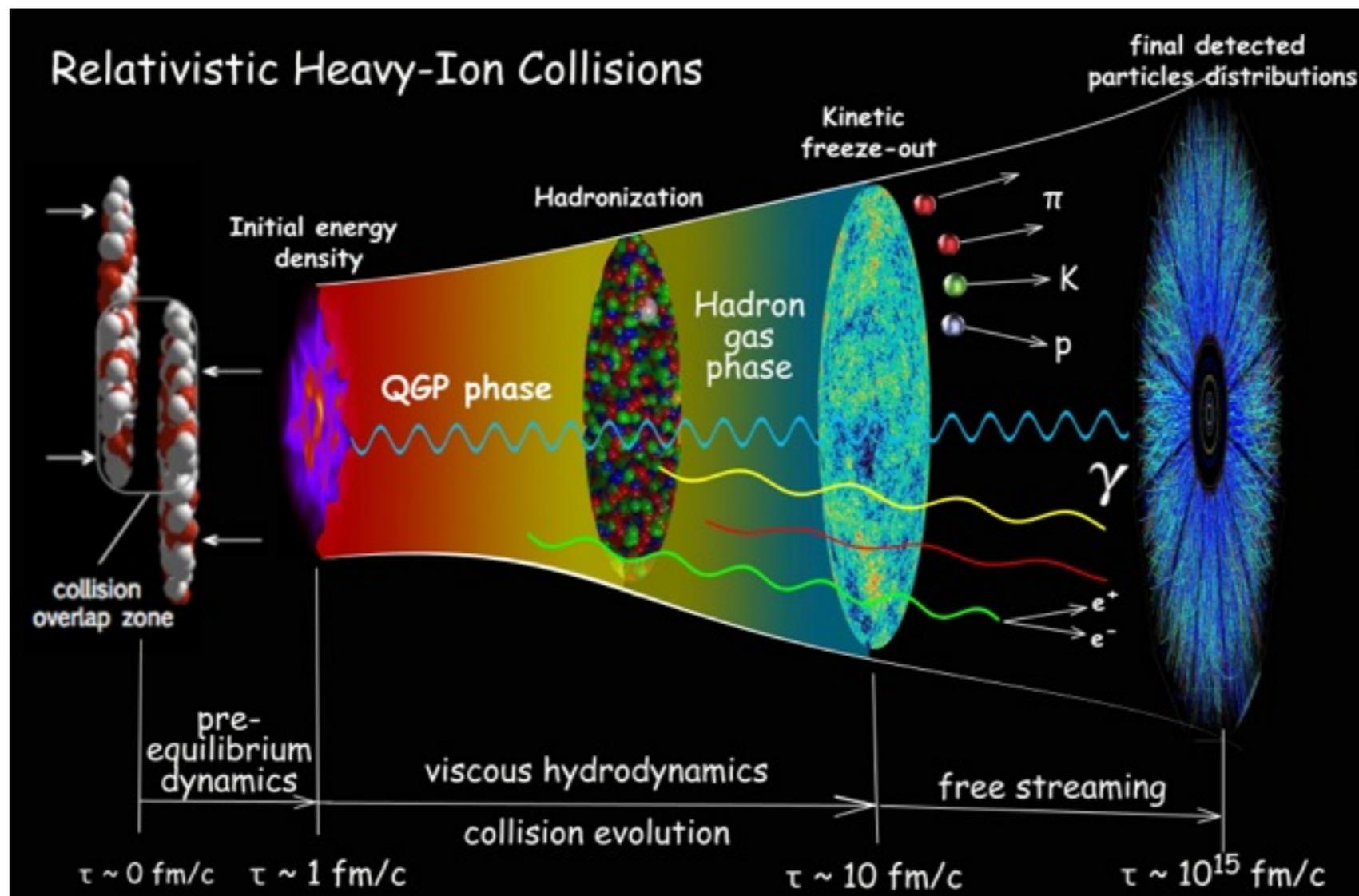
1. Head north on Pennsylvania St toward Cornell Ave

 Restricted usage road

417 ft

Heavy-Ion collisions and hydrodynamics

- The thermal equilibrium stage of heavy-ion collisions is well-described by relativistic hydrodynamics.



Hydrodynamics and heavy-ion collisions

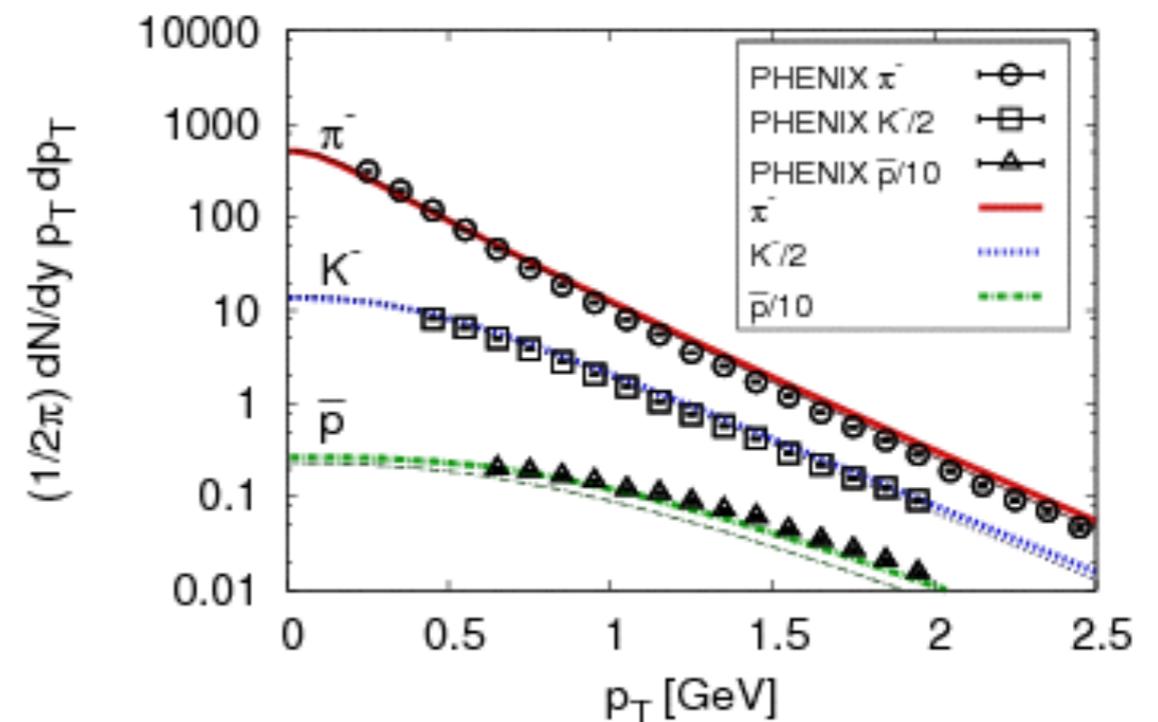
- Hydrodynamics:

- Low energy and low momentum effective theory.
- Degree of freedom: conserved quantities such as energy density, momentum density, charge density.

- Dynamics: conservation equations.

$$\partial_\mu T^{\mu\nu} = 0 \quad \partial_\mu J^\mu = 0$$

- Relativistic hydrodynamic simulations provide a decent description of particle spectrum in the experiment.



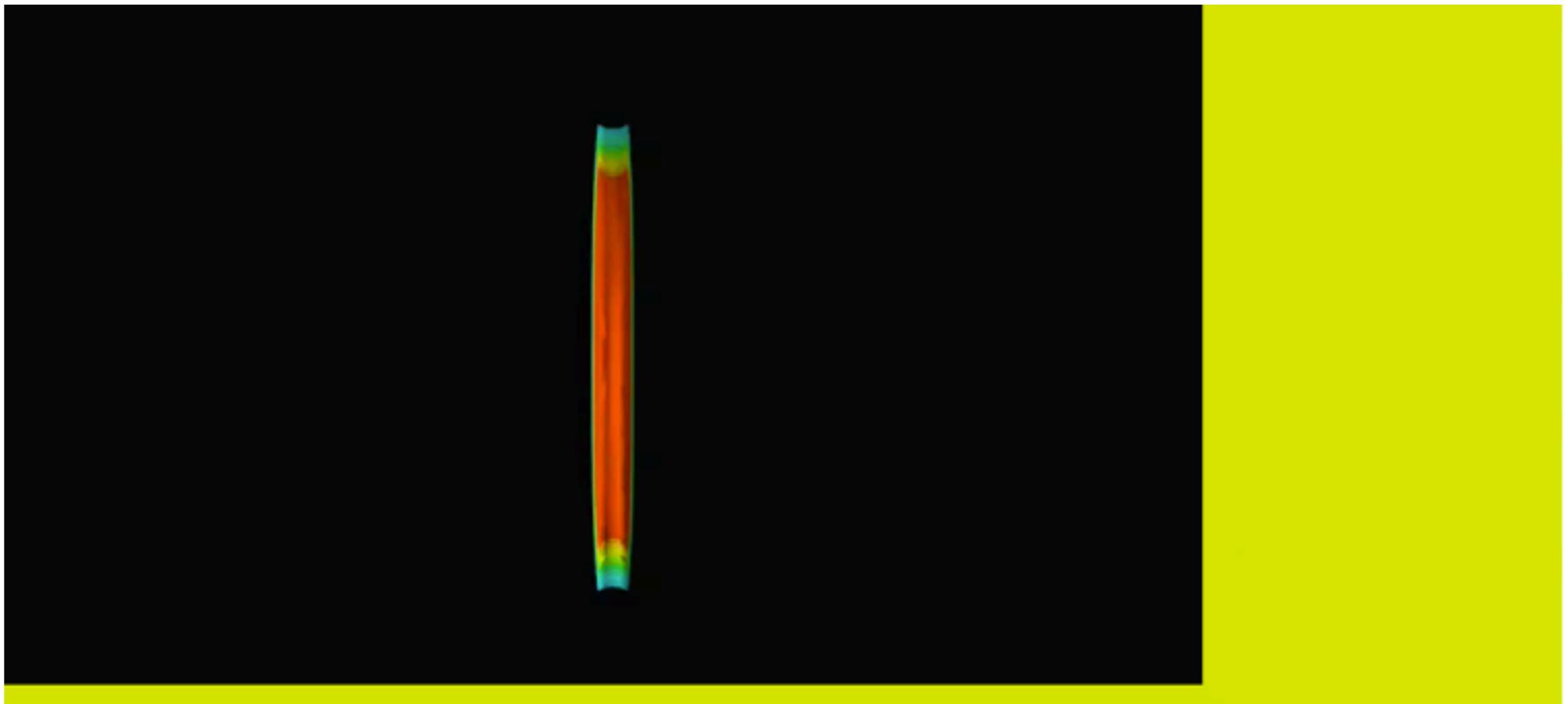
(Results of 3d hydro. simulations by McGill group)

We know hydrodynamic evolution of the fireball !

- Given a fluid cell labelled by its temporal and spatial coordinates, its temperature flow profile are known:

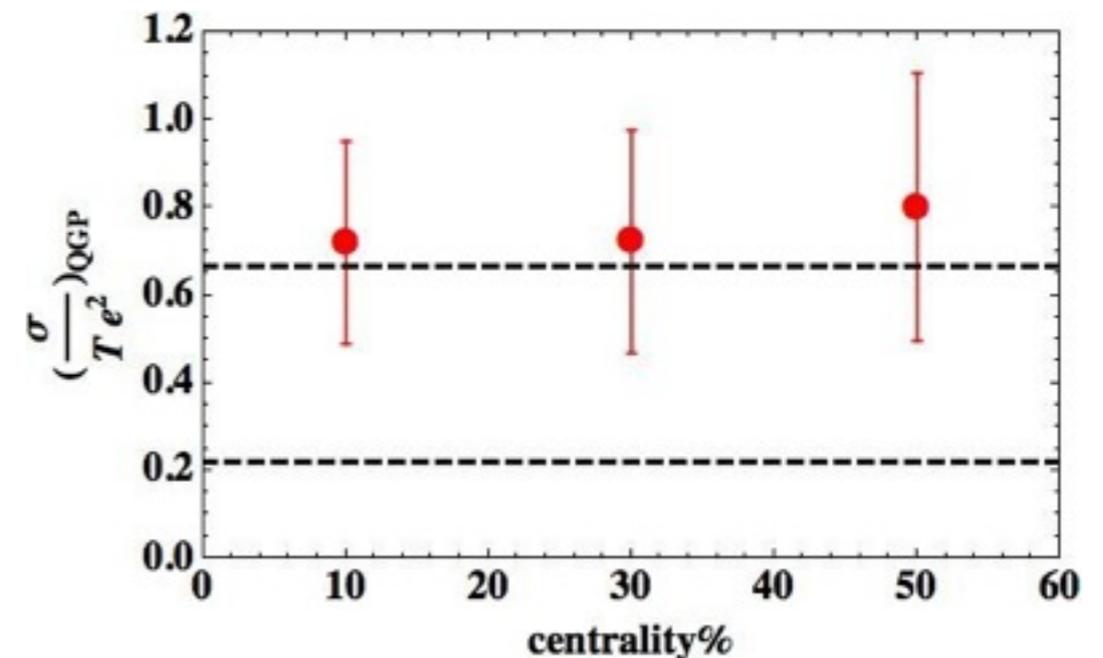
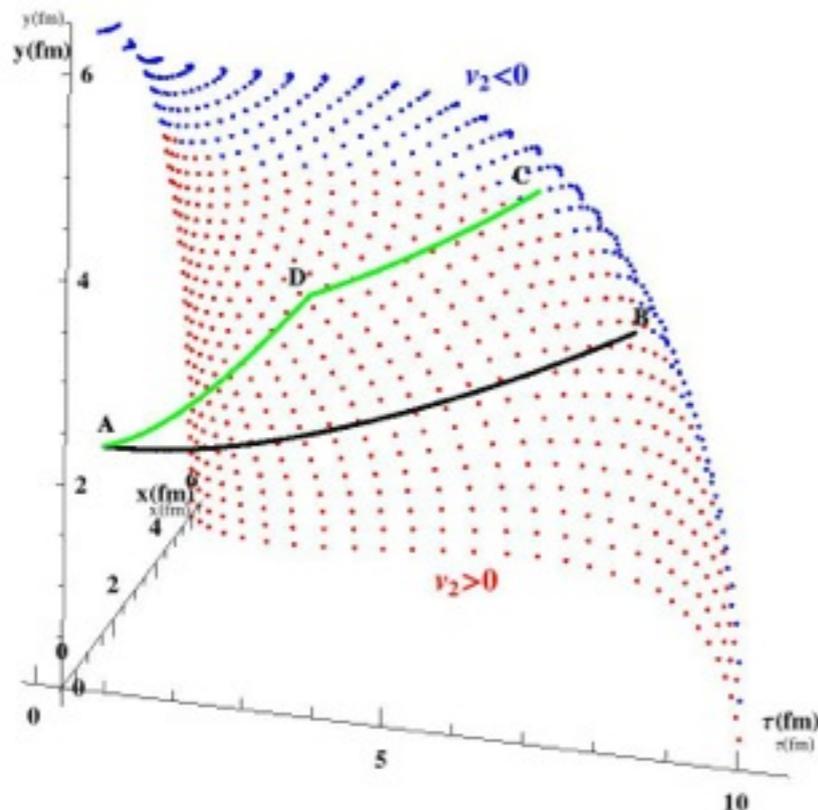
$$T(t, \vec{x}), u^\mu(t, \vec{x})$$

(Animation by Bjoern)



Exploring properties of QGP with realistic hydrodynamic background

- Taking temperature-flow profile as known background, how much information one could infer from the experiment?
- This talk: two examples.
 - ◆ Story I: Effects of chiral anomaly
 - ◆ Story II: Transport properties of QGP



Story I: chiral anomaly in heavy-ion collisions

Chiral anomaly and hydrodynamics

- The constitutive relation of a chiral fluid is modified by chiral anomaly:

$$J_V^\mu = n_V u^\mu + \Delta J_{V,\text{anom}}^\mu + \text{dissipative terms}$$

$$\Delta J_{V,\text{anom}}^\mu = \frac{\mu_A}{2\pi^2} B^\mu \quad (B^\mu = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} u_\nu F_{\alpha\beta}) \quad \partial_\mu J_A^\mu = \frac{1}{2\pi^2} \vec{E} \cdot \vec{B}$$

- Nielsen-Ninomia argument(1983):

$$\mu_A \frac{dn_A}{dt} = \frac{\mu_A}{2\pi^2} \vec{E} \cdot \vec{B} = \vec{E} \cdot \Delta \vec{J}_{V,\text{anom}}$$

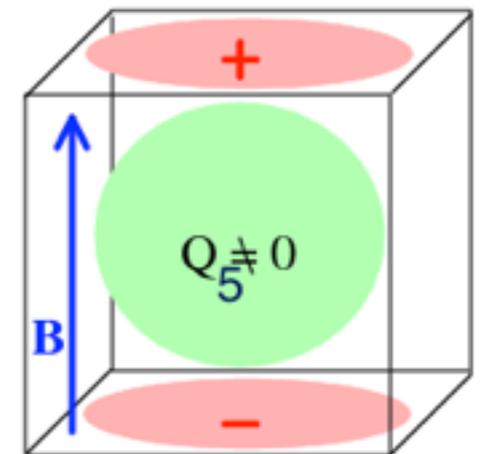
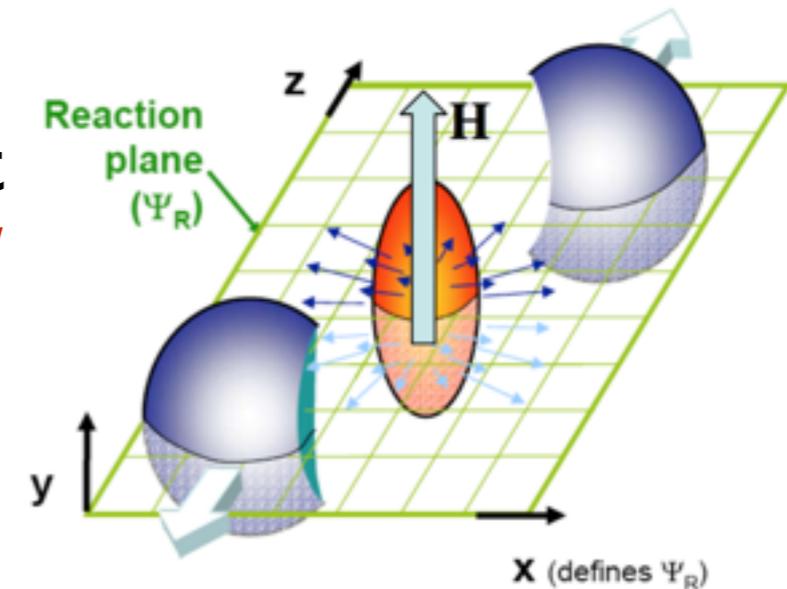
Energy = Work

- Microscopically: effective velocity of a chiral fermi is modified by Berry curvature(e.g. M. Stephanov,YY PRL 2012).

- Similarly: $\Delta J_{A,\text{anom}}^\mu = \frac{\mu_V}{2\pi^2} B^\mu$

Chiral anomaly and heavy-ion collisions

- Basic ingredients for the anomalous current induced by anomaly: **magnetic field**, and **axial/vector charge density**.
- **Strong magnetic field** is created in heavy-ion collisions. $eB \sim m_\pi^2 \sim 10^{17}$ Gauss
- **Axial charges** are generated by topological fluctuations.
- Experimentalist could select events with non-zero **vector charges**.



Anomaly induced charge asymmetry

- Anomalous current leads to charge quadrupole.

$$\Delta J_{R,\text{anom}}^\mu = \frac{\mu_R}{4\pi^2} B^\mu \quad \Delta J_{L,\text{anom}}^\mu = -\frac{\mu_L}{4\pi^2} B^\mu$$

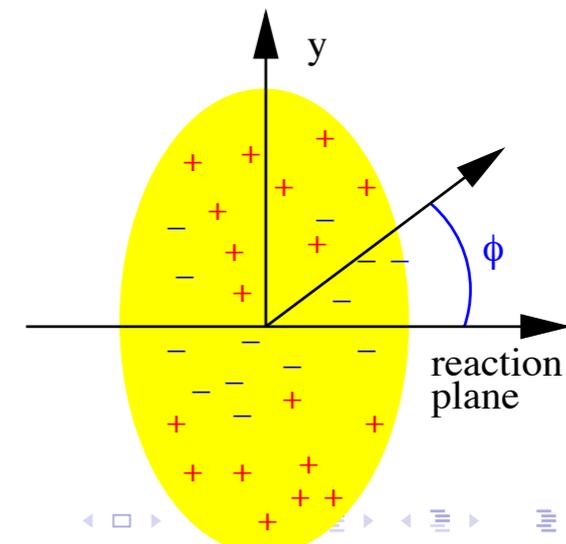
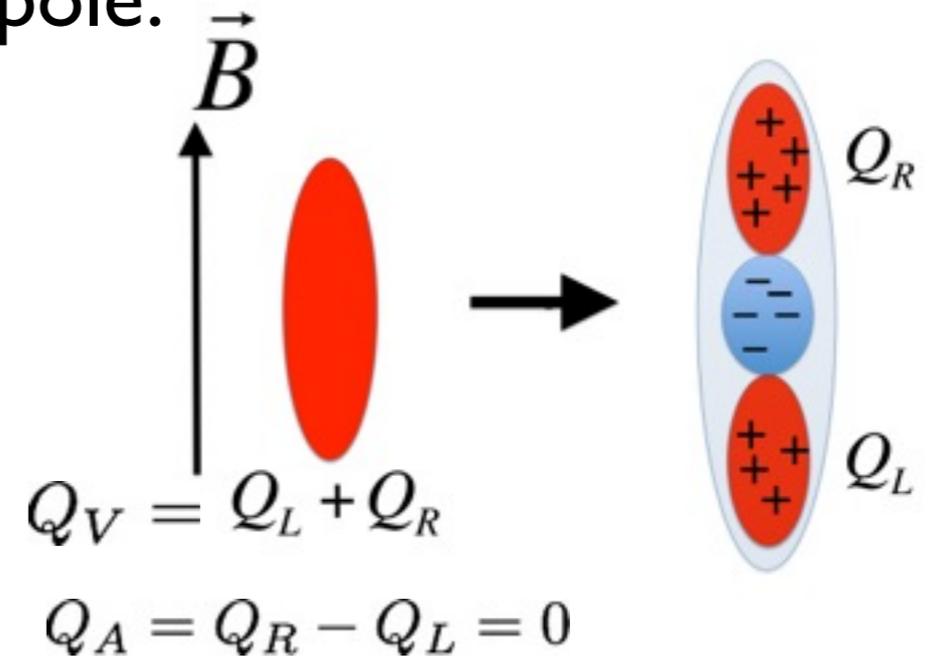
- There will be more positive charges around the polar than around equator.

Around the polar: $\cos(2\phi) < 0$

Around the equator: $\cos(2\phi) > 0$

- Quadrupole distribution leads to charge asymmetry in pion elliptic flow:

$$v_2^-(\pi) > v_2^+(\pi)$$



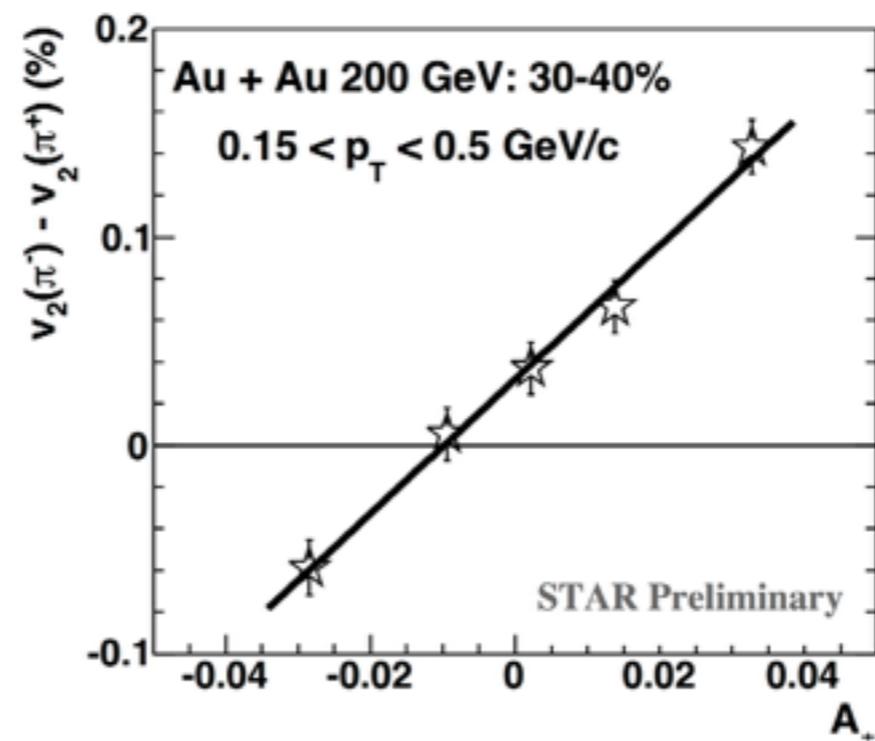
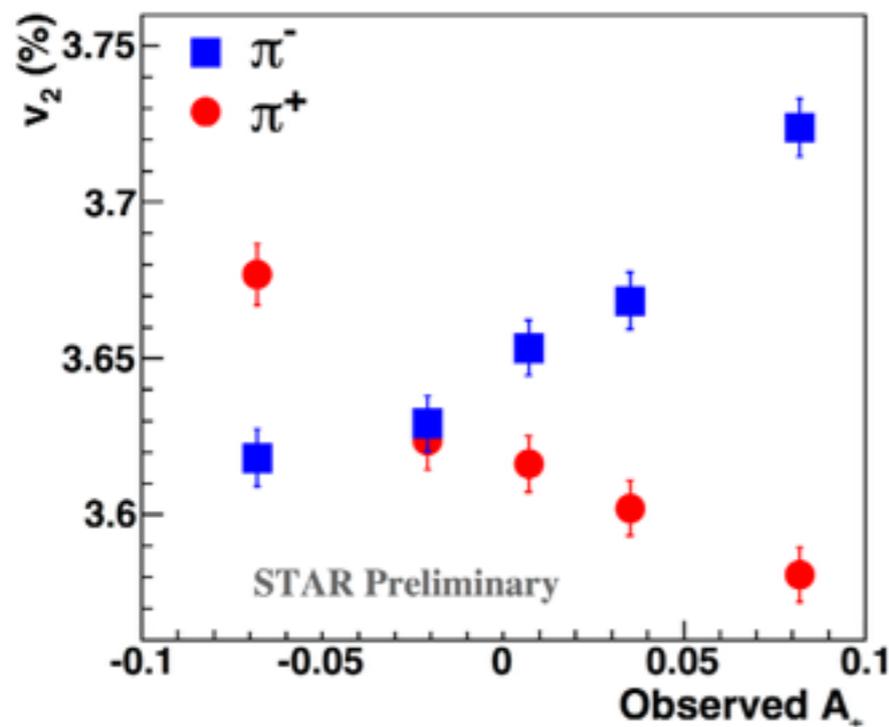
Charge-dependent in experiment

- In experiment, $v_2^-(\pi) - v_2^+(\pi) > 0$ and the difference is approximately proportional to charge asymmetry.

$$\frac{d^3 N}{p_\perp dp_\perp d\phi_p dY} \Big|_{Y=0} = N^\pm(p_\perp) (1 + 2v_2^\pm(p_\perp) \cos(2\phi_p) + \dots)$$

$$v_2^- - v_2^+ = r A_\pm \quad A_\pm = \frac{N^+ - N^-}{N^+ + N^-} \quad (\text{relative charge difference})$$

(Slope parameter)



How much would chiral anomaly contribute to the data?

- The effects of anomaly is qualitatively consistent with the data.
- **Quantitatively**, how much would chiral anomaly contribute to the data?
- Two groups have opposite conclusions:
 - Burnier-Kharzeev-Liao-Yee(2012), solving anomalous equation in static temperature and flow background.
 - Hongo-Hirono-Hirano(2013), solving full set of hydro. equation with separating the initial effects and the effects of evolution.
- To resolve the controversial, a realistic computation is imperative.

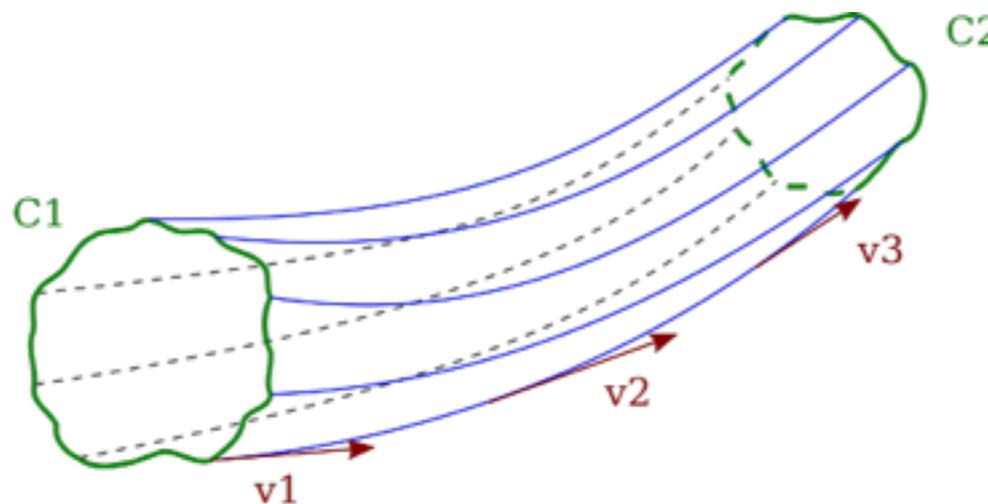
Realistic implementation of chiral charge transportation

- Ho-ung, Yee and I revisited this problem(1311.2574). The slope parameter can be interpreted as a response to charge asymmetry. We then treat charge density as a probe.
- The back-reaction from charge density to the temperature flow profile can be neglected. We solve anomalous hydrodynamic equation linearized around the realistic hydrodynamic background. (hydro. background from P&R Romatschke).
- We choose the initial condition in such way that the slope parameter is zero in the absence of chiral anomaly.
- A lot of numerics? No if we formulate the problem in a physical intuitive way.

Ideal hydrodynamics without anomaly and streamline

- What does current conservation imply for a normal fluid?

$$\partial_{\mu} (J_{\text{norm}}^{\mu}) = \partial_{\mu} (n u^{\mu}) = u^{\mu} \partial_{\mu} (n) + n (\partial^{\mu} u_{\mu}) = 0$$



- Total charge is fixed in a **fluid cell** along the **streamline**. Change of charge density is compensated by the expansion of the volume.

Streamlines of anomalous hydrodynamics

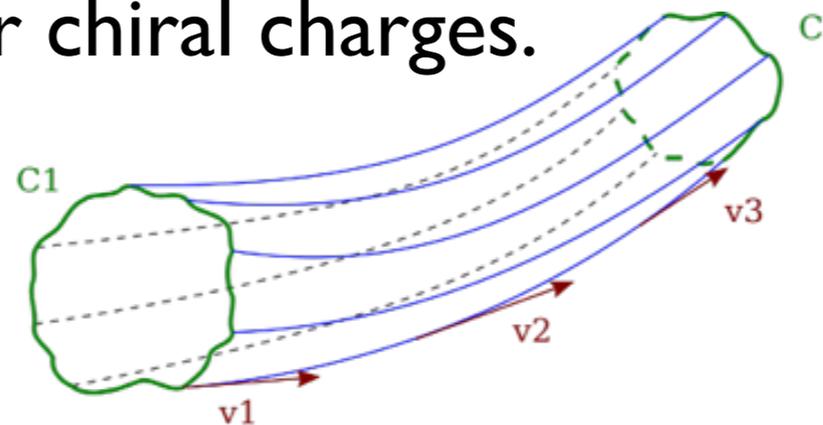
- Rewrite the constitute relation of chiral current similar to normal current as much as possible.

$$J_{L,R}^\mu = n_{L,R} u^\mu \mp \frac{qN_c}{4\pi^2} \mu_{L,R} B^\mu = n_{L,R} (u^\mu \mp u_{\text{anom}}^\mu) = \tilde{n}_{L,R} \tilde{u}_{L,R}^\mu$$

$$\partial_\mu J_{L,R}^\mu = 0 \qquad u_{\text{anom}}^\mu = \frac{qN_c B^\mu}{4\pi\chi} \quad (n_{L,R} = \chi\mu_{L,R})$$

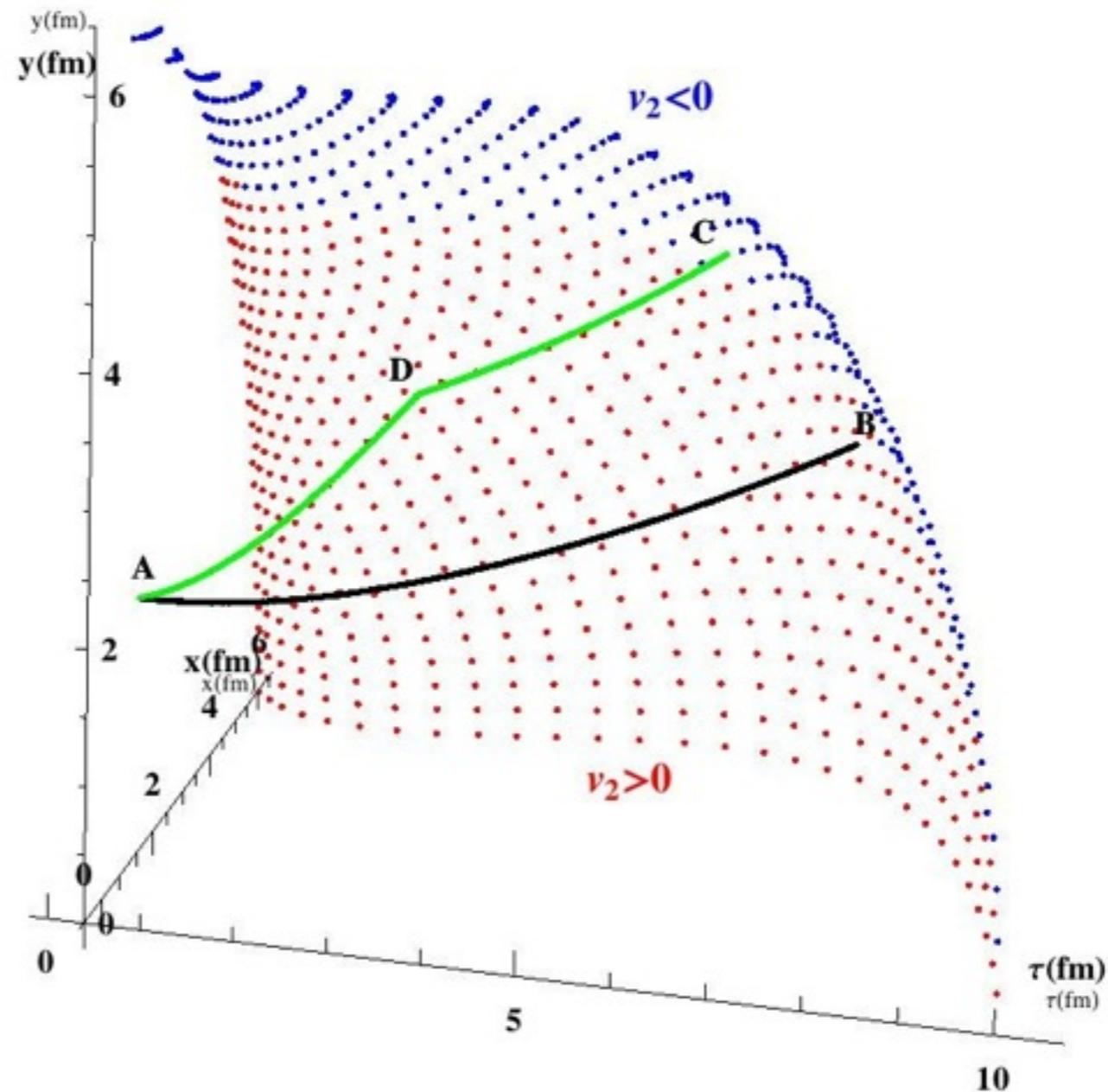
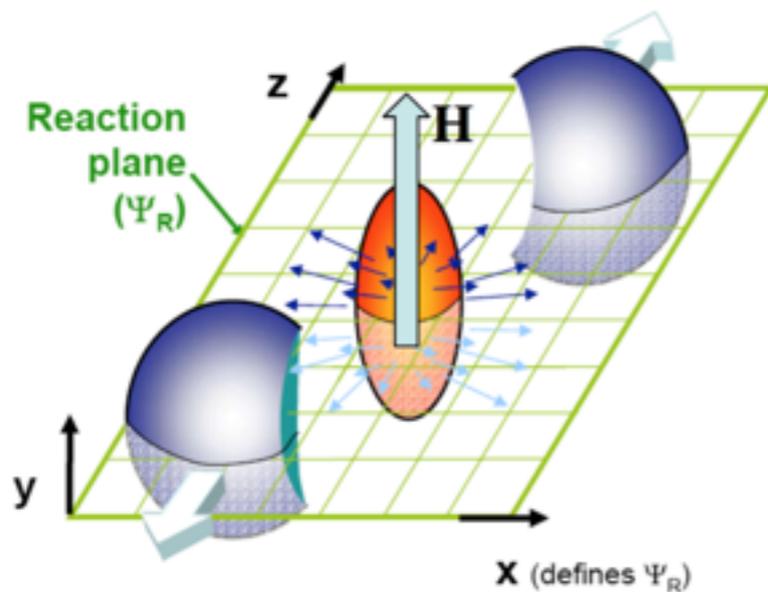
- Solving anomalous hydrodynamic equation is translated into the determination of streamline for chiral charges.

$$\frac{dX^\mu(\lambda)}{d\lambda} = \tilde{u}_{L,R}^\mu(X^\mu)$$



- The streamlines for left and right charge are **different** due to chiral anomaly.

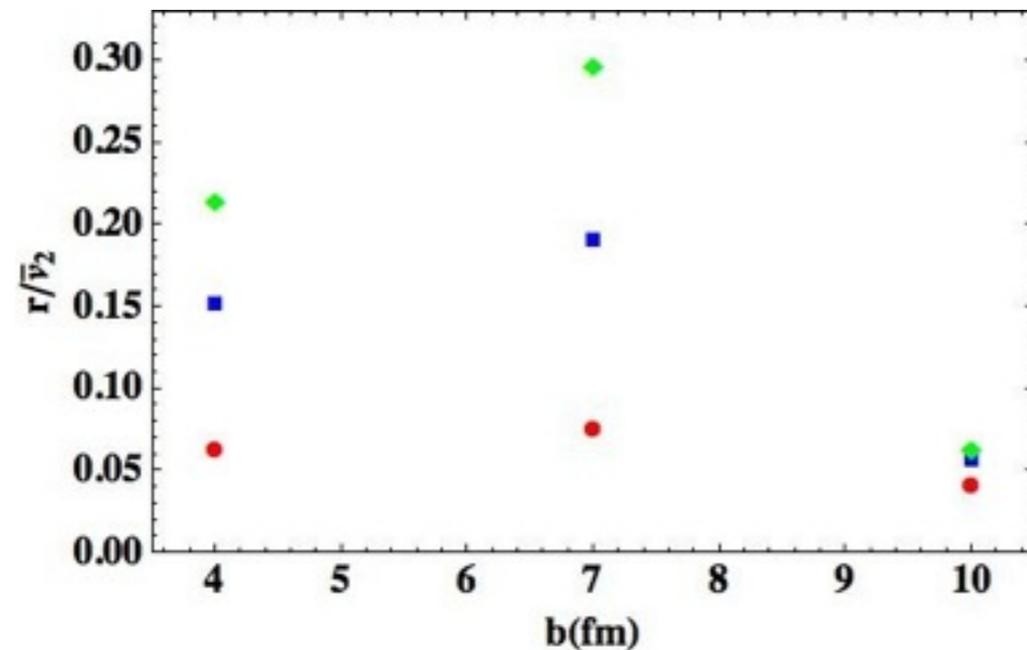
Streamline in heavy-ion collisions



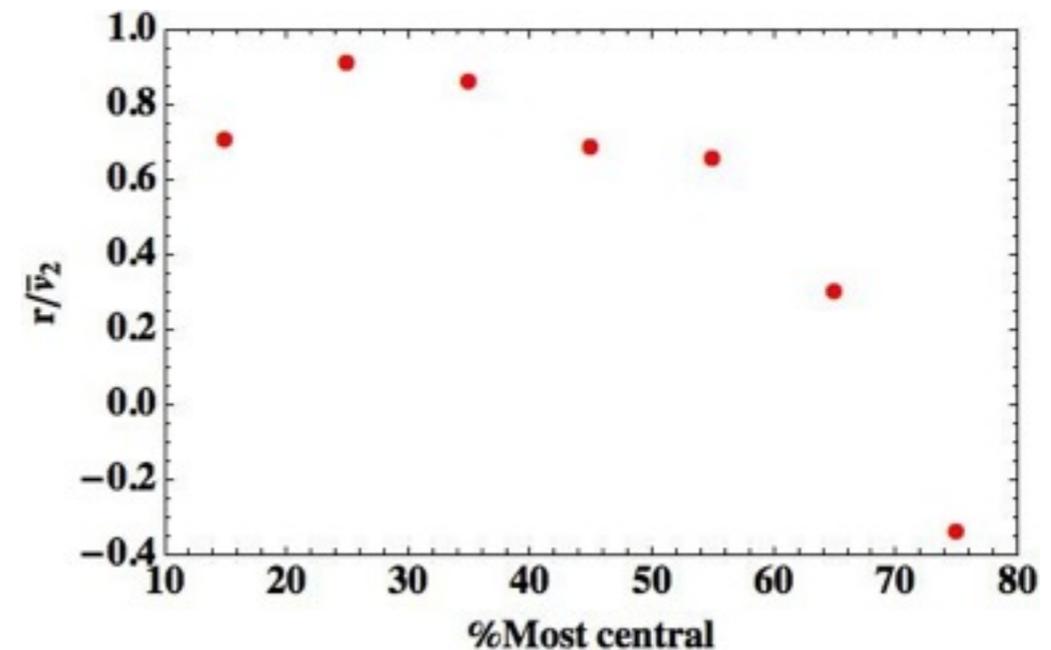
- Anomalous fluid velocity changes the streamline of chiral charges. This leads to $v_2^-(\pi) - v_2^+(\pi) > 0$.

Results

- Chiral anomaly will give sizable contribution to the data if life time of magnetic field is a few fm.



Red, blue, green dots are corresponding to
 $\tau_B = 2, 4, 6$ fm



STAR data

- Otherwise, it is more likely that the difference in elliptic flow it is due to initial distribution. (See Bzdak-Bozek(2013) for other sources of contributions)

Summary of Story I:

- We have developed a physically intuitive method to solve anomalous hydrodynamics in realistic background.
- As an application, we have studied charged dependent elliptic in experiment.
- The method based streamline gives a transparent picture on the chiral charge transportation in heavy-ion collisions.
- We could solve the anomalous hydrodynamic equation with arbitrary initial conditions.

Quantification of chiral magnetic effect

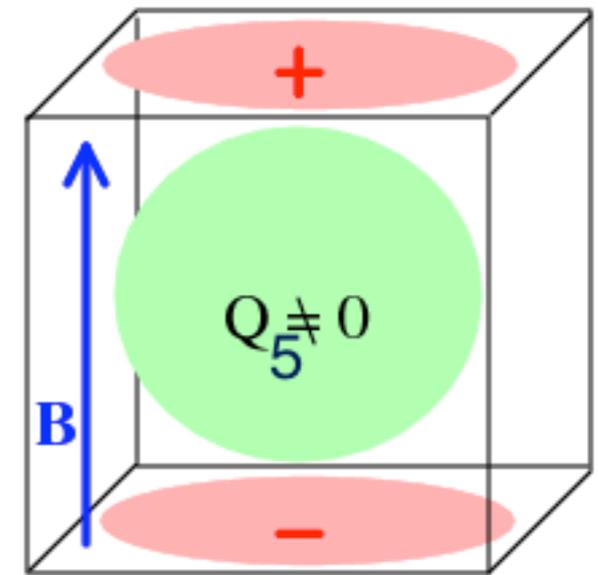
- Axial charge created in heavy-ion collisions will lead to vector current(chiral magnetic effect) thus charge dipole.

$$\Delta J_{V,\text{anom}}^\mu = \frac{\mu_A}{2\pi^2} B^\mu$$

- Charge dipole implies non-trivial correlations between particle pairs(Kharzeev-McLerran-Warringa).

$$\langle \cos(\phi_{1,+} + \phi_{2,+} - 2\Psi_{\text{RP}}) \rangle < 0$$

$$\langle \cos(\phi_{1,+} + \phi_{2,-} - 2\Psi_{\text{RP}}) \rangle > 0$$

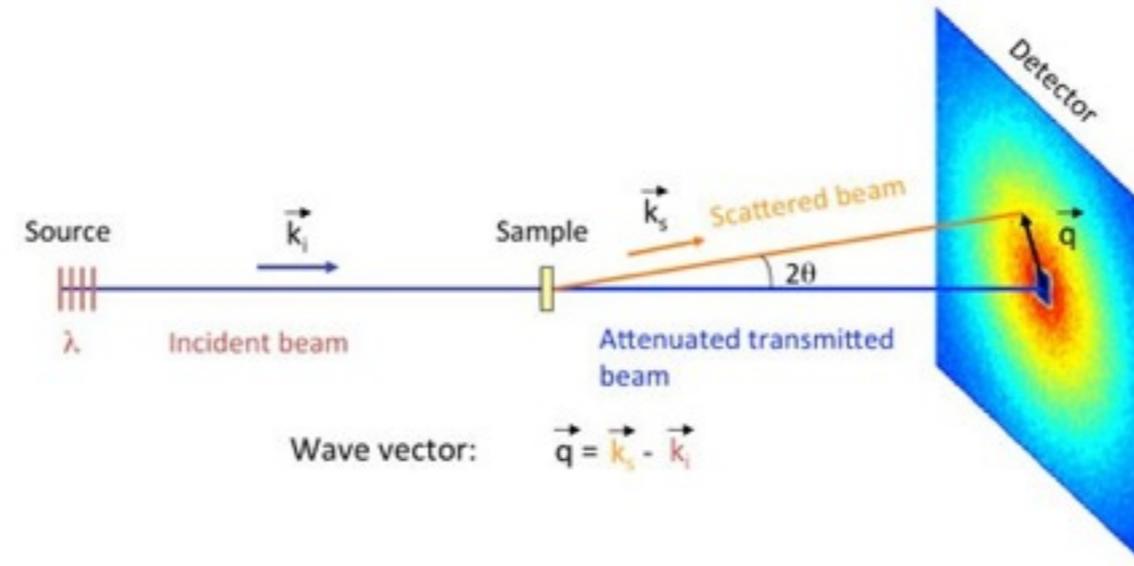


- Due to large background effects, quantify such effects in realistic background is imperative(in progress with Jinfeng, Liao).

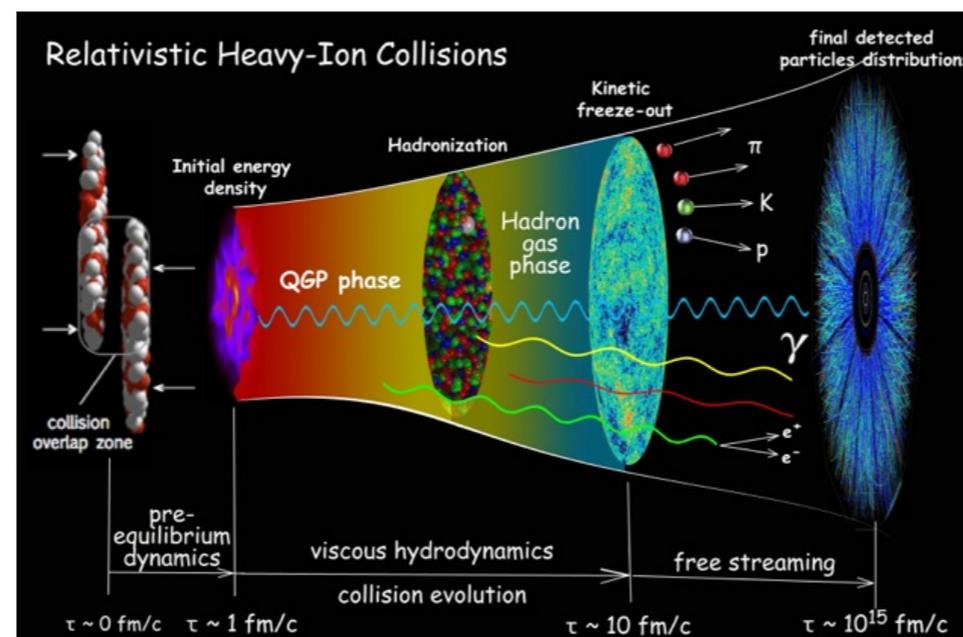
Story II: Electrical conductivity of QGP from photon production

Photon as a probe of QGP

- The properties of condensed matter can be measured by scatterings experiments.



- Quark-gluon plasma is radiating. Photon produced in heavy-ion collisions is a probe of QGP. ($\alpha_{EM} \ll \alpha_s$)



Photon as a probe of QGP

- Photon produced per volume per time is related to Green's function of QGP:

$$\omega \frac{d\Gamma_\gamma}{d^3\vec{p}} = -\frac{n_B(\omega/T)}{(2\pi)^3} P_T^{ij} \text{Im} [G_{ij}^R(\omega = |\vec{p}|)]$$

$$P_T^{ij} = \delta^{ij} - \hat{p}^i \hat{p}^j$$

Retarded correlator: $G_R^{\mu\nu} \sim \langle J^\mu J^\nu \rangle$

- Photons are produced during the full evolution of the fireball with shifted frequency:

$$\frac{dN_\gamma}{p_t dp_t d\phi_p dY} = \int_{T \geq T_f} d^4x \omega_{\text{shift}} \frac{d\Gamma_\gamma}{d^3p'} \Big|_{\omega_{\text{shift}} = |\vec{p}|}$$

$$\omega_{\text{shift}} = p_{\text{Lab}}^\mu u_\mu$$

Photon frequency is shifted in the rest frame of fluid.

- To study photon spectrum, we need i) theoretical understanding of photon rate and ii) hydrodynamic evolution.

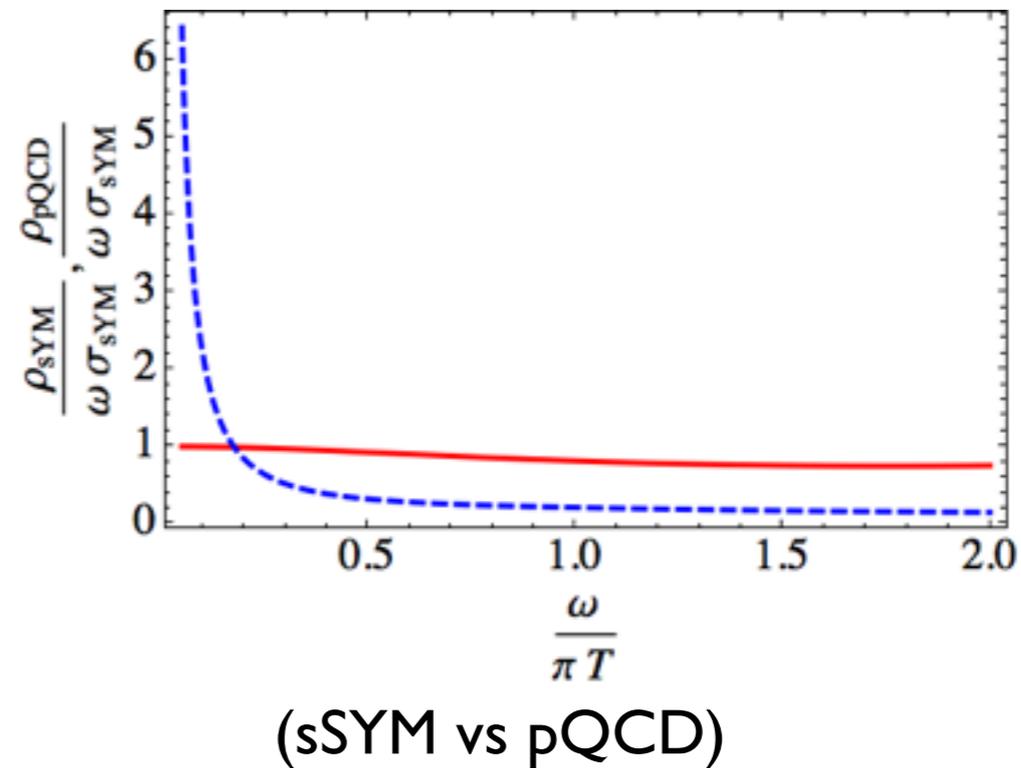
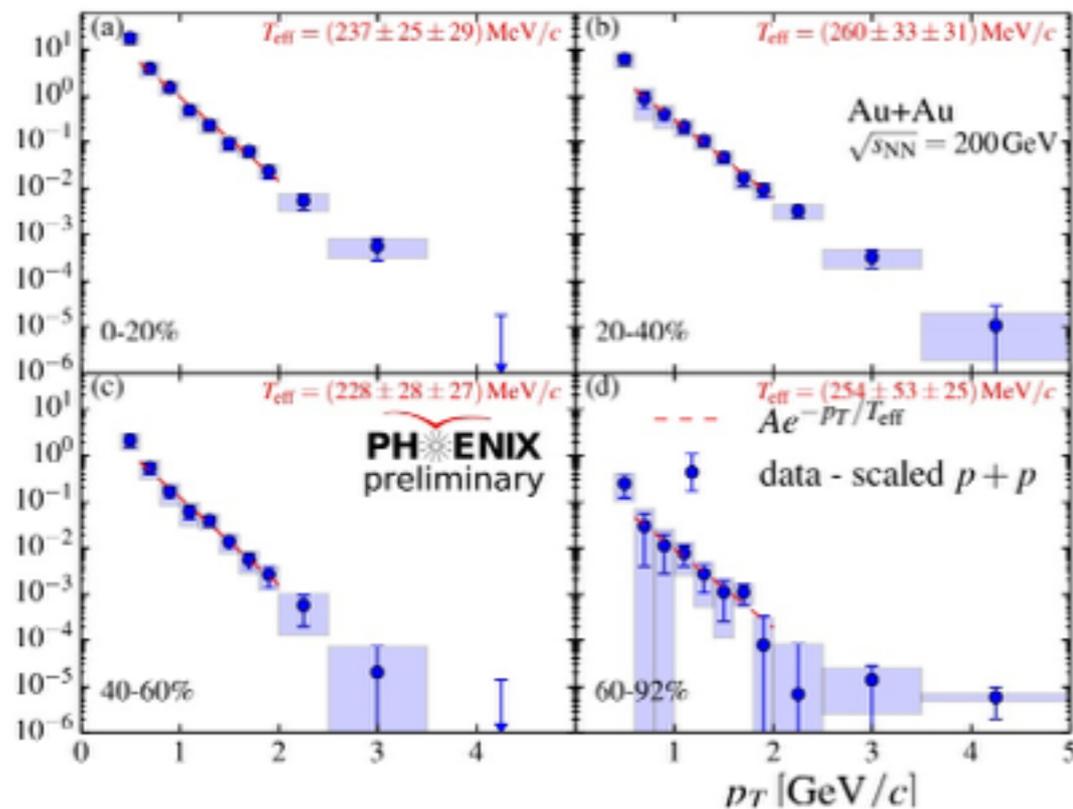
Photon and thermal correlators

- Photon production in heavy-ion collisions has been studied by evolving pQCD photon rate with hydrodynamic simulations.
- QGP is strongly coupled! pQCD rate may not be applicable for photon energy below a few GeV!
- Determination of correlation functions from microscopic theory is challenging.
- In low frequency limit, a macroscopic description is possible!

$$P_T^{ij} \text{Im} [G_{ij}^R(\omega = |\vec{p}|)_{B=0}] = \frac{\alpha_{\text{EM}} \omega \sigma_0}{\pi^2 e^2}$$

Soft Photon and conductivity

- Lowest pt in experiment: 0.5 GeV. (by PHENIX) $\pi T_{QGP} \approx 1\text{GeV}$



- Strongly interacting system has a wider hydrodynamic regime!
- We will use the hydrodynamic approximation to study the lowest pt photon data

$$P_T^{ij} \text{Im} [G_{ij}^R(\omega = |\vec{p}|)]_{B=0} = \frac{\alpha_{\text{EM}} \omega \sigma_0}{\pi^2 e^2}$$

Electrical conductivity of QGP

- We write down the rate

$$\frac{dN_\gamma}{2\pi p_t dp_t dY} = \frac{\alpha_{\text{EM}}}{\pi^2} \int_0^{2\pi} \frac{d\phi_p}{2\pi} \int_{T \geq T_f} d^4x \frac{\omega_{\text{shift}} \sigma}{\exp(\omega_{\text{shift}}/T) - 1} \quad \omega_{\text{shift}} = p_{\text{Lab}}^\mu u_\mu$$

- We estimate conductivity at QGP temperature by computing the following ratio.

$$\left\langle \frac{\sigma}{e^2 T} \right\rangle \equiv \frac{\frac{dN_\gamma}{2\pi p_t dp_t dY}}{\frac{\alpha_{\text{EM}}}{\pi^2} \int_0^{2\pi} \frac{d\phi_p}{2\pi} \int_{T \geq T_f} d^4x \frac{\omega_{\text{shift}} T}{\exp(\omega_{\text{shift}}/T) - 1}}$$

← From the data

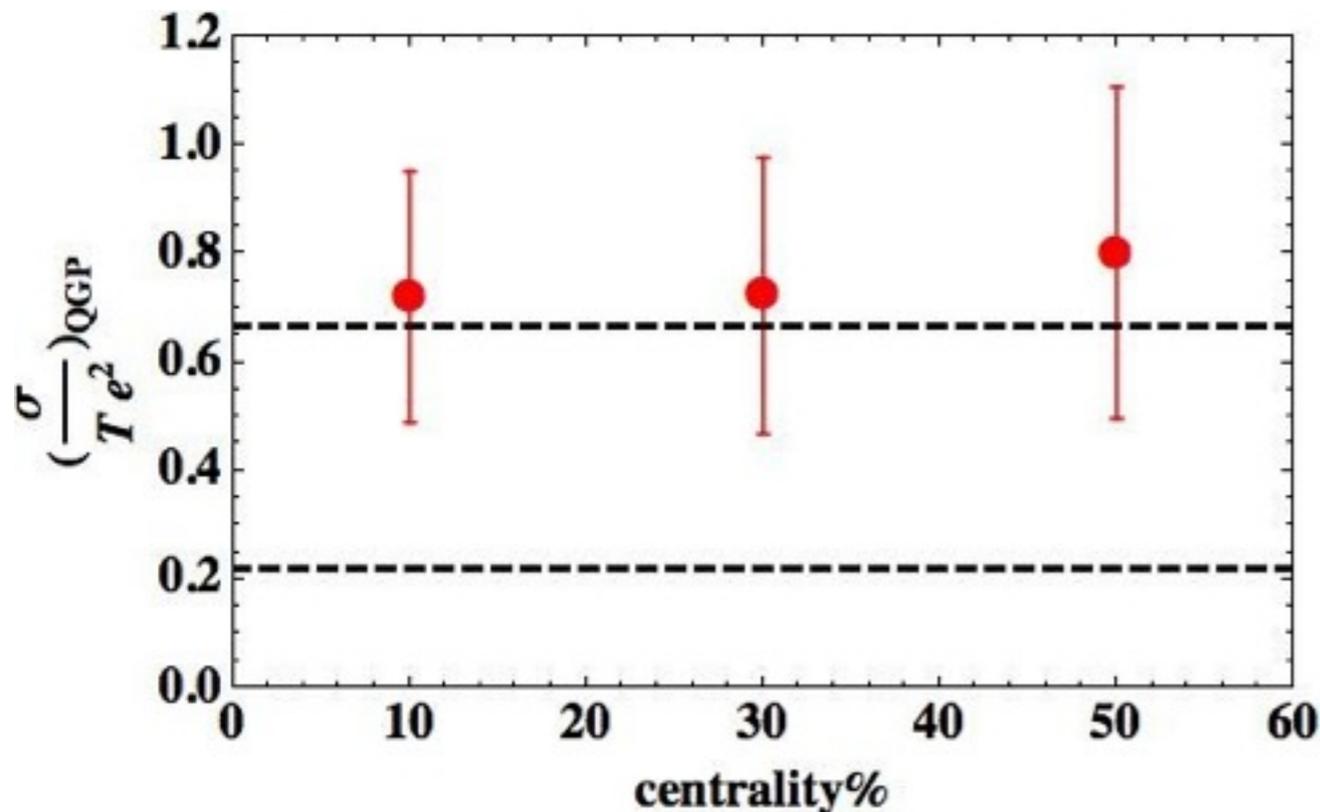
(NB: σ/T is dimensionless, similar to η/s .)

- We evolve the integral with realistic hydrodynamic background.

(The realistic hydrodynamic background is from Heinz' group, available online:
https://wiki.bnl.gov/TECHQM/index.php/Main_Page)

Electrical conductivity of QGP

- PHENIX data, different for **different** centralities.
- Hydrodynamic background, different for **different** centralities.
- The ratio has a **weak-dependence** on centralities! Conductivity is the properties of QGP!



$$\left\langle \frac{\sigma}{e^2 T} \right\rangle \equiv \frac{\frac{dN_\gamma}{2\pi p_t dp_t dY}}{\frac{\alpha_{\text{EM}}}{\pi^2} \int_0^{2\pi} \frac{d\phi_p}{2\pi} \int_{T \geq T_f} d^4x \frac{\omega_{\text{shift}} T}{\exp(\omega_{\text{shift}}/T) - 1}}$$

Lattice extraction by Bielefeld-BNL group(2011) at $T=1.45T_c$

Summary of Story II

- We have estimated electrical conductivity, for the first time in literature, from soft photon data.
- Thermal correlations function at other pt window? More information from dileptons?

The spherical cow

- Theorists always use highly simplified scientific models to study complex real phenomena.



The actual cow

- In heavy-ion collisions, studying the problems in realistic environment is important!



Conclusion and outlook

- We have shown two examples on exploring the properties of QGP with realistic hydrodynamic background.
 - Effects of chiral anomaly and pion elliptic flow.
 - conductivity and photon spectrum.
- More fruitful results in future.