



# A Unicorn at RHIC: the (s)QGP

At *very* high temperatures, expect hadrons (QCD) behaves qualitatively different:  
new type of matter, Quark Gluon Plasma (QGP).

To see experimentally, collide heavy ions (AA) at high energies:

Relativistic Heavy Ion Collider, RHIC at Brookhaven; SPS & LHC at CERN.

Today:

Theory of the QGP, especially numerical “experiments” on a lattice

RHIC: Suppression of jets in AA collisions. Robust signal of new physics.

Non central collisions exhibit strong “flow”

Even heavy (“charm”) quarks are suppressed, flow!

Not a perturbative QGP, instead “sQGP” (s?)

“Most perfect fluid on earth”?

Theory of Ads/CFT: hot “QCD” at infinite coupling.

Theorists all agog about a *wealth* of experimental data.

Myths can come true, but maybe not the ones you expect...



# QCD: Quantum ChromoDynamics

Gluons: spin-1, three colors, local SU(3) symmetry. Classically, *no* mass scale.

One dimensionless coupling,  $\alpha_s$ , which “runs” with momentum.

Theory *uniquely* specified by value of  $\alpha_s$  at *one* (finite) momentum.

Asymptotically Free (AF):  $\alpha_s$  *vanishes* at *infinite* momenta

$$\alpha_s(p) \sim \# / \log(p) , p \rightarrow \infty$$

Perturbative > 1-2 GeV

powers of  $\alpha_s \sim 1/\log(p)$

Non-perturbative < 1-2 GeV:

powers of (non-pert. mass)<sup>2</sup>/p<sup>2</sup>

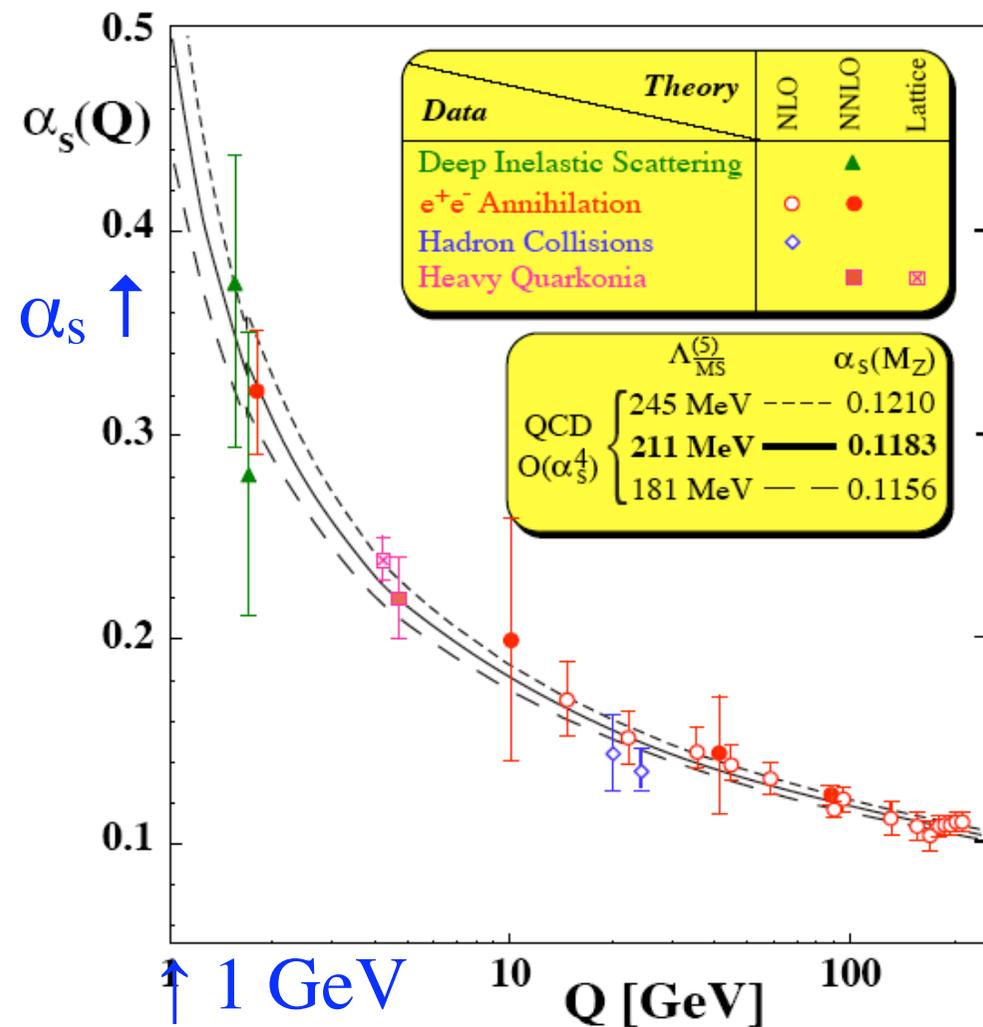
AF => *ideal gas* at *infinite* temperature.

Natural to normalize to ideal gas values.

Quarks: six “flavors”, coupled to glue

Light: up, down (u,d). Kinda: strange (s)

Heavy: charm (c), bottom, top.



# QCD & Confinement

Non-perturbative effects => confinement. *Only color singlets observable.*

Quark + anti-quark = meson.

Very light: pions ( $\bar{u}d$ ...), 140 MeV. kaons, ( $\bar{u}s$ ...), 490 MeV

= Spin waves of global (“chiral”) symmetry for light quarks, u, d, & s.

Three quarks = baryon.

Neutrons, protons ( $ddu$ ,  $uud$ ), 940 MeV; also strange ( $uds$ ...), 1-2 GeV

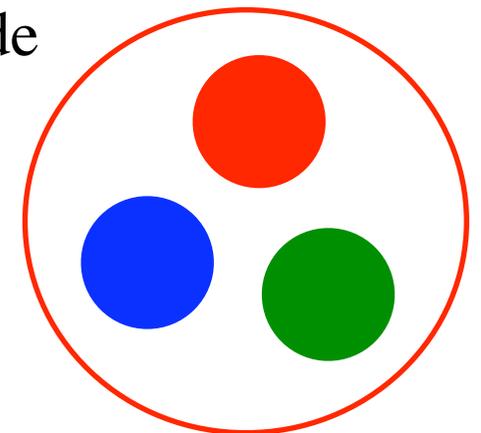
Charm mesons, J/Psi, ( $\bar{c}c$ ) 3.1 GeV. *Glueballs?*

One model: MIT bag. Confine quarks & gluons inside bag =>

Costs Bag constant  $B \sim (200 \text{ MeV})^4$ ;  $\sim$  perturbative inside

Dynamics of bag? *Surface has zero width*

Spectrum OK,  $\sim$  non-relativistic quark model



# Large N Expansion

Instead of SU(3) color, consider SU(N). **Claim: N=3 is “near” infinite N!**

Gluons: SU(N) matrix, so #  $\sim N^2-1$ . Quarks: SU(N) vector, so #  $\sim N$ .  
 $\Rightarrow$  large N exp. is dominated by glue.

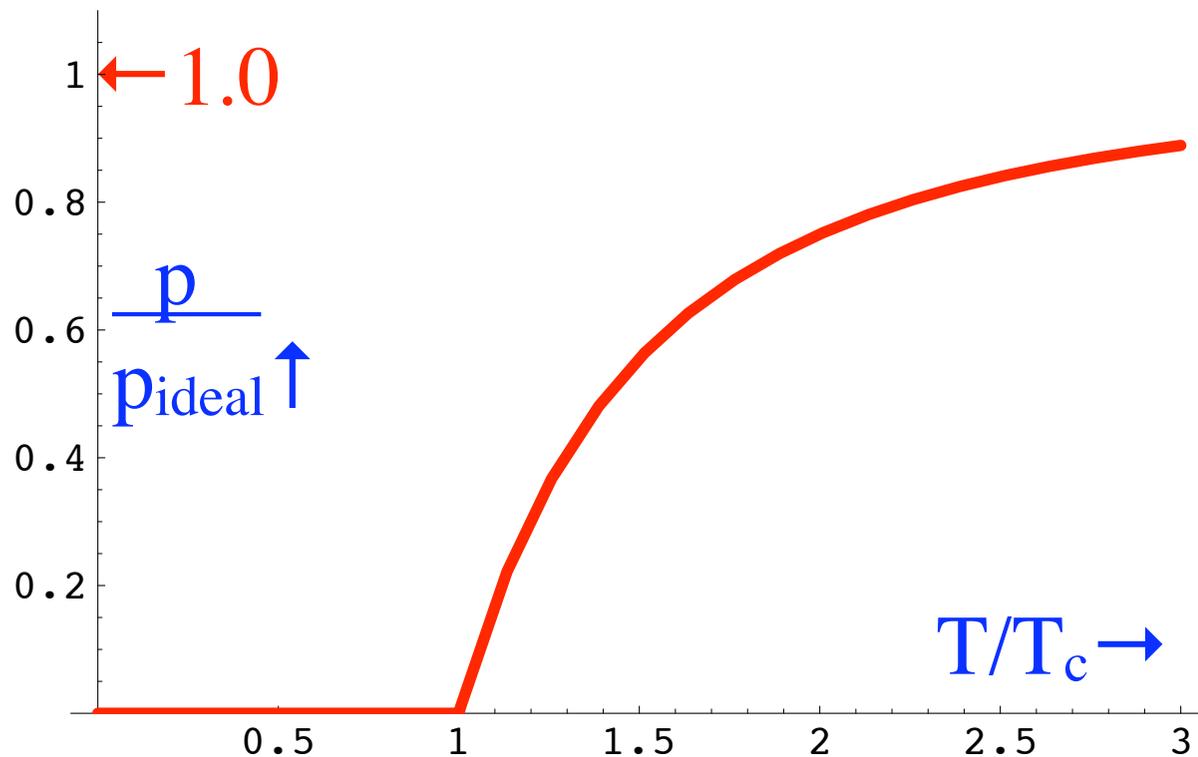
AF  $\Rightarrow$  *ideal* gas of  $\sim N^2$  gluons at  $T = \infty$ . Presumably for all  $T > T_c$

Confinement  $\Rightarrow$  only color singlets, so # hadrons  $\sim 1$  for  $T < T_c$

At *infinite* N:  
pressure/(ideal gas pressure)  
is an *order parameter*:

Ratio = 0 below  $T_c$  ( $\sim 1/N^2$ ),  
 $\neq 0$  above  $T_c$ ,  $\rightarrow 1$  as  $T \rightarrow \infty$ .

Possible form for this ratio  $\Rightarrow$



# Numerical “experiments”: the lattice

How to compute non-perturbatively? Put QCD on a lattice, with spacing “ $a$ ”.

In QCD, AF + renormalization group  $\Rightarrow$  *unique* answer as  $a \rightarrow 0$ , continuum limit

Approximate  $\infty$ -dim. integral by finite dim.: “Monte Carlo”

In practice: need to “improve” to control lattice effects  $\sim a^2$ , etc.

For pure gauge (*no* quarks) near continuum limit from mid-’90’s!

Compute equilibrium thermodynamics: pressure(temperature) =  $p(T)$

Lattices: 6-8 steps in time, 24-32 in space  $\Rightarrow 8 \cdot 32^3 \cdot 8 = 10^6$  dim. integral

With dynamical quarks, much harder: today, *not* near continuum for light quarks

Quarks non-local, difficult to treat (global!) “chiral” symmetry of light quarks

Today: only  $T \neq 0$ , *not* dense quarks at  $T=0$

With quarks: *orders* of magnitude more difficult

**BIG** calculations done on **BIG** machines...

To the right: “QCDOC”, QCD on a chip

RIKEN/BNL: 12,000 nodes, 20 teraflops...



# Lattice: SU(3) gluons deconfine

Confined phase at  $T=0$ : only “glueballs”. *Lightest* glueball 1.5 GeV

Deconfinement at  $T_c \sim 270\text{MeV} \pm 10\%$  *Small* relative to glueball masses.

Pressure *very* small below  $T_c$ : *very* much like large  $N$ !

MIT bag model does *not* work  $p_{\text{MIT}}(T) = \#T^4 - B$

“Fuzzy” bag does:

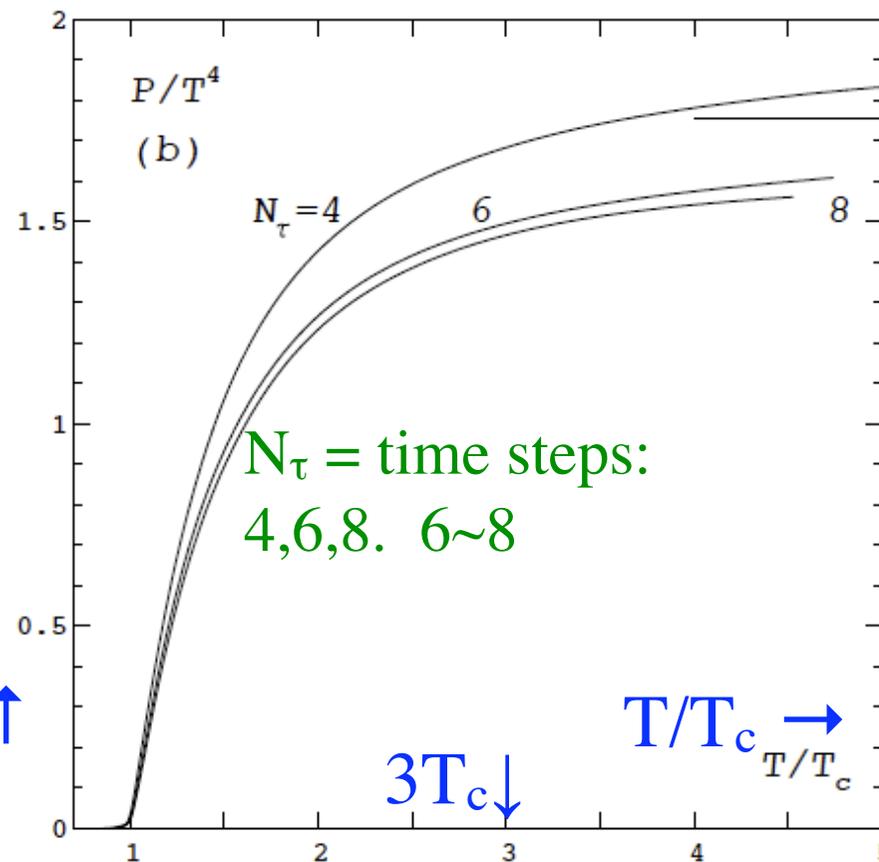
$$p_{\text{fuzzy}}(T) = \#(T^4 - T_c^2 T^2)$$

$\# = 90\%$  of ideal gas  $\Rightarrow T^4 \sim$  perturbative

Term  $\sim T^2$  non-pert. power law correction

Transition weakly first order

$\sim Z(3)$  Potts model



$p(T)/T^4 \uparrow$

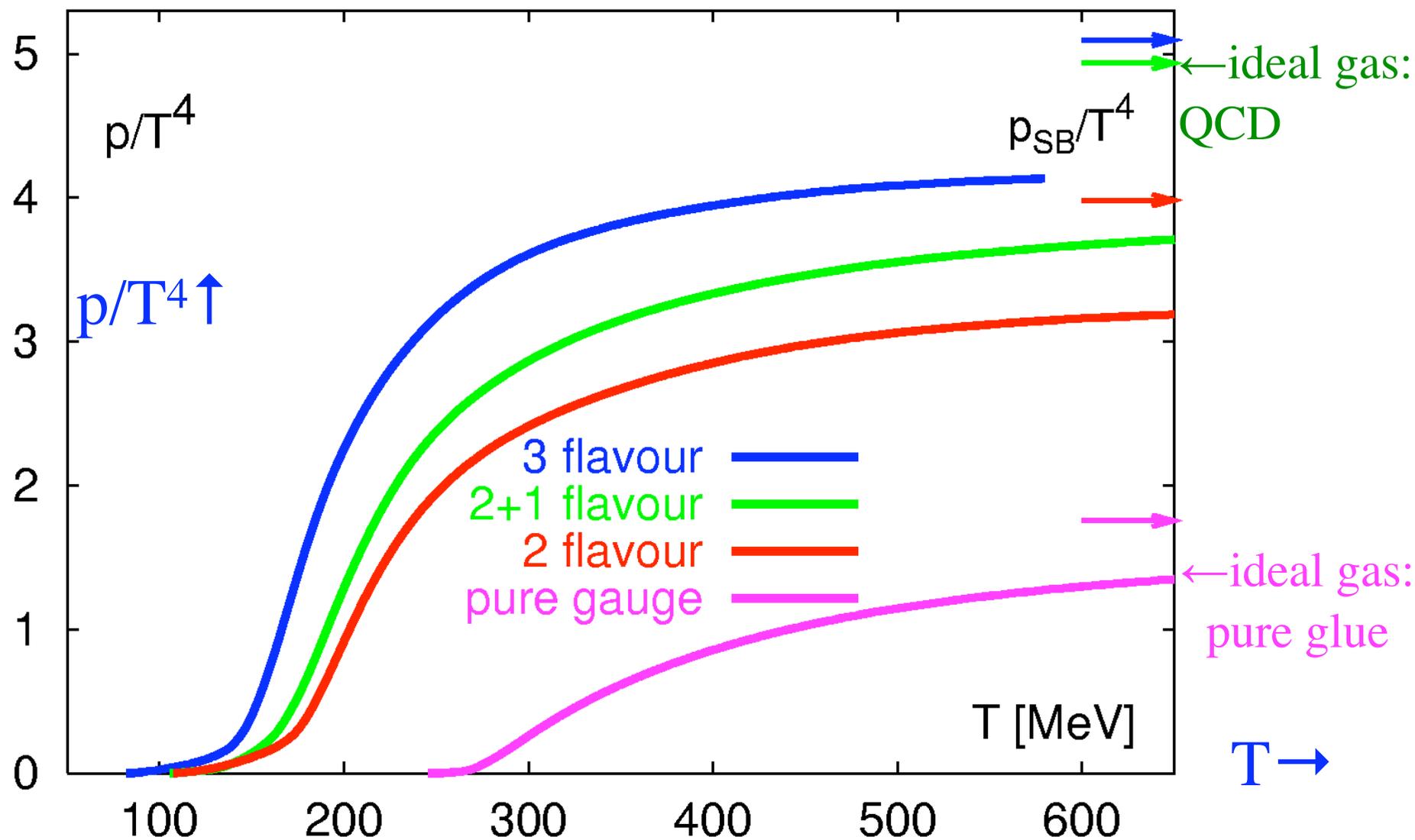
# Lattice: QCD, with quarks

Adding quarks: ideal gas value increases by  $\sim 3$

“ $T_c$ ” decreases. Today:  $\sim 190 \text{ MeV} \pm ?$  1st order transition washed out, *crossover*.

Pressure nonzero  $< T_c$ : from pions, kaons..., masses  $\sim T_c$

Pressure within 20% of ideal by  $\sim 3 \times T_c$ . Different behavior?



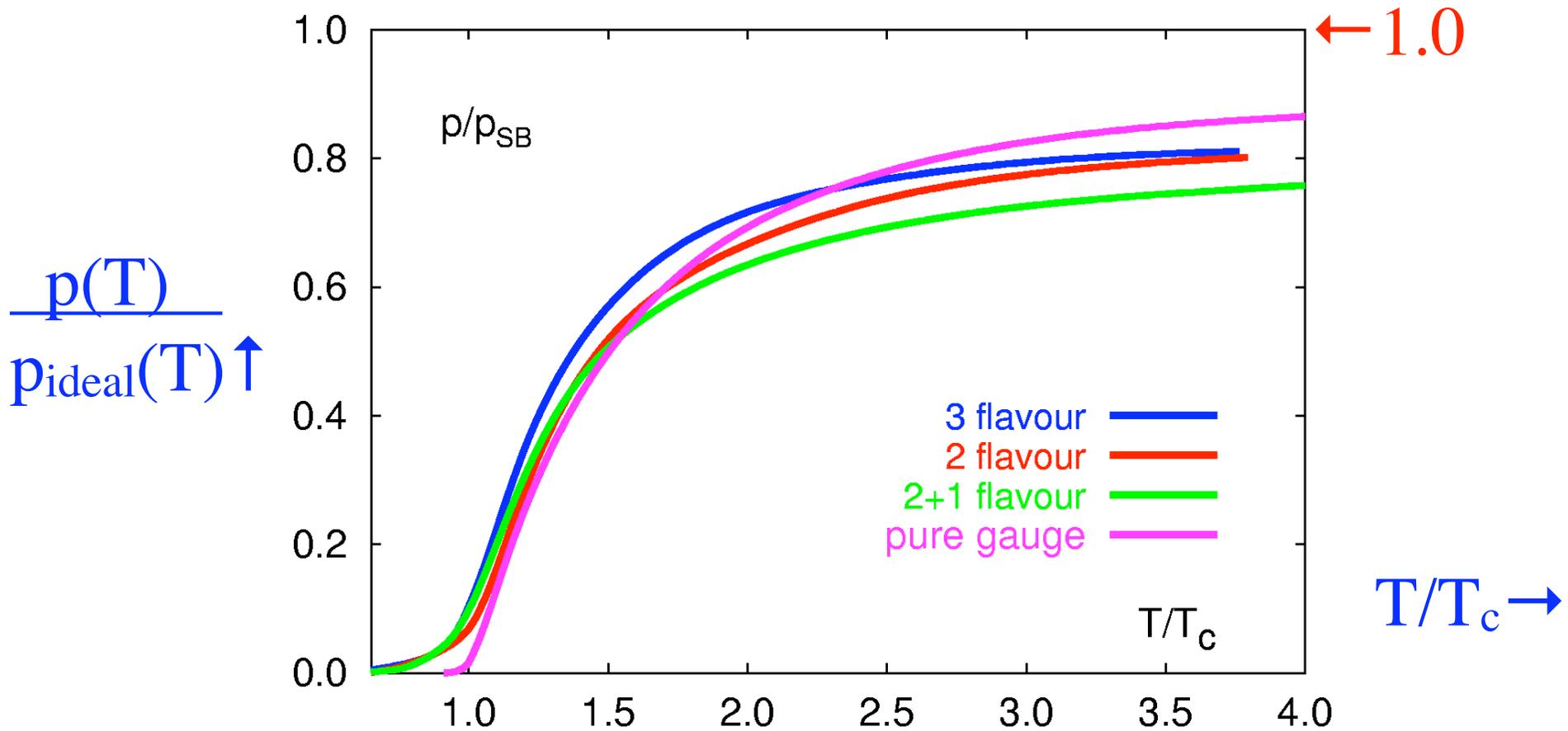
# Lattice thermo.: QCD & “flavor independence”

Bielefeld: properly scaled,  $\approx$  *universal* pressure  
Transition dominated by deconfinement

$$\frac{p}{p_{ideal}} \left( \frac{T}{T_c} \right) \approx \text{const.}$$

In *all* cases so far, chiral symmetry restored at *same*  $T_c$  as deconfinement: *why?*

Is transition from confined phase, directly to nearly ideal QGP?

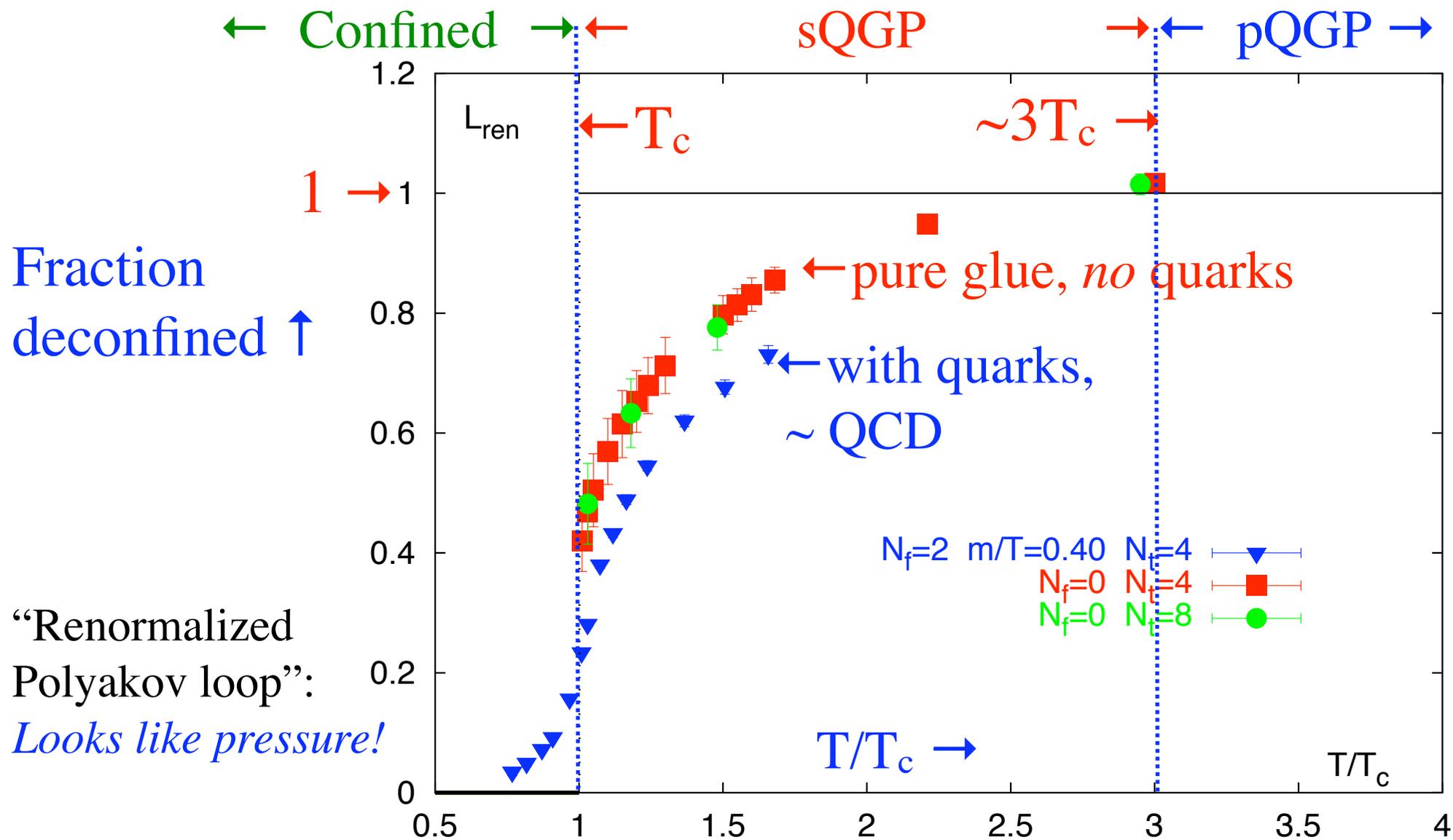


# Partially deconfined, $T_c$ to $\sim 3 T_c$ : “sQGP”

On lattice, can measure *fraction of deconfinement*

$T > 3 T_c$  : fraction  $\sim 1$ ,  $\Rightarrow$   $\sim$  perturbative QGP

$T_c \rightarrow 3 T_c$  : fraction  $< 1$ , partially deconfined. “sQGP” (s?)



# The “sQGP”: *not* strong coupling

“Helsinki Dogma”: resum pert. thy. to 4 loop order. *Fails below  $\sim 3 T_c$*

Sure! Coupling big below  $3 T_c$  !

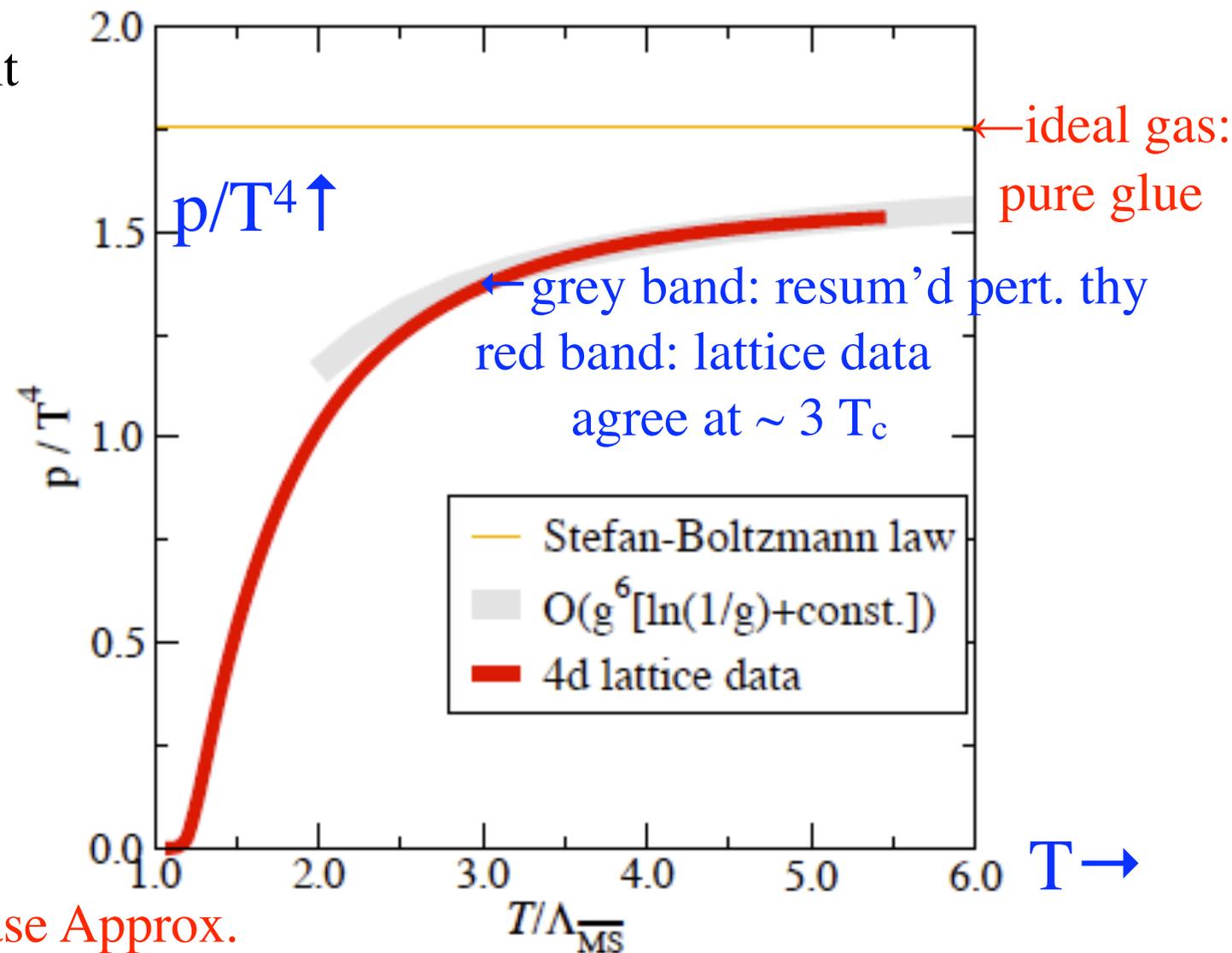
No! While deconfinement only partial at  $T_c$ ,  
*coupling moderate*

$\alpha_s(9T_c) \sim \alpha_s(1.6 \text{ GeV})!$

So what *is* “sQGP”,  
 $T_c$  to  $\sim 3 T_c$ ?

Semiclassical, dominated  
by solitons.

Eventually - NaRPA:  
Non-abelian Random Phase Approx.



# Hunting for the “Unicorn” in Heavy Ion Collisions



“Unicorn” & the QGP: D. Scott, R. Stock, M. Gyulassy, M. Veltman...

Experimentalists = hunters, so...all theorists are dogs...

# AA collisions at high energies

Collide: pp, protons on protons. Benchmark for “ordinary” QCD.

AA, nuclei on nuclei. Atomic # “A”: 60 => 200, Cu -> Au, Pb. “Hot” nuclei.

pA, proton (or deuteron) on nucleus. Another check: QCD in “cold” nuclei

**Why AA?** Baryons are like hard spheres, so for A: 60 - 200,  
transverse size  $A^{2/3}$  : 15 - 35  $\times$  proton. **Big nuclei are big!**

Total energy in the center of mass,  $E_{\text{cm}} = \sqrt{s}$  (GeV); per nucleon,  $\sqrt{s}/A = \sqrt{s_{\text{NN}}}$ .

SPS @ CERN	5 => 17 GeV
RHIC @ BNL	20 => 200 GeV
LHC @ CERN	5500 GeV
SIS @ GSI	2 => 6 GeV

SPS, SIS Fixed Target

RHIC, LHC Colliders

LHC > '08

SIS @ GSI, Darmstadt > '12



# Geometry of AA collisions at high energy

At *high energies*: nuclei Lorentz contracted along beam (15 fermi  $\Rightarrow$  0.3 fermi)

AF  $\Rightarrow$  nuclei don't stop, pass through each other.

Collider: lab = center of mass frame

Momenta of produced particles: along beam,  $p_z$ ; transverse to beam,  $p_t$

Baryons in original nuclei go down beam pipe, at large  $\pm p_z$

For pp collisions: particles  $\sim$  constant for some range in  $p_z$  about  $p_z = 0$  :

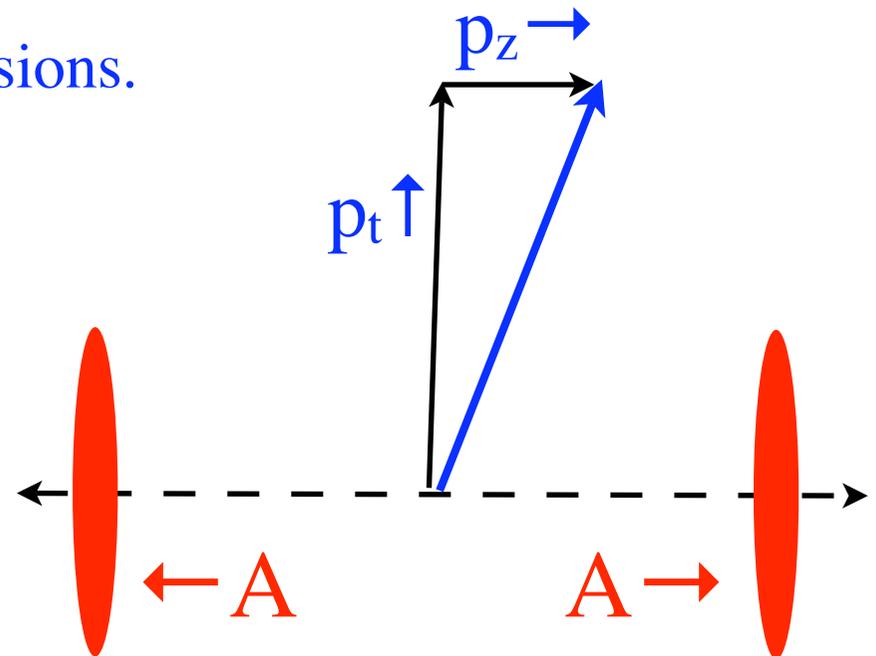
“rapidity plateau” (rapidity  $\sim p_z$ ; boost invariance)

Bjorken: look at rapidity plateau in AA collisions.

Rapidity plateau  $\sim$  free of incident baryons.

$\Rightarrow$  most likely to be at nonzero temperature,  
zero (quark) density.

Collider: central plateau  $90^\circ$  to beam



# Typical Au-Au collision @ RHIC

Experiments @ RHIC:

“Big”: ~ 400 people. STAR & PHENIX

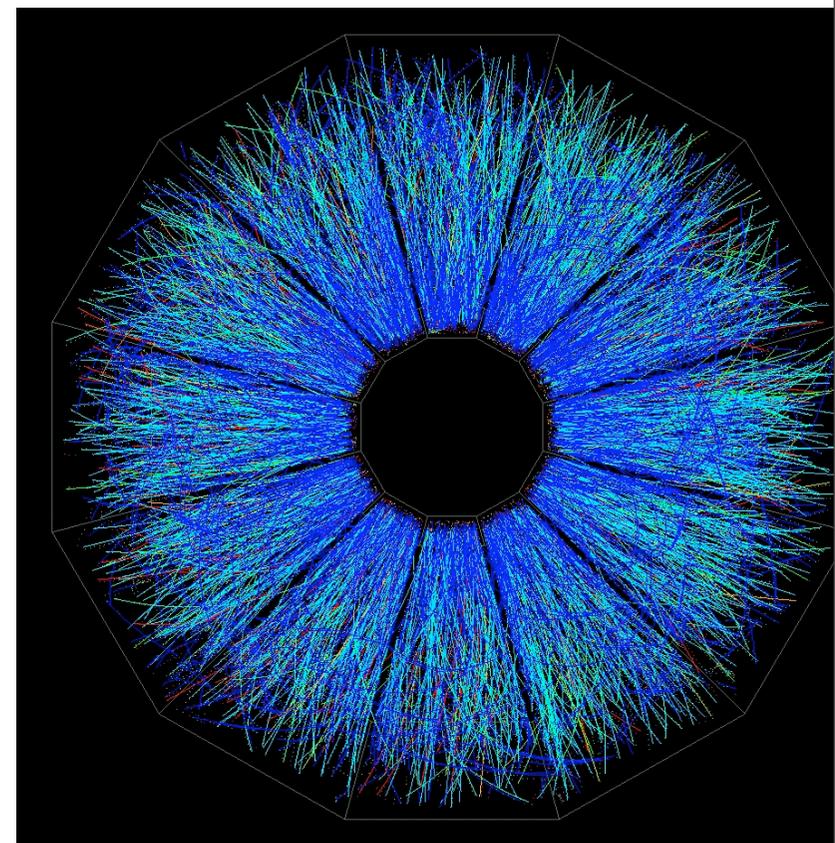
“Small”: ~ 50 people. PHOBOS & BRAHMS

*No surprises in total multiplicity; ~ 1.3 A × pp:*

total # particles ~ ~ log(total energy)

~total # experimentalists

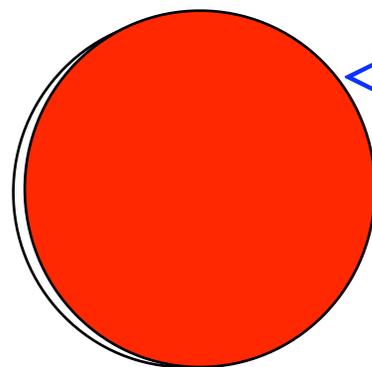
# theorists ~ log(log(total energy)).



Need hunters more than dogs...

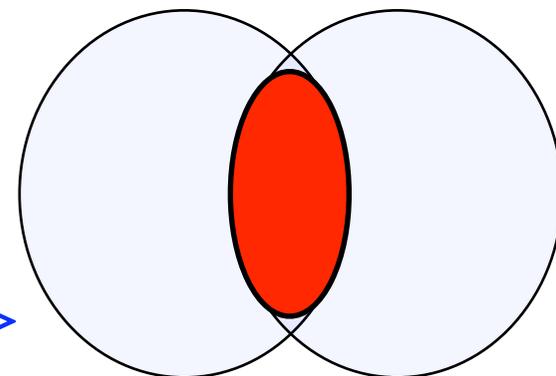
Total # particles/unit rapidity ~900↑

Also: can exp.'y measure how much nuclei overlap in transverse plane



<= central

peripheral =>



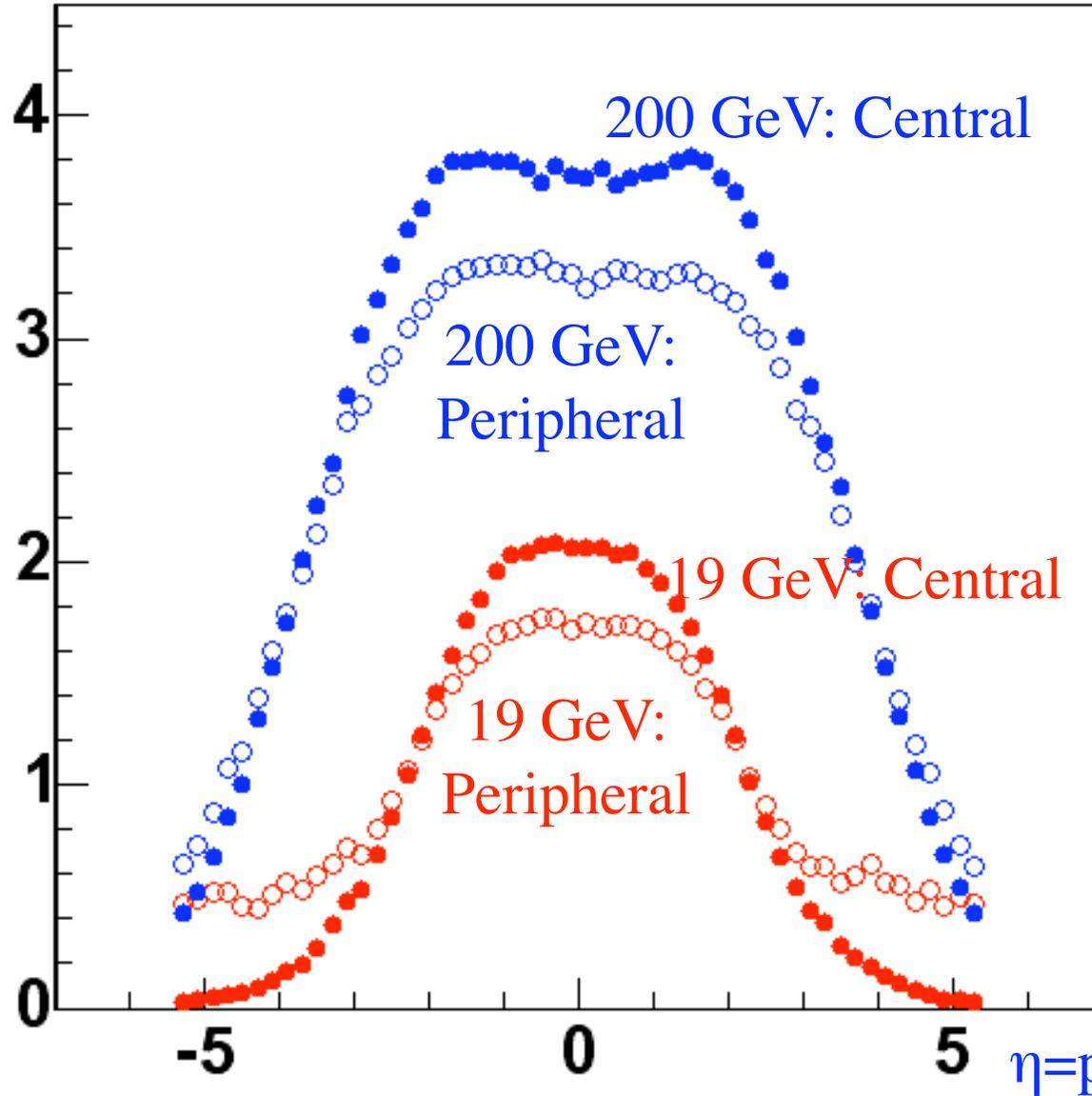
# Overall multiplicity: slow growth

$dN/d\eta/\uparrow$

$N = \#$  particles

$\eta =$  "pseudo"-  
rapidity  
(no particle ID)

/ by #  
"participants"



200 GeV, RHIC  
900 particles  
/unit  $\eta$

19 GeV, SPS  
600 particles  
/unit  $\eta$

No big increases in multiplicity, as predicted by cascade models.

Rapidity plateau  $\pm .5$  (out of  $\pm 5.0$ ) for  $dN/dy$  ( $y =$  true rapidity)

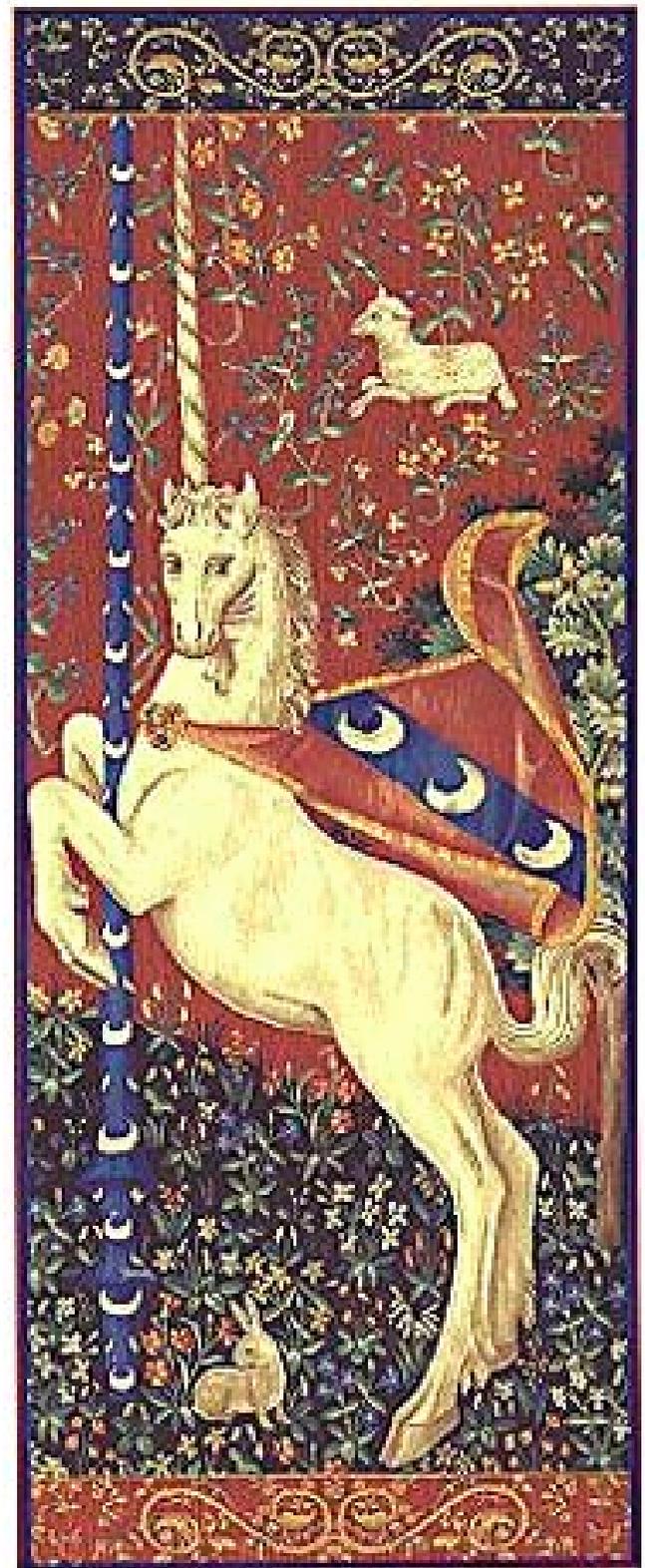
# The Tail Wags the (Dog) Unicorn

Body of the “Unicorn”:

Majority of particles at small momenta  
 $p_t < 2 \text{ GeV}$ .

Tail of the “Unicorn”:

Look at particles at HIGH momentum,  
 $p_t > 2 \text{ GeV}$ , to probe the body.

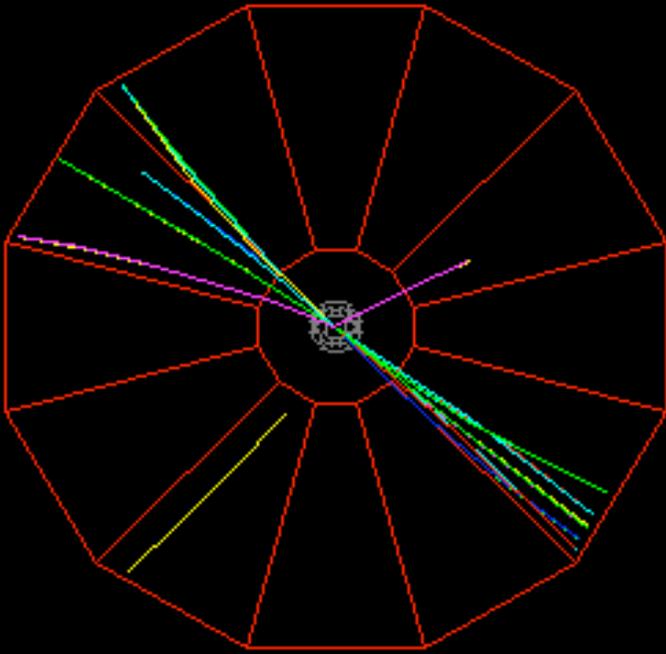


# Jets: “seeing” quarks and gluons in QCD

At high transverse momentum ( $p_t$ ), instead of hadrons, have  $\sim$  quarks & gluons: *jets*.

$\leq 2$  jets from pp collision at RHIC.

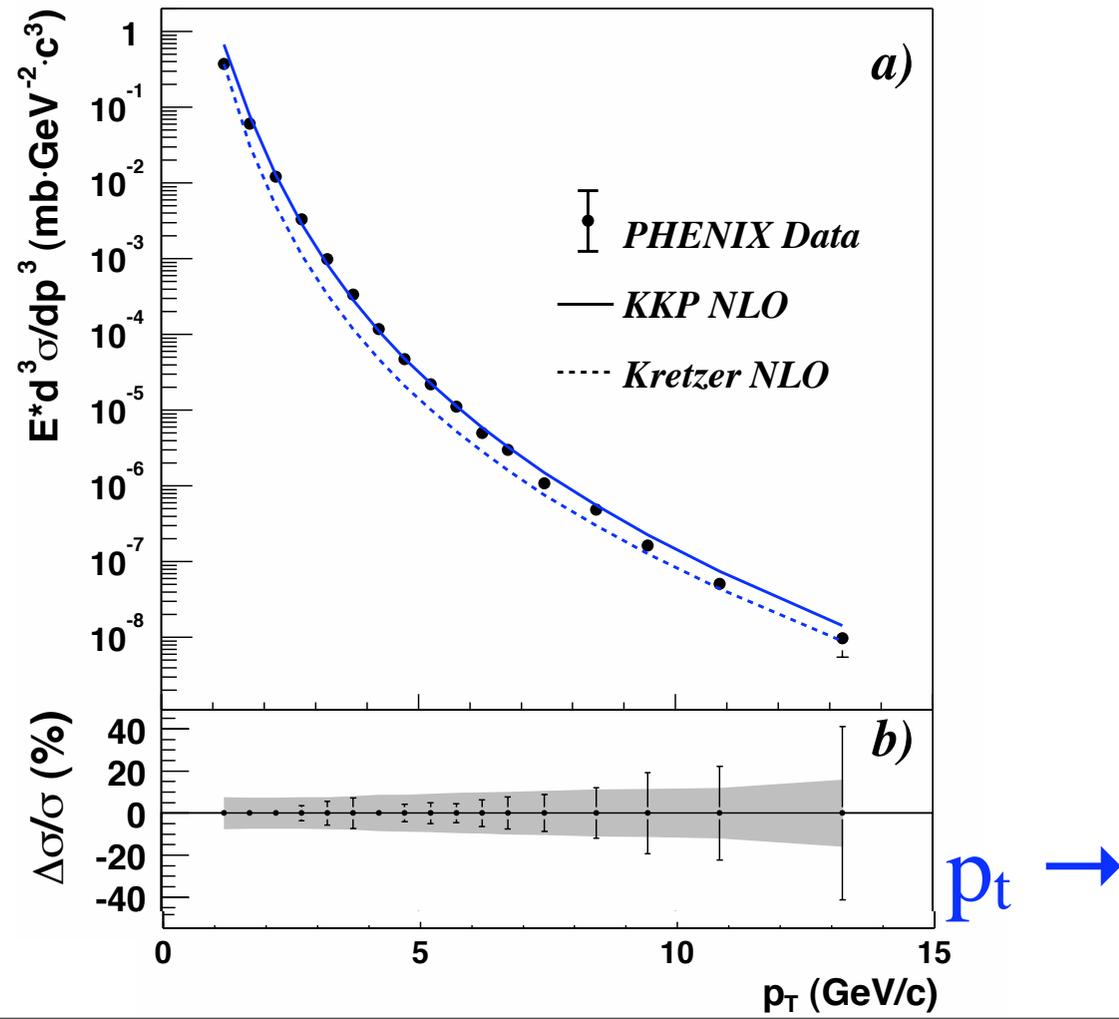
By momentum conservation, for each jet, there is a backward jet.



In pp coll.'s, jets can be computed for large  $p_t$ , down to  $p_t \sim 1$  GeV.

*Jets rare*: particles at  $p_t \sim 2$  GeV  $\sim .1$  % of total!

Look at jets in AA



# “ $R_{AA}$ ”: *robust* signal of new physics

$R_{AA}$  = for a given  $p_t$ , # particles central AA / ( $A^{4/3}$  # particles pp)

For  $\pi^0$ 's,  $p_t : 2 \rightarrow 20$  GeV,  $R_{AA} \sim 0.2$ . As if jets emitted *only* from surface!

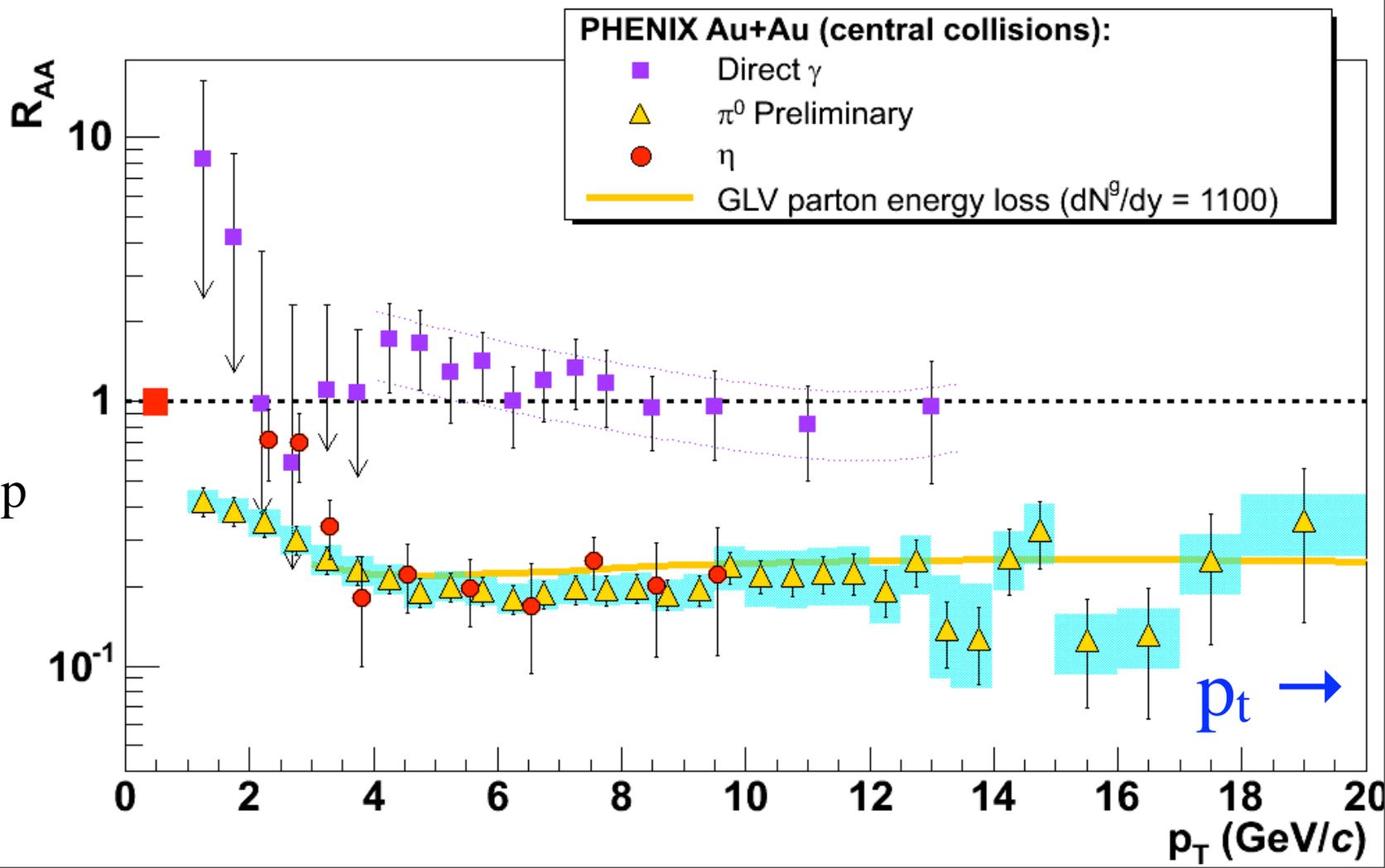
Due to “energy loss” in thermal medium?

$A^{4/3}$ : *experimentally*: for  $\gamma$ 's,  $R_{AA} \sim 1.0$   $\pi^0$ 's “eaten”,  $\gamma$ 's not

$R_{AA}$ :  $\uparrow$

# particles  
central AA/  
# particles pp

$A=200 \Rightarrow$

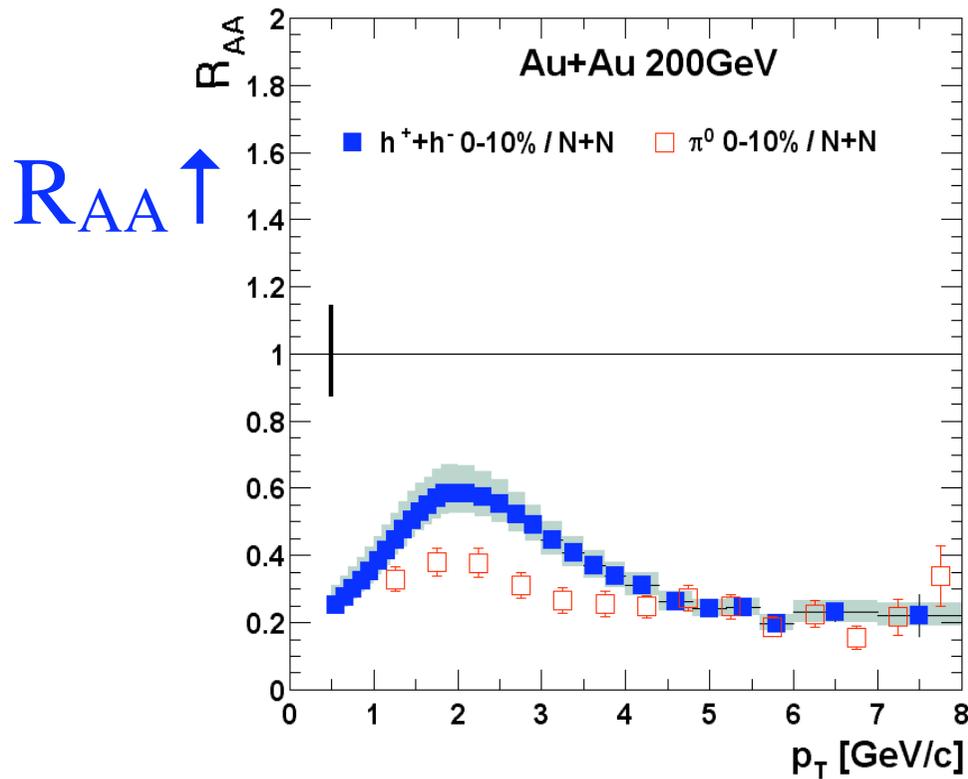


# $R_{AA}$ final state effect: *not* in $R_{dA}$

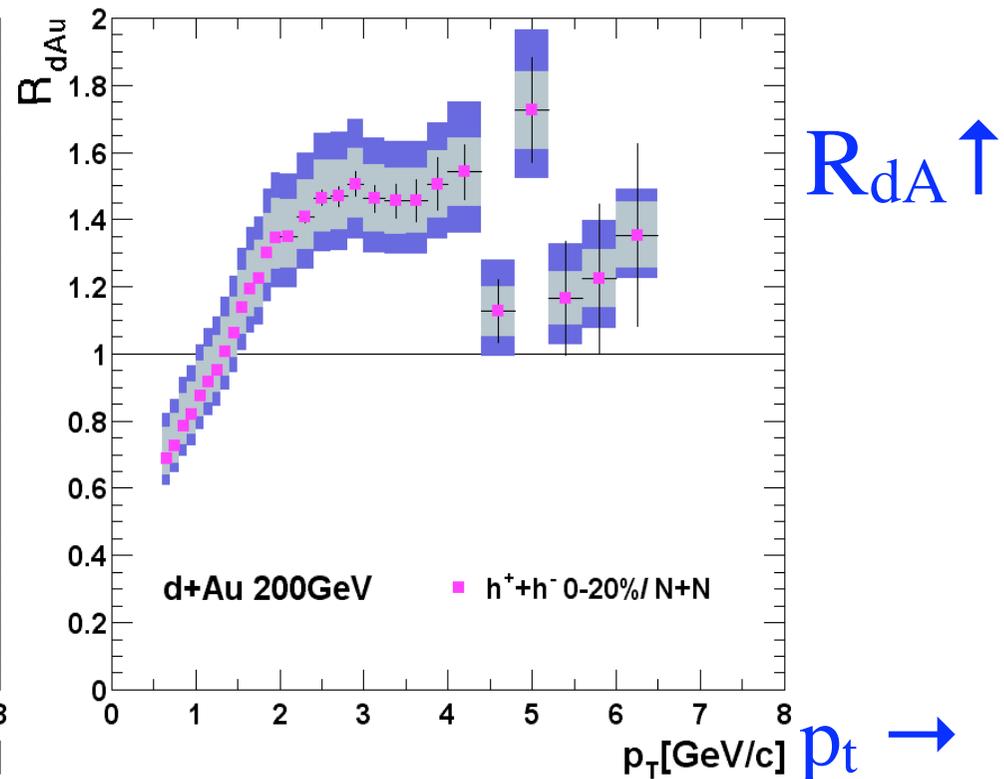
$R_{dA}$ : like  $R_{AA}$ , but for dA coll.'s/pp coll.'s. At zero rapidity:

dA: *enhancement*, from initial state “Cronin” effect ( $\Rightarrow 1$  @  $p_t > 8$  GeV)

AA: *suppression*  $\Rightarrow$  final state effect



Suppression in AA  $\uparrow$   
 $R_{AA} \sim 0.4$  @ 3 GeV



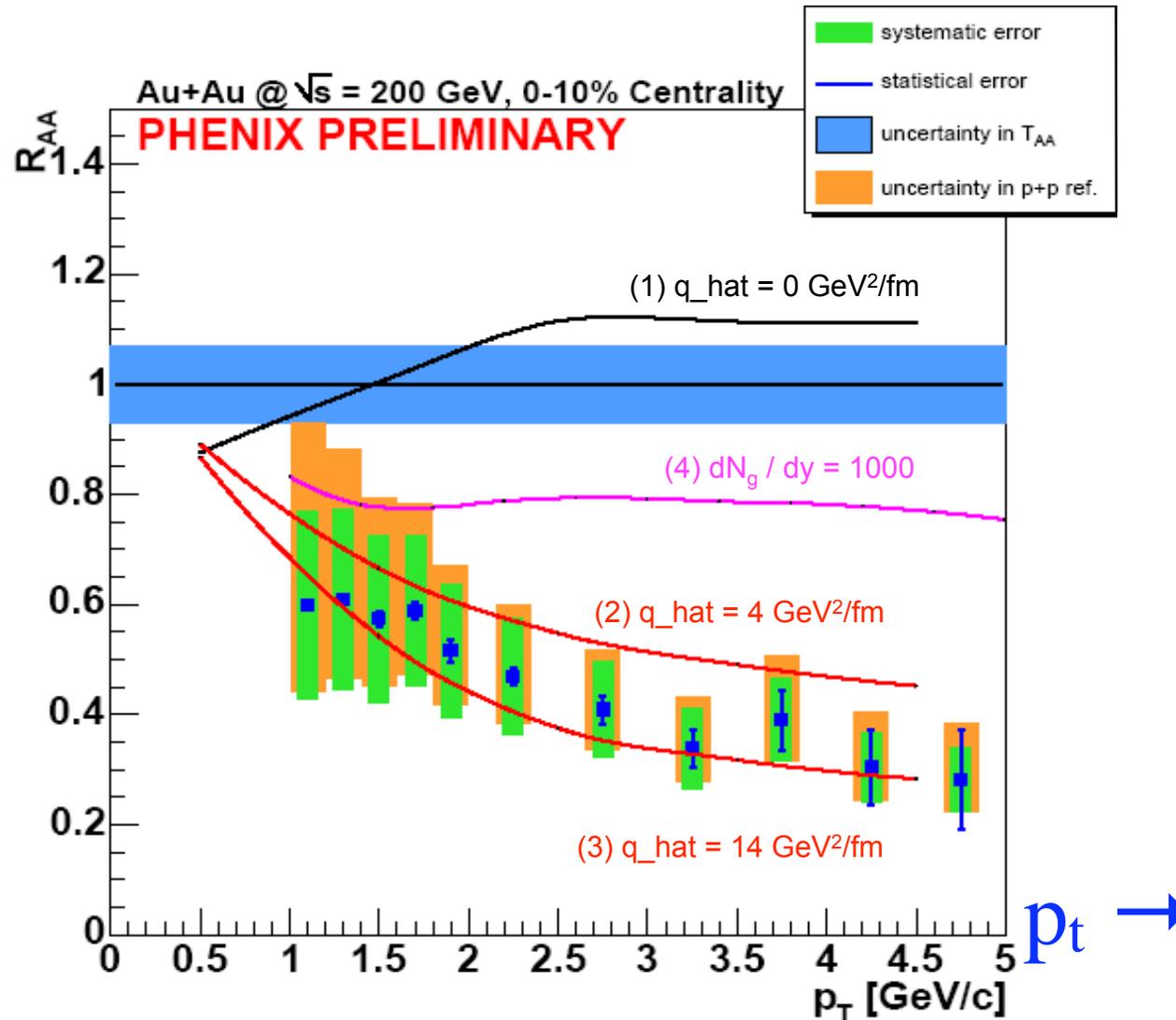
Enhancement in dA  $\uparrow$   
 $R_{dA} \sim 1.4$  @ 3 GeV

# $R_{AA}$ for heavy quarks: also suppressed!

PHENIX:  $R_{AA}$  for charm quarks  $\sim$  light quarks!

Theoretically? Heavy quarks should feel medium less by  $T/m \sim 1/10$  for charm.  
No sign of that!

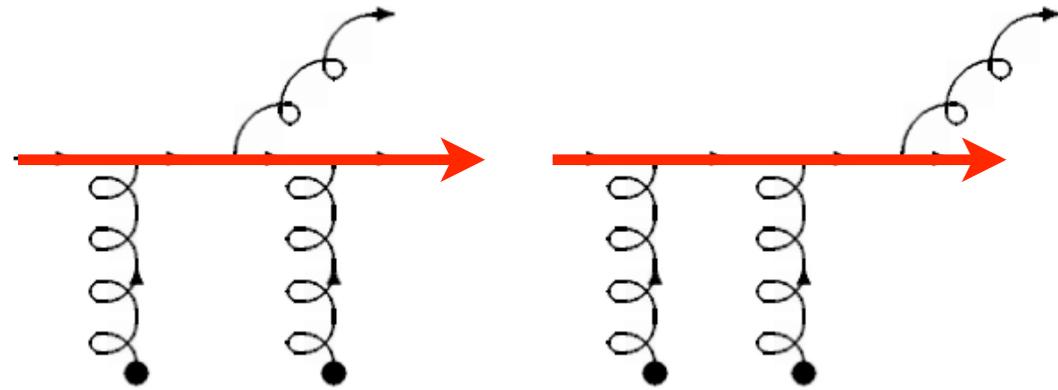
$R_{AA} \uparrow$



# Theory of $R_{AA}$ : energy loss

Fast quark (or gluon) emits radiation, scatters off of thermal bath.

Like QED, but qualitatively different in QCD: emitted gluon emits more gluons, etc.



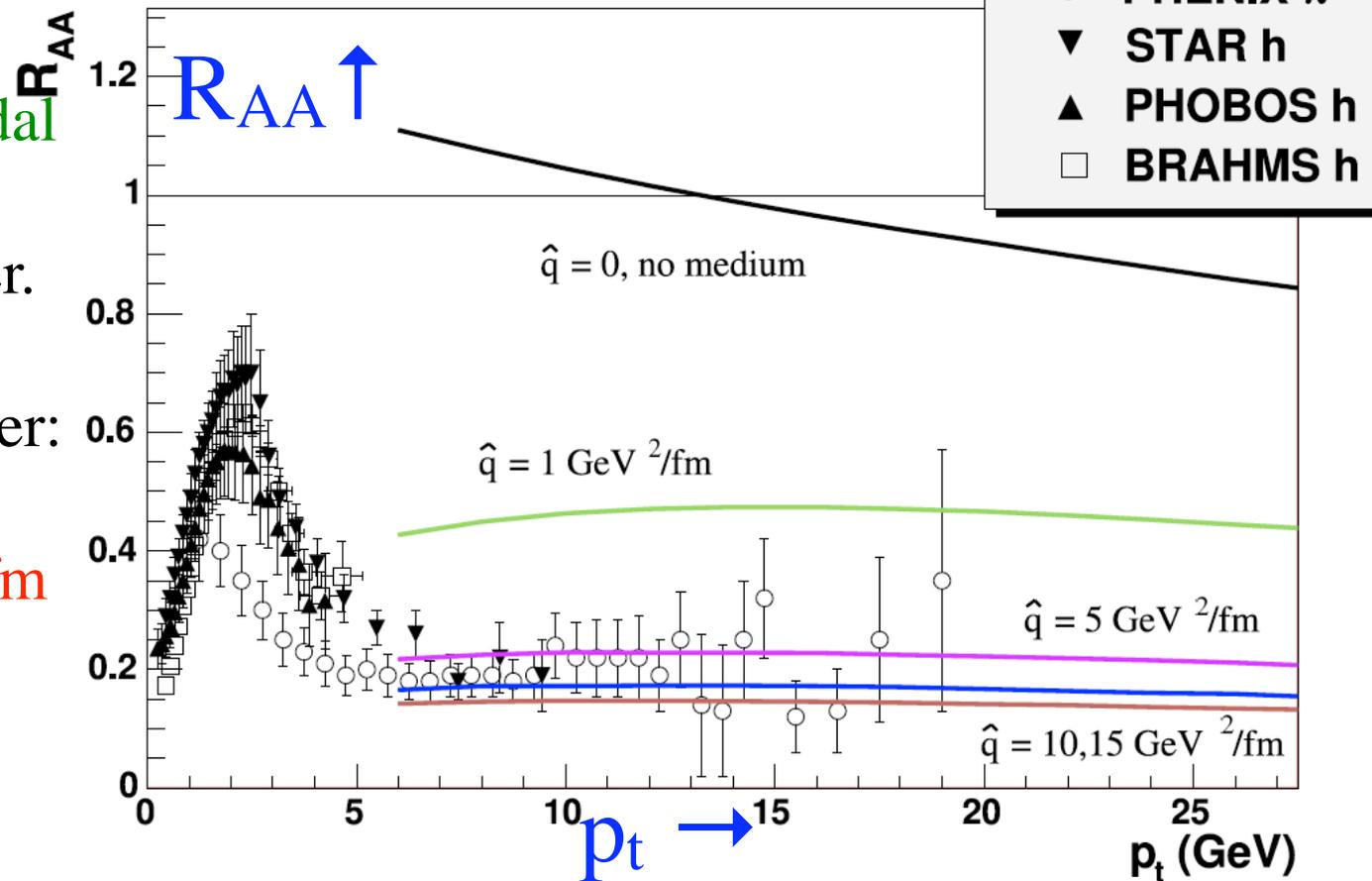
Involves LPM:  
Landau-Pomeranchuk-Migdal

Parametrized by one number.

Theorists disagree on number:

“weak” coupling  $\sim 2 \text{ GeV}^2/\text{fm}$

or “strong”  $\sim 15 \text{ GeV}^2/\text{fm}$ ?



# Central AA collisions “eat” jets!

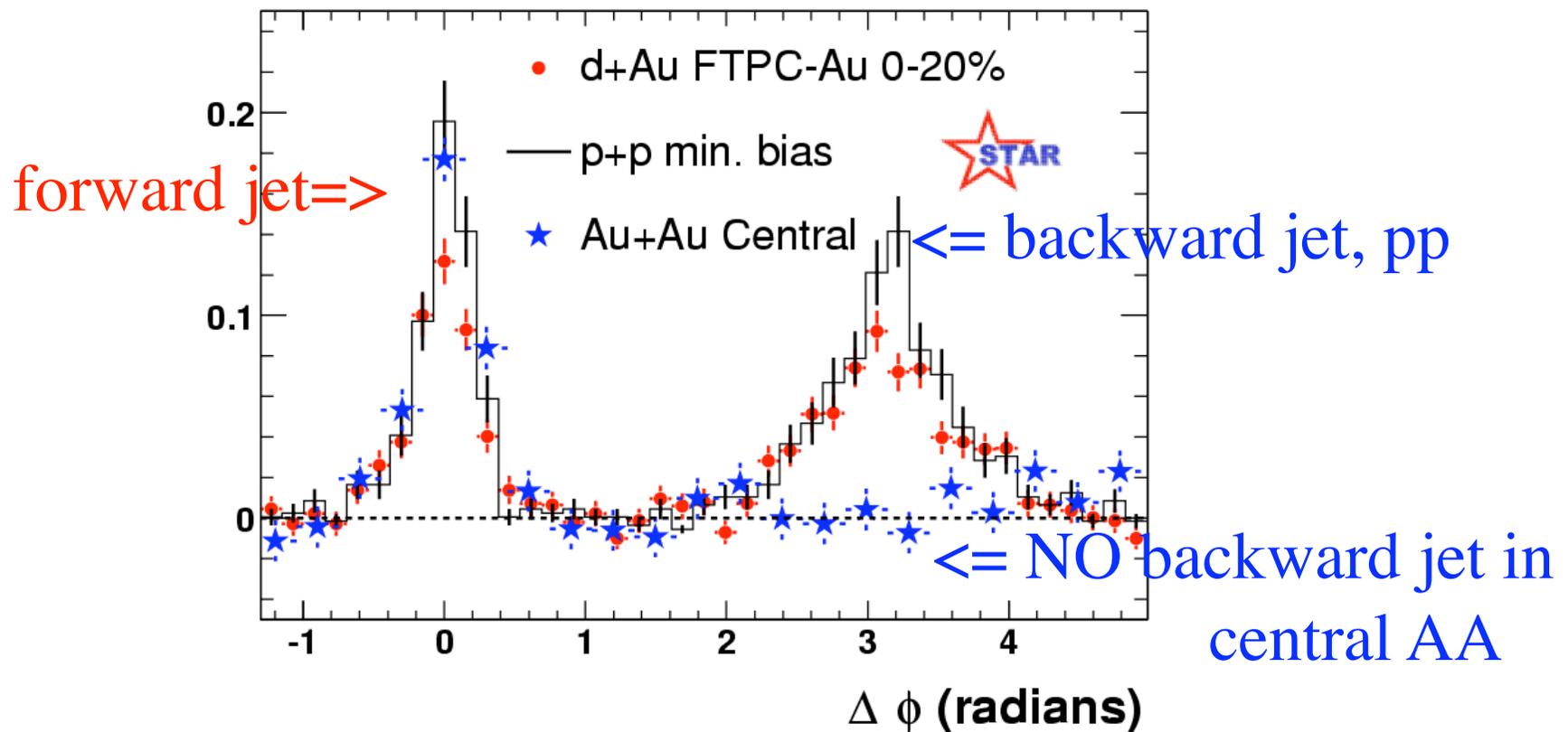
Unlike pp, in central AA, cannot trigger on individual jets. Can:

Trigger on “hard” particle,  $p_t: 4 \rightarrow 6$  GeV. Look for “away” side jet,  $p_t > 2$  GeV

In pp or dAu collisions, *clearly* see away side jet.

In central Au-Au, away side jet gone: “*stuff*” in central AA “*eats*” jets!

Adams *et al.*, Phys. Rev. Let. 91 (2003)



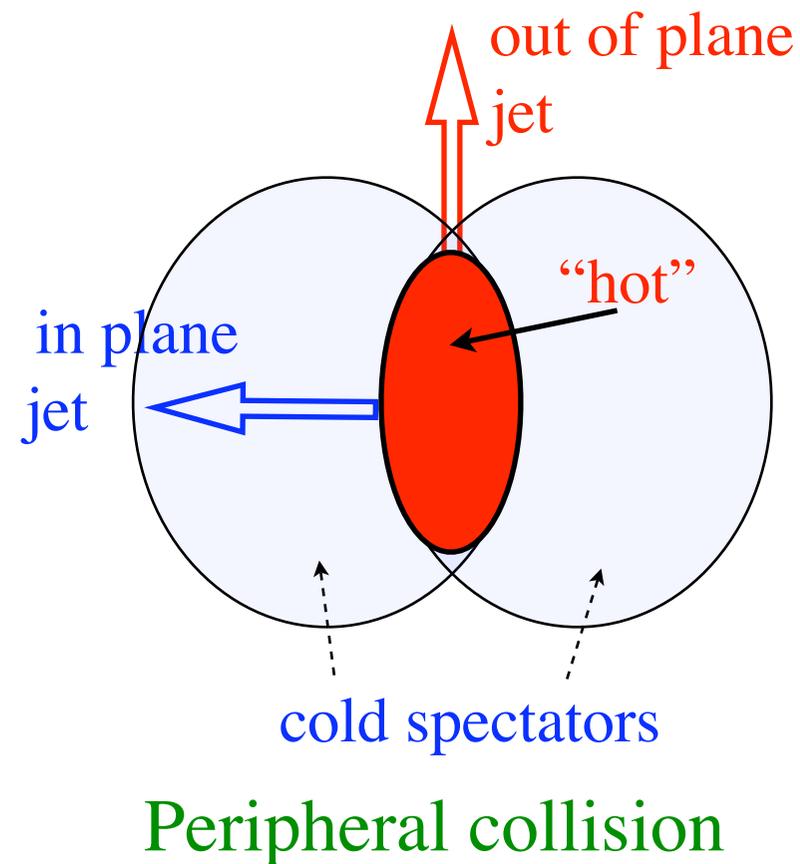
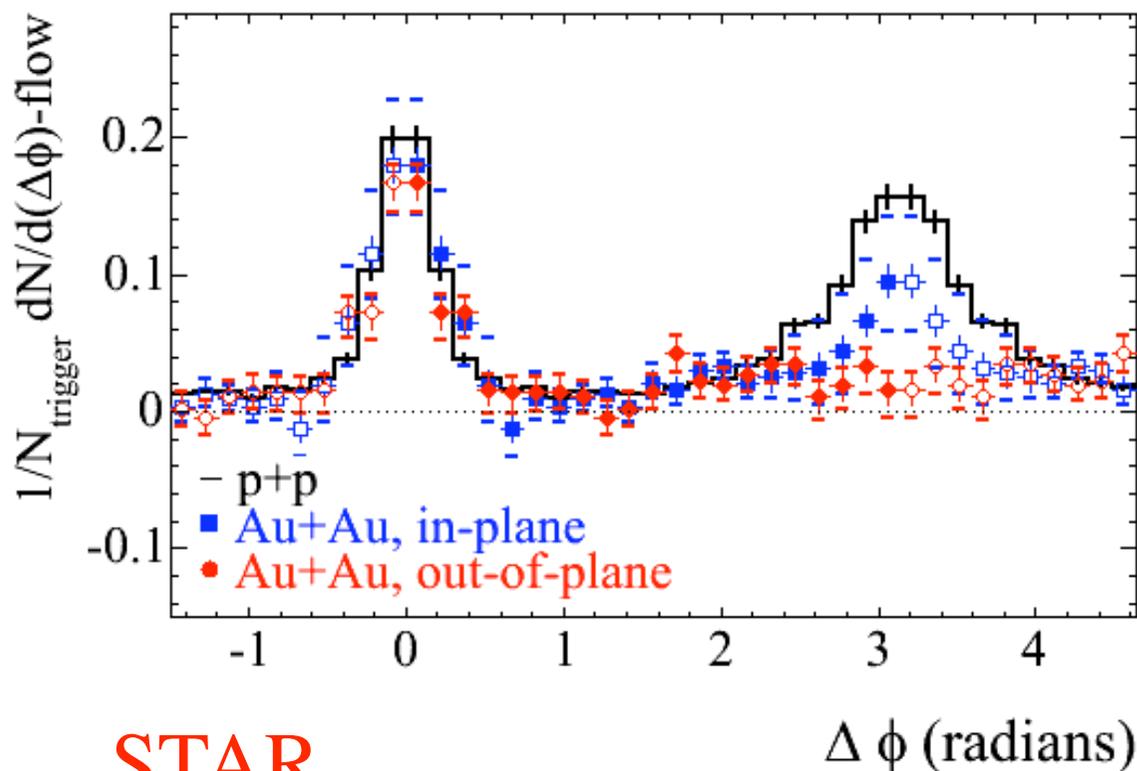
# Peripheral coll.'s: more jets eaten *out* of plane

Peripheral collisions: “hot stuff” forms “almond”:

*Out* of the reaction plane: more “hot stuff” than *in* plane

In the reaction plane, less hot stuff, but more cold nuclear matter

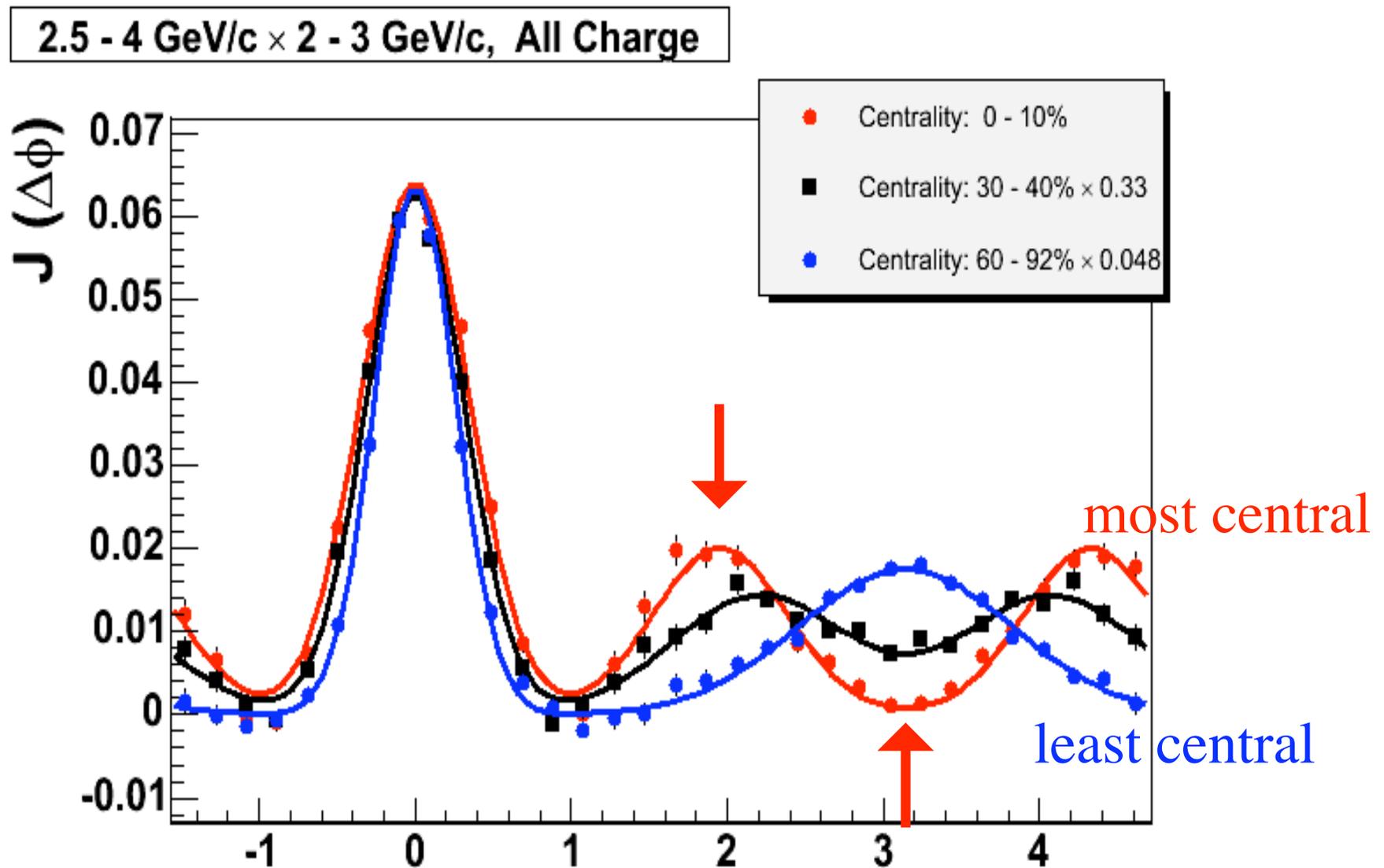
Exp.’y: backward jet more strongly suppressed *out* of plane than *in* plane =>  
*geometrical* test that central AA “eats” jets preferentially



# AA collisions modifies jet shapes

PHENIX: shape of away side jet appears to be modified by “stuff”:

Mach cone or Cerenkov radiation?



# The Body of the Unicorn: the sQGP

Particles peaked about zero transverse momentum

$T_c \sim 200$  MeV: expect thermal to  $p_t \sim 2$  GeV.

Thousands of particles, maybe hydrodynamics ok?

me →



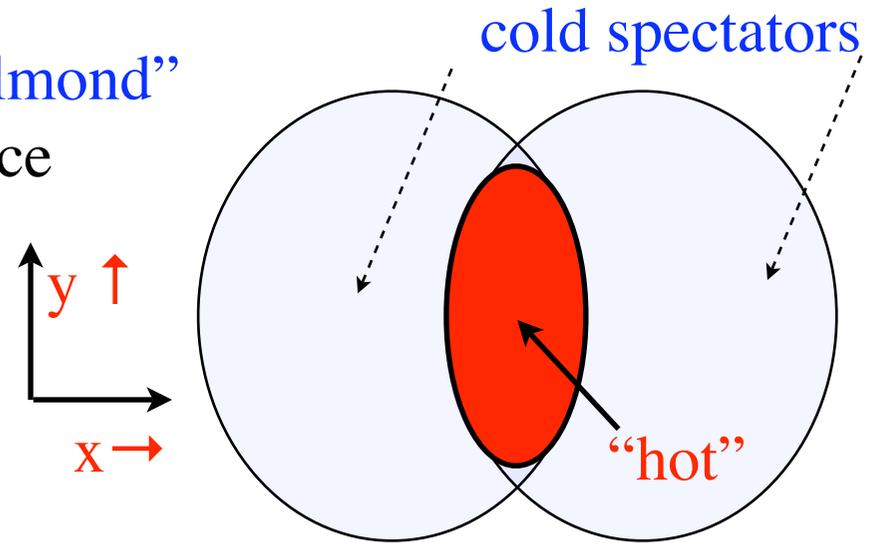
← Unicorn

# “Elliptic Flow”

For peripheral collisions, overlap region is “almond” in coordinate space, sphere in momentum space

So start with spatial anisotropy,

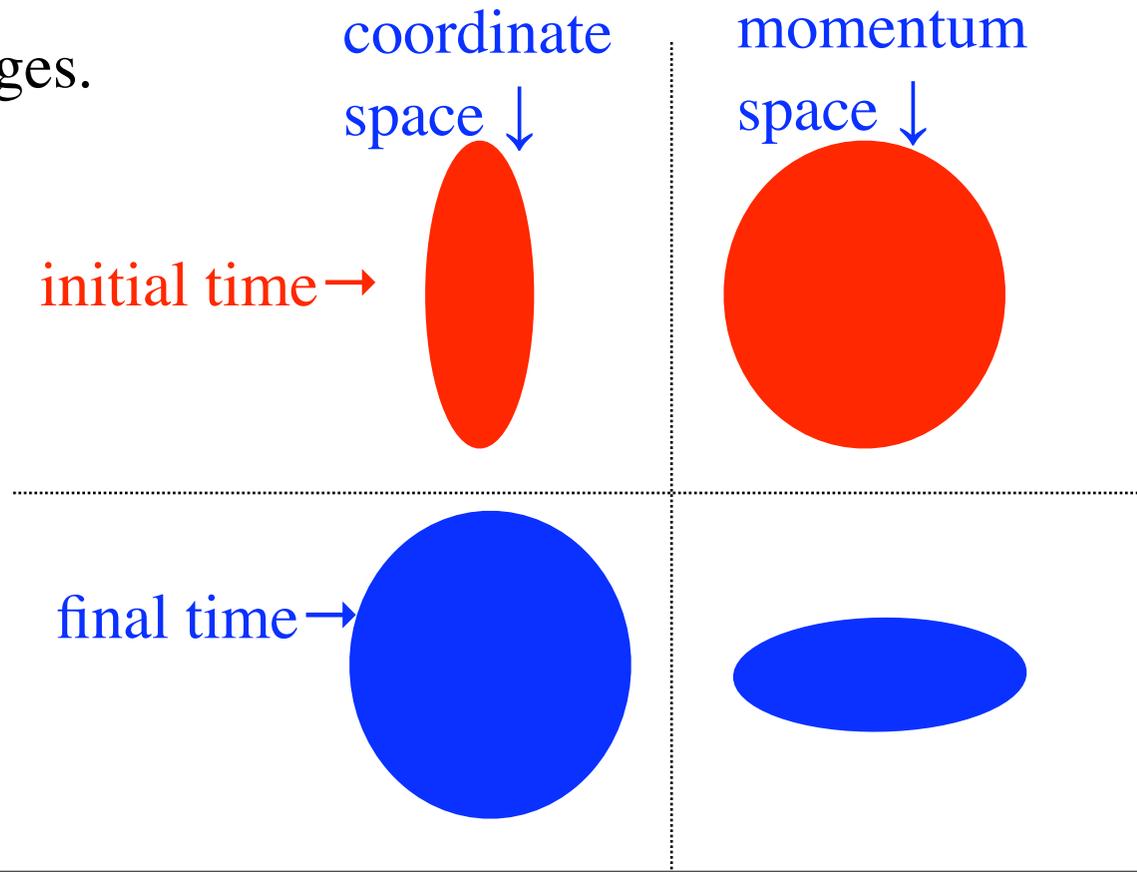
$$\epsilon = \frac{\langle y^2 - x^2 \rangle}{\langle x^2 + y^2 \rangle}$$



If particles free stream, nothing changes.

*If collective effects present*, end up with sphere in coordinate space, almond in momentum space: “elliptic flow”

$$v_2 = \frac{\langle p_y^2 - p_x^2 \rangle}{\langle p_x^2 + p_y^2 \rangle}$$



# “Most perfect fluid on earth”

Try *ideal* hydrodynamics: large # particles

*Bag* Equation of State.

Zero viscosity in QGP phase

Large viscosity in hadronic “afterburner”

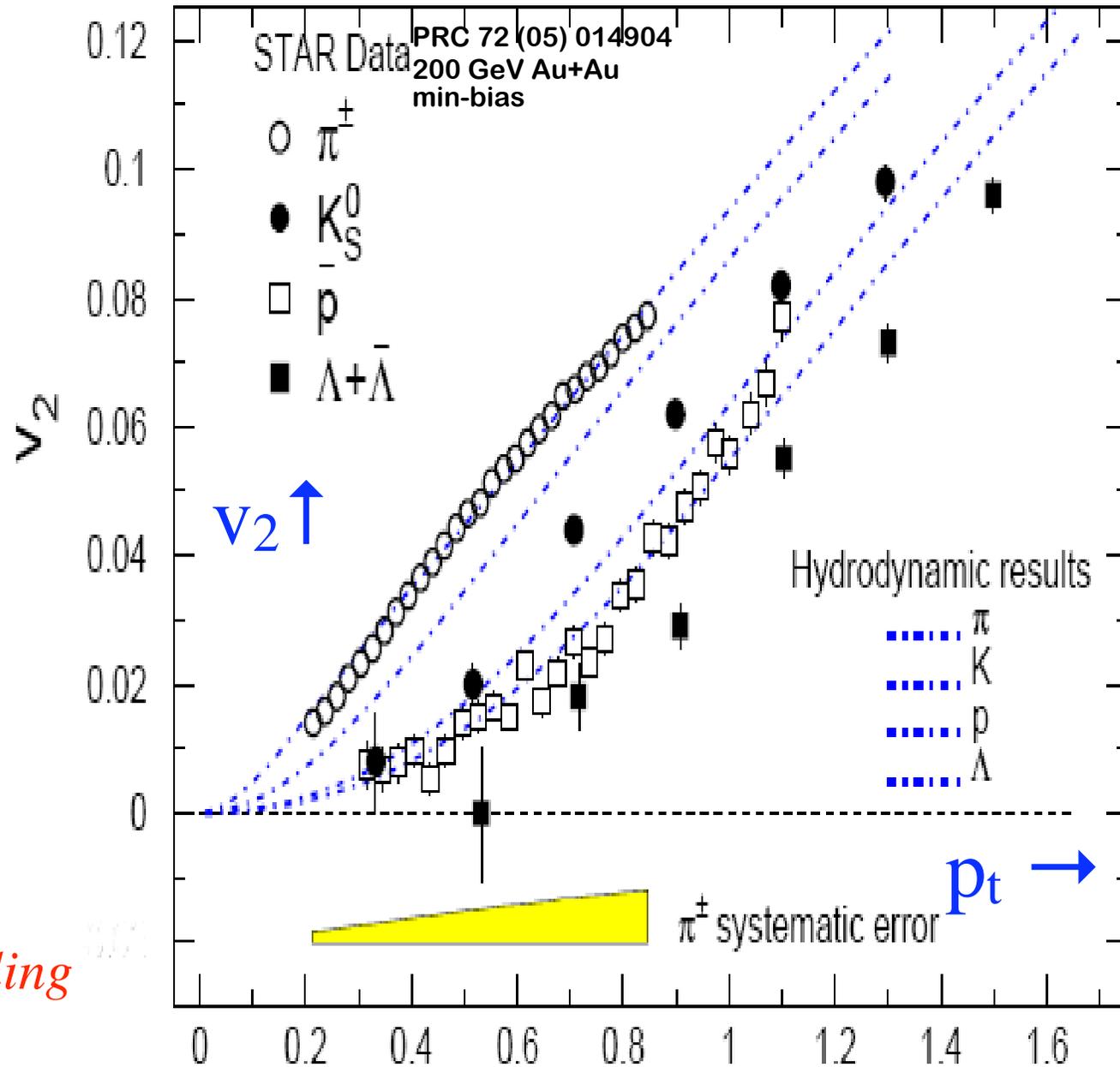
Short initial time (tune)

“Good” fit to  $\pi$ 's, K's, p's.

So far: *no* bound on viscosity in QGP.

Viscous, relativistic hydro. *very* hard.

Viscosity  $\sim 1/\text{cross section}$ : so *small viscosity*  $\Rightarrow$  *strong coupling* in sQGP?

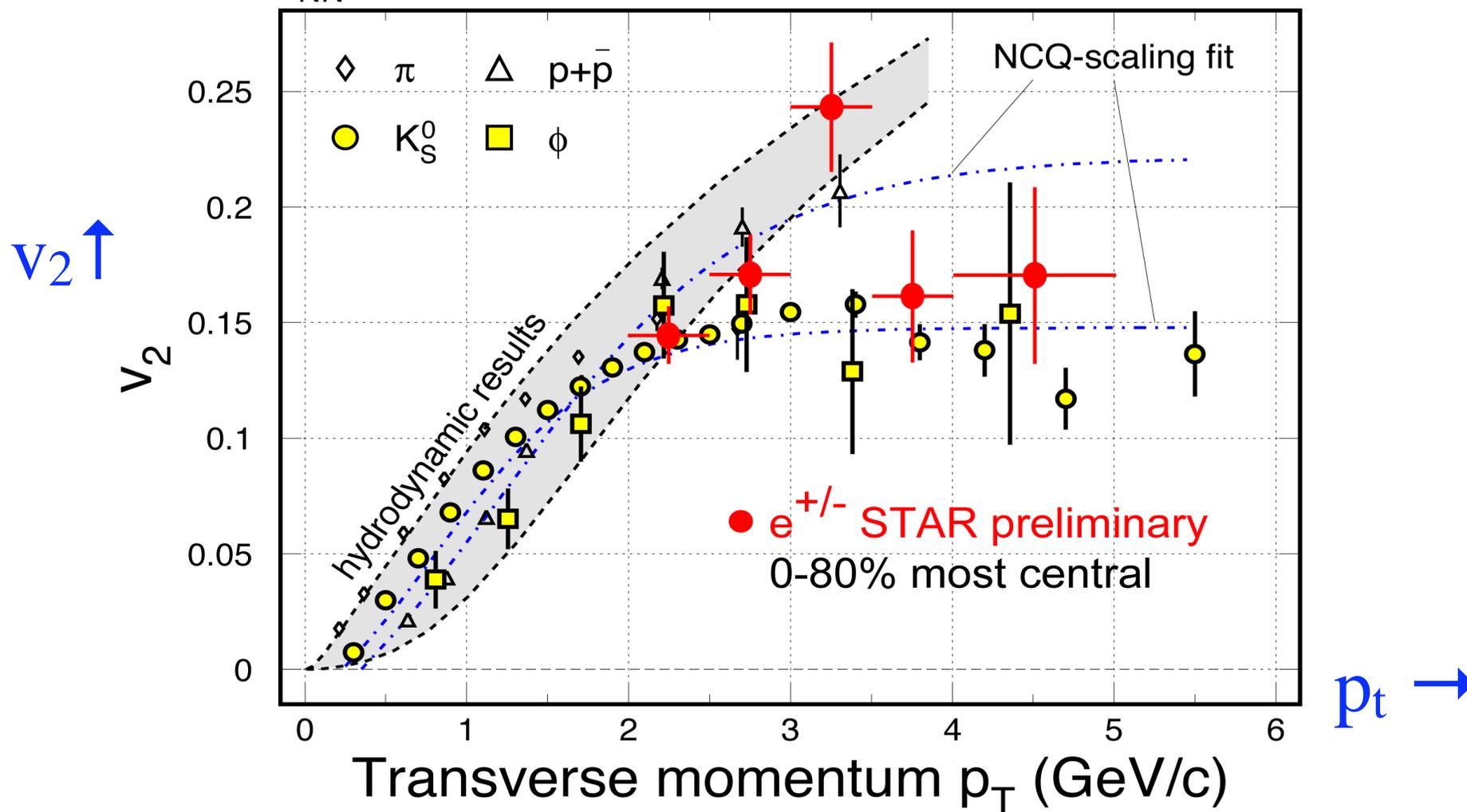


# Even charm quarks flow

Look at charm quarks through single electrons. See large elliptic flow.

No suppression due to large mass. Experimental evidence for “sQGP”.

$\sqrt{s_{NN}} = 200 \text{ GeV } ^{197}\text{Au} + ^{197}\text{Au} \text{ Collisions at RHIC (IV)}$

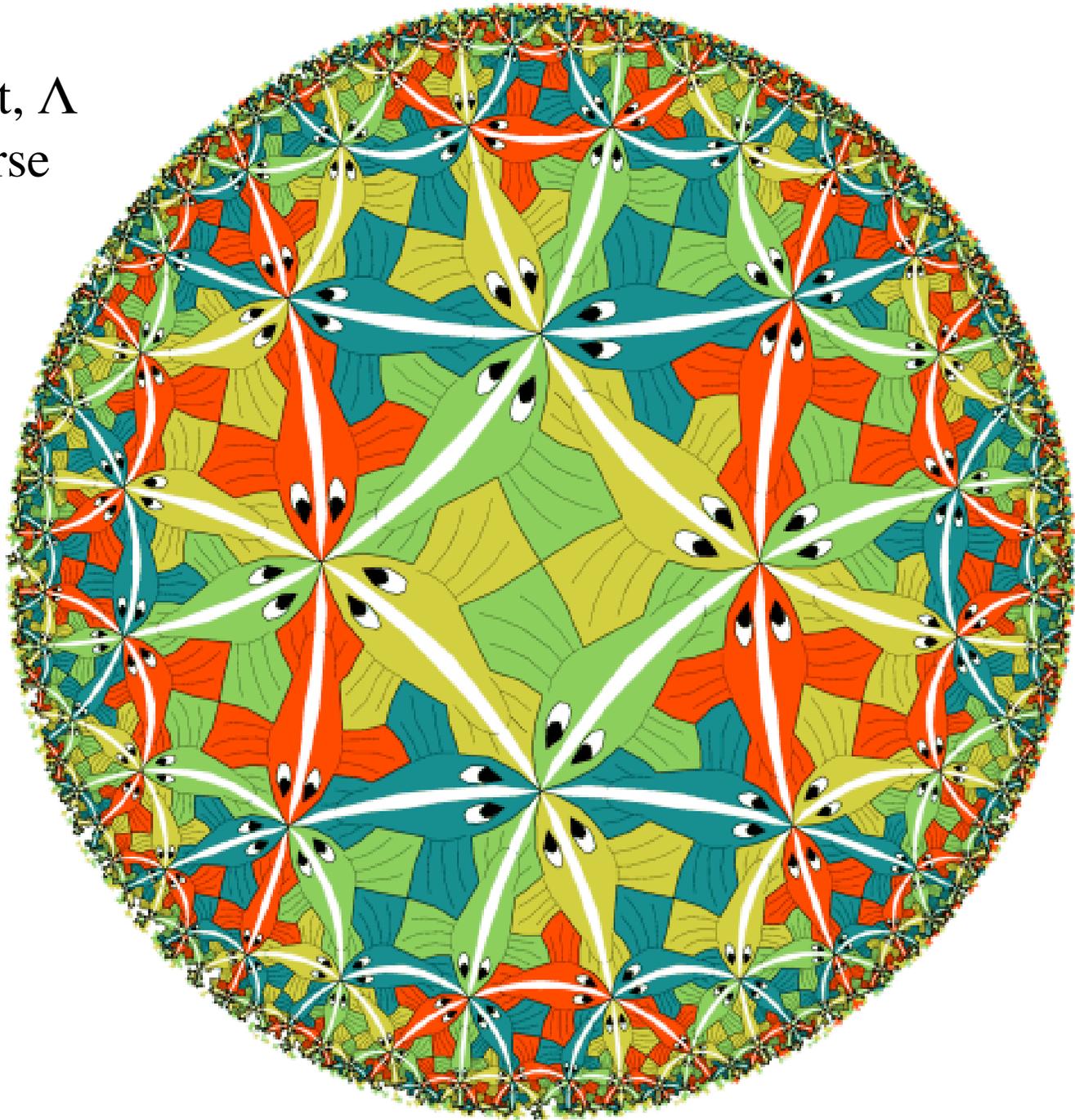
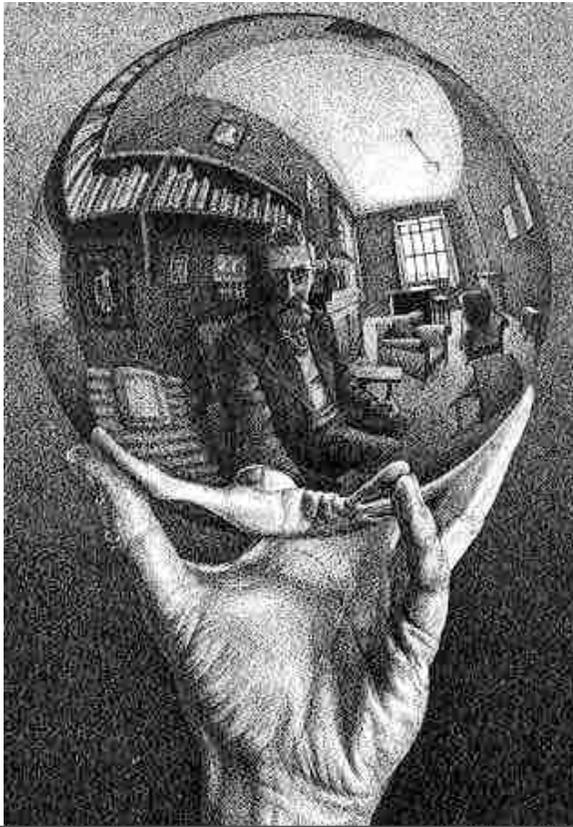


# AdS: Anti de Sitter space

De Sitter space: Gravity with positive cosmological constant,  $\Lambda$   
Accelerated, expanding universe

Anti de Sitter:  $\Lambda < 0$

Spatial cross section of AdS = hyperbolic space



M. K. Escher, courtesy of J. Maldacena

# AdS/CFT Duality

*Most supersymmetric QCD: “ $\mathcal{N}=4$ ” SUSY for SU(N) gluons: 4 supercharges.*

Gluons (spin 1) + 4 spin 1/2 + 6 spin 0, all adjoint rep. No quarks.

One dimensionless coupling,  $\alpha_s$ , but does *not* run! Extraordinary theory:

*No mass scale, both classically and quantum mechanically!*

Conformal Field Theory (CFT). Probably exactly soluble.

*Maldacena’s Conjecture:  $\mathcal{N}=4$  SU( $\infty$ ) dual to string theory on AdS<sub>5</sub> x S<sup>5</sup>*

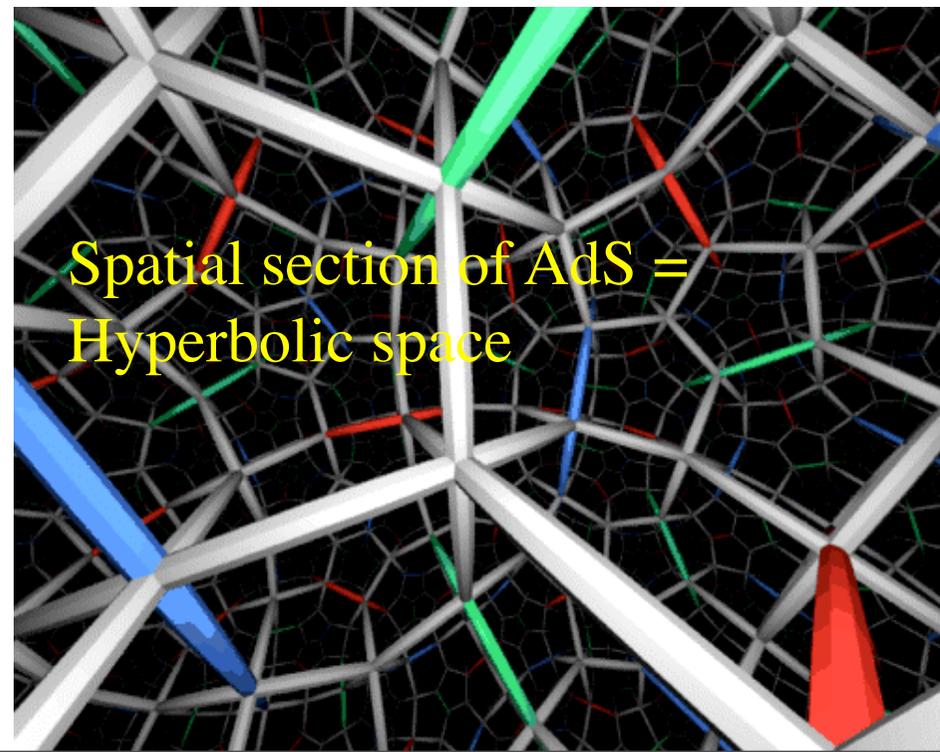
AdS<sub>5</sub> (AdS in 5 dimensions) + S<sup>5</sup> (five sphere)

= Type IIB string in 10 dim.’s

*AdS/CFT duality:* Strong coupling in one theory is weak coupling in the other.

So *strong* coupling for  $\mathcal{N}=4$  SU( $\infty$ ),  $\alpha_s = \infty$  same as *weak* coupling on AdS<sub>5</sub> x S<sup>5</sup>.

Weak coupling string theory = *classical* supergravity!



Spatial section of AdS =  
Hyperbolic space

# AdS/CFT at $T \neq 0$ : “Most perfect fluid”

If one can compute with AdS/CFT,  
often easier for  $\alpha_s = \infty$  than  $\alpha_s \approx 0$ !

Results for  $\mathcal{N} = 4$  SU( $\infty$ ), infinite  $\alpha_s$  :

Pressure 3/4 ideal! CFT  $\Rightarrow$  flat with T.

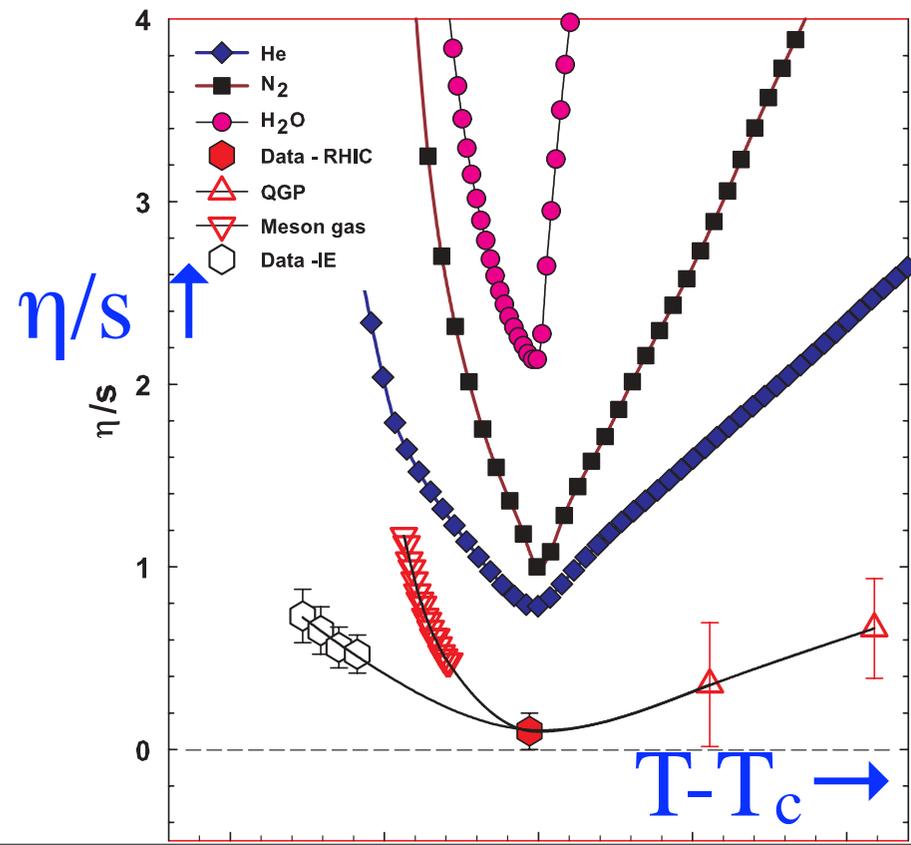
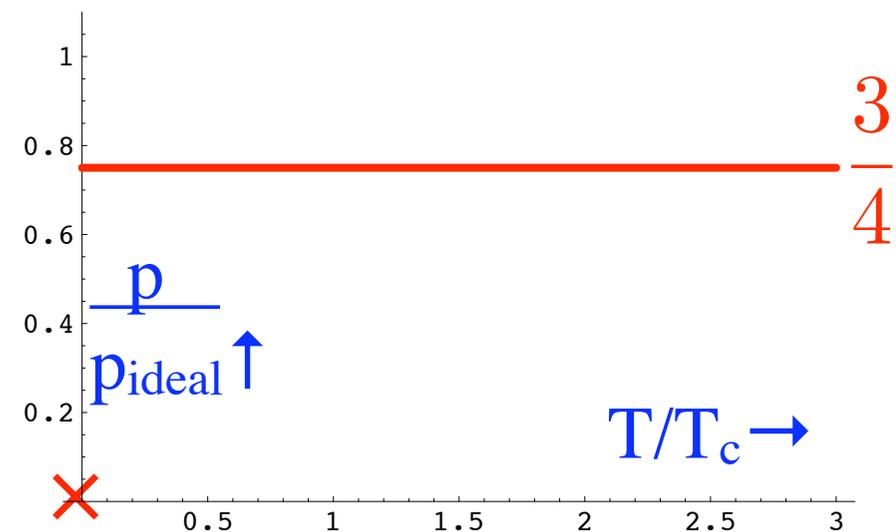
Not only equi., but also transport properties!

$\eta/s = \text{viscosity/entropy} = 1/(4\pi)$

Universal lower bound?

Perhaps.  $\eta \sim 1/\alpha_s^2$  for  $\alpha_s \approx 0$

In many systems,  $\eta/s$  has a minimum  
near  $T_c$ . Maybe in the sQGP?



# Shootout in the Unicorn Corral!

At RHIC, central AA  $\neq$  A (pp) collisions

“Tail wags the Unicorn”

Clearest signal from “high” pt:  
 $R_{AA}$ , jet suppression...

Body of the Unicorn: sQGP

Two packs of fighting dogs (aka theorists):

**Strong coupling!** AdS/CFT rules!

*Weak coupling:*

*need* lattice for equilibrium thermo.  
*need* to extend to non-equi.: NaRPA!

Nothing better than a good dog fight...





*"A possible eureka."*