



New Phases at RHIC: what's the “s” in the sQGP?

Theory:

At high temperature, hadrons (QCD) form new state of matter:

Quark Gluon Plasma (QGP). Near critical temperature: a “s”QGP?

Numerical “experiments” on a lattice: *semi-QGP*?

AdS/CFT duality: super-QCD and the *strong-QGP*?

Experiment:

Collide heavy ions at high energies: RHIC @ BNL; SPS & LHC @ CERN

Suppression of jets in AA collisions. Robust signal of new physics.

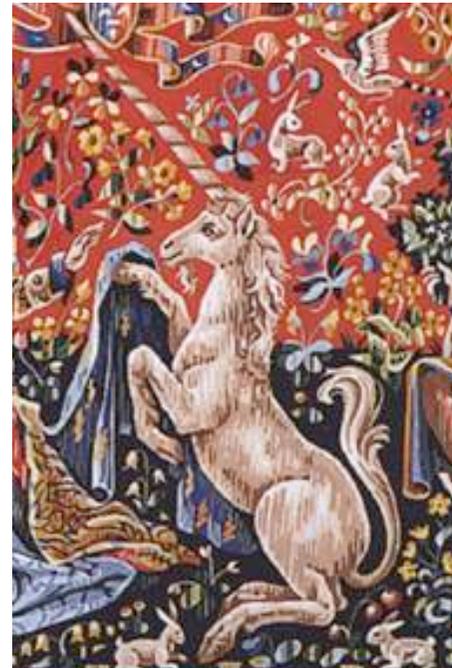
Non central collisions exhibit strong “flow”

Even heavy quarks suppressed, “flow”

Not an ordinary QGP; “Most Perfect Fluid on Earth”?

Wealth of precise experimental data. Theorists agog.

Myths can come true, but maybe not the ones you expect...



QCD & Confinement

QCD (Quantum Chromo Dynamics): quarks and gluons carry *color*.

In vacuum: permanently bound - “*confined*” - into states without color.

Quarks are fermions, come in six “flavors”

Basic mass scale: GeV \sim 12 trillion degrees. GeV = 1000 MeV

Mesons: quark + anti-quark, color + anti-color cancel

Up, down, strange quarks \Rightarrow pions, kaons, eta, eta' + *many* others

Lightest hadrons: masses 140, 500, 550, 960 MeV.

Light because of “chiral” symmetry: up and down q's very light, strange kinda.

Charm, bottom, top quarks \Rightarrow J/Psi, Upsilon... Masses 3.1, 10... GeV

Baryons: three quarks. Color neutrality \Rightarrow three colors.

Up, down quarks \Rightarrow neutrons, protons. 940 MeV

Strange quarks \Rightarrow Lambda...Omega. 1-2 GeV

Symmetries of QCD

Why three colors?

Assume that for *each* quark, multiply by $q \rightarrow z q$, where $z = \exp(2 \pi i/3)$.

Since $z^3 = 1$, three quarks are *invariant* under this $Z(3)$ symmetry.

QCD: Symmetries *local* $SU(3)$, *global* $Z(3)$ (Just like 3-state Potts model)

Quarks: spin 1/2, $SU(3)$ vector. Carry $Z(3)$ charge.

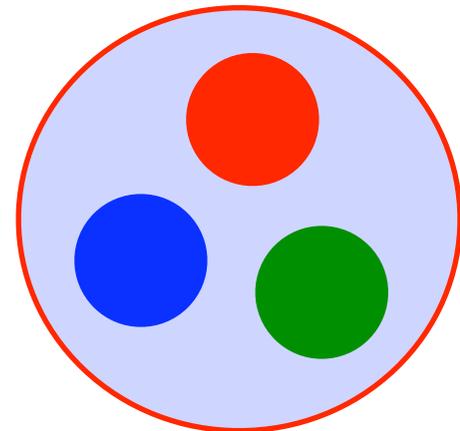
Gluons: spin 1, $SU(3)$ matrix. $3^2 - 1 = 8$ types of gluons. $Z(3)$ charge indirect.

Gluons interact with each other, quarks...*very* complicated!

Confinement related to “hidden” $Z(3)$ symmetry.

“Easy” to see quarks, gluons more obscure.

One model: quarks & gluons confined inside MIT “bag” =>



QCD & Asymptotic Freedom

Unlike a photon, the 8 gluons in QCD interact with one another

Complicated at large distances, *very simple at short distances* (large momenta)

Asymptotically Free: QCD coupling, α_s , “runs”, vanishes as momentum $k \rightarrow \infty$:

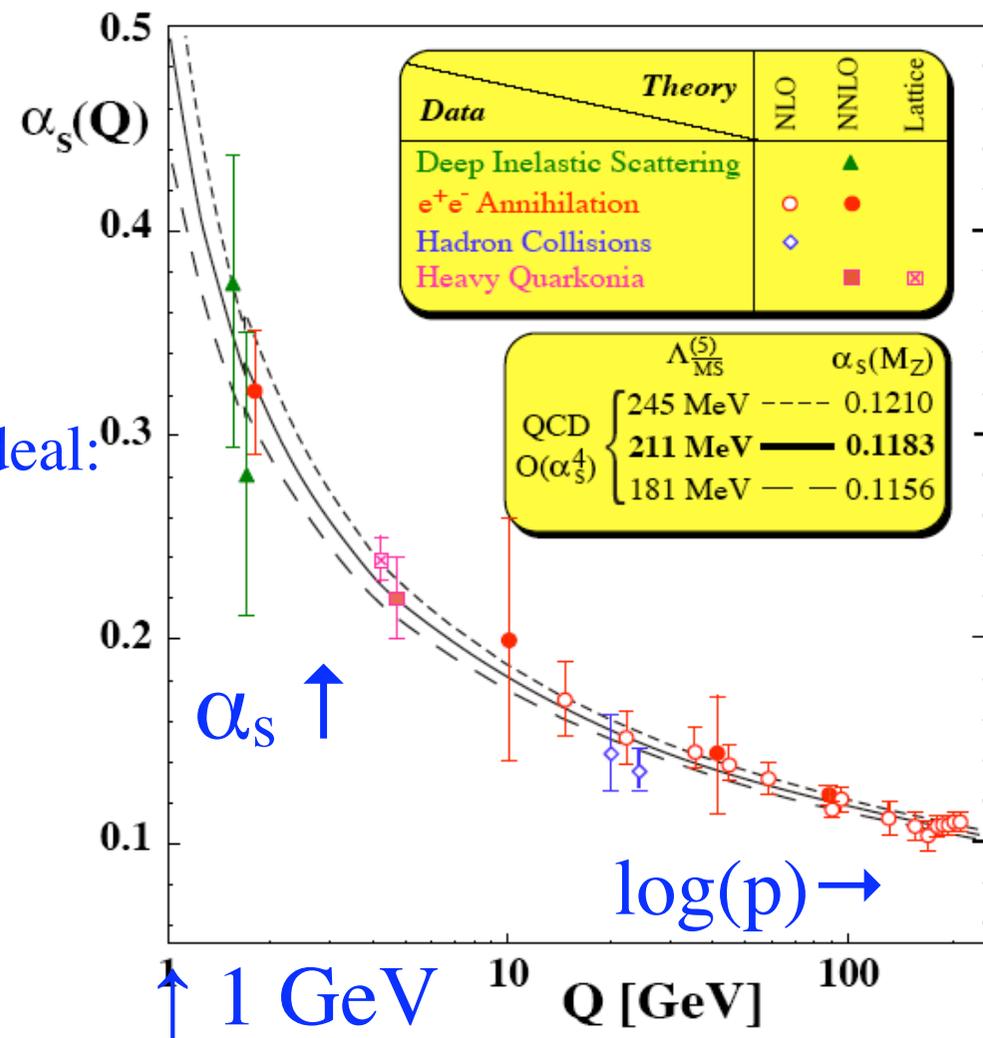
$$\alpha_s(k) \sim \# / \log(k), \quad k \rightarrow \infty$$

Conversely, $\alpha_s(k)$ *increases* as momentum decreases (large distances). Experiment:

As temperature $T \rightarrow \infty$, pressure $p(T) \rightarrow$ ideal:

$$p_{ideal}(T) = \left(8 + \frac{7}{8} \cdot 18 \right) \frac{\pi^2}{45} T^4$$

At *high T*, QCD forms *nearly ideal*
Quark-Gluon Plasma, QGP: *deconfined*

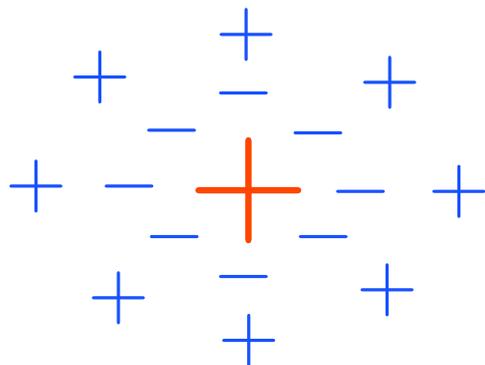


Coulombic Plasmas, *partial* ionization

Ordinary matter (gas, liquid, solid) composed of electrically neutral objects.

“Fourth” state: plasma. No net charge, but charges *free* to move about.

Coulomb int. with test charge, $+$, *shielded* by free charges, $-$ $+$, with density “n”:


$$\frac{e^2}{r} \rightarrow e^{-m_{Debye} r} \frac{e^2}{r} \quad ; \quad m_{Debye}^2 \sim \frac{e^2 n}{T}$$

Ordinary matter: if *no* ionization, Coulomb int. remains $1/r$, charge modified.

In plasma, distinguish between *complete* or *partial* ionization.

Interesting dynamics in regime with partial ionization. Hard to treat cleanly.

Z(3) charges in Gluonic Plasmas

QCD plasma: Coulomb int.'s + much more. Consider Z(3) charges:

Use propagator of test quark, ~ “Polyakov Loop”

$$\ell(x) = \frac{1}{3} \text{tr} \mathcal{P} \exp \left(ig \int_0^{1/T} A_0(x, \tau) d\tau \right)$$

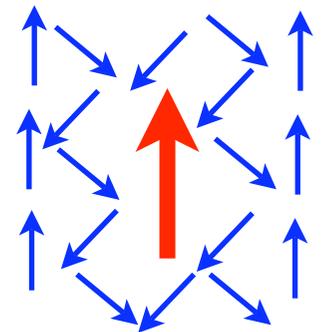
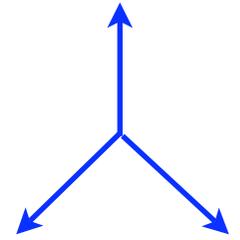
In purely gluonic plasma, *only* Polyakov Loops carry Z(3) charge. (*t Hooft)

Polyakov Loop ~ Z(3) magnetization of Potts model

Use Loop as probe of Z(3) test charge.

High temperature: $g \sim 0$, so Loop ~ 1 . Debye screening.

Loop is order parameter for deconfinement!



Deconfinement in Gluonic Plasma

Zero temperature: confinement \Rightarrow $Z(3)$ charge can't propagate.

$\langle \text{Loop} \rangle = 0$: strict order parameter for confinement

Confined phase: *symmetric* phase of $Z(3)$ spins

No ionization of $Z(3)$ charge.

$T \rightarrow \infty$: test quark \sim free $\Rightarrow \langle \text{Loop} \rangle \sim 1$; $Z(3)$ charge *completely* ionized.

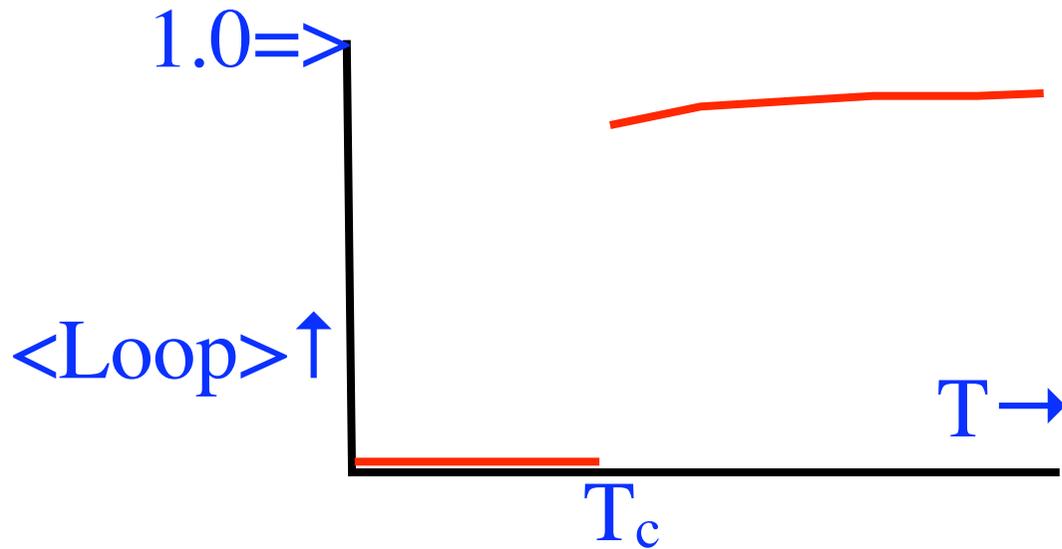
$\Rightarrow \langle \text{Loop} \rangle$ *nonzero* at some T_c = deconfinement transition temp.

Deconfined phase, $T > T_c$: *broken* phase of $Z(3)$ spins.

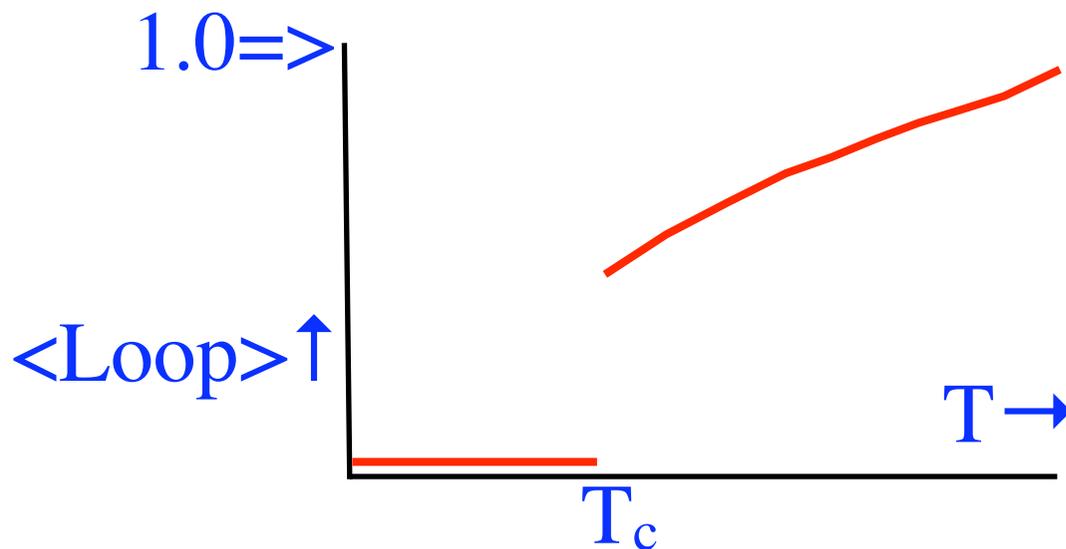
$Z(3)$ charge ionized.

How big is $\langle \text{Loop} \rangle(T_c^+)$? At T_c^+ , is ionization complete, or partial?

Ordinary or “Semi”-Gluonic Plasma?



$\langle \text{Loop} \rangle(T_c^+) \sim 1$:
complete ionization,
ordinary Gluonic Plasma,
for all $T > T_c^+$



$\langle \text{Loop} \rangle(T_c^+) < 1$:
partial ionization at T_c^+

If $\langle \text{Loop} \rangle(\# T_c) \sim 1$

“Semi”-Gluonic Plasma,
 T_c^+ to $\# T_c$

Ordinary Gluonic Plasma,
 $T > \# T_c$.

Numerical “experiments”: the lattice

How to compute non-perturbatively? Put QCD on a lattice, with spacing “ a ”.

In QCD, AF + renormalization group \Rightarrow *unique* answer as $a \rightarrow 0$, continuum limit

Approximate ∞ -dim. integral by finite dim.: “Monte Carlo”

In practice: need to “improve” to control lattice effects $\sim a^2$, etc.

For pure glue near continuum limit from mid-’90’s!

Compute equilibrium thermodynamics: pressure(temperature) = $p(T)$, loop(s)

Lattices: 6-8 steps in time, 24-32 in space $\Rightarrow 8 \cdot 32^3 \cdot 8 = 10^6$ dim. integral

With dynamical quarks, much harder: today, *not* near continuum for light quarks

Quarks non-local, difficult to treat (global!)

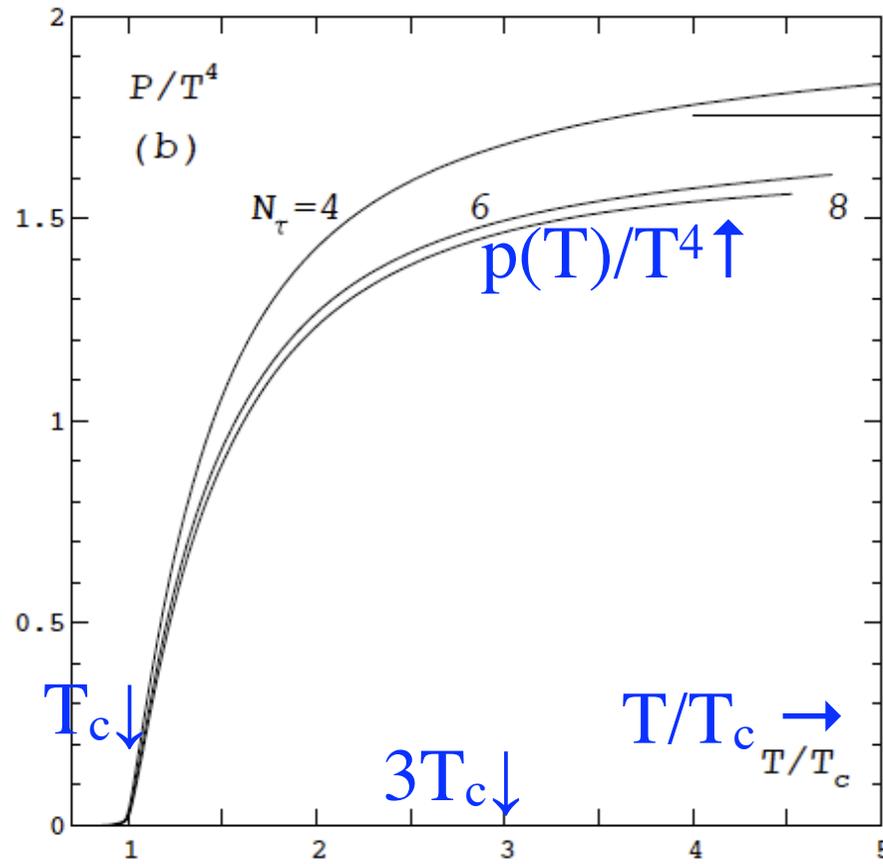
“chiral” symmetry of light quarks

Today: only $T \neq 0$, *not* dense quarks at $T=0$

With quarks: *orders* of magnitude more difficult



Lattice: SU(3) gluonic plasma



$N_\tau =$ time steps:
4,6,8. 6~8

arXiv:hep-lat/9602007, Boyd et al

Lattice: SU(3) Gluonic Plasma (no quarks)

Confined phase at $T=0$: only “glueballs”. *Lightest* glueball 1.5 GeV

Deconfinement at $T_c \sim 270\text{MeV} \pm 10\%$ *Small* relative to glueball masses.

Pressure *very* small below T_c , large above T_c (like large # colors)

Transition weakly first order ($\sim Z(3)$ Potts model)

Pressure: MIT bag model does *not* work

$$p_{\text{MIT}}(T) = \#T^4 - B$$

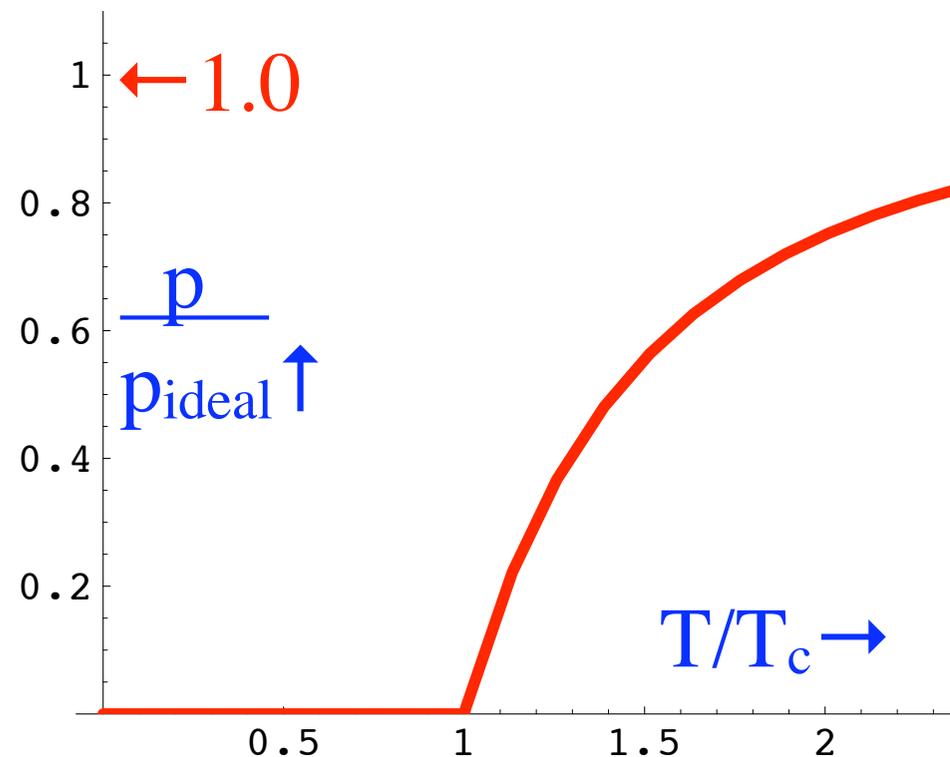
“Fuzzy” bag does:

$$p_{\text{fuzzy}}(T) = \#(T^4 - T_c^2 T^2)$$

$\sim T^4$ perturbative: $\# \sim 0.9$ ideal gas value

$\sim T^2$ *non*-perturbative.

Pressure $\sim 85\%$ ideal by $\sim 4 T_c$



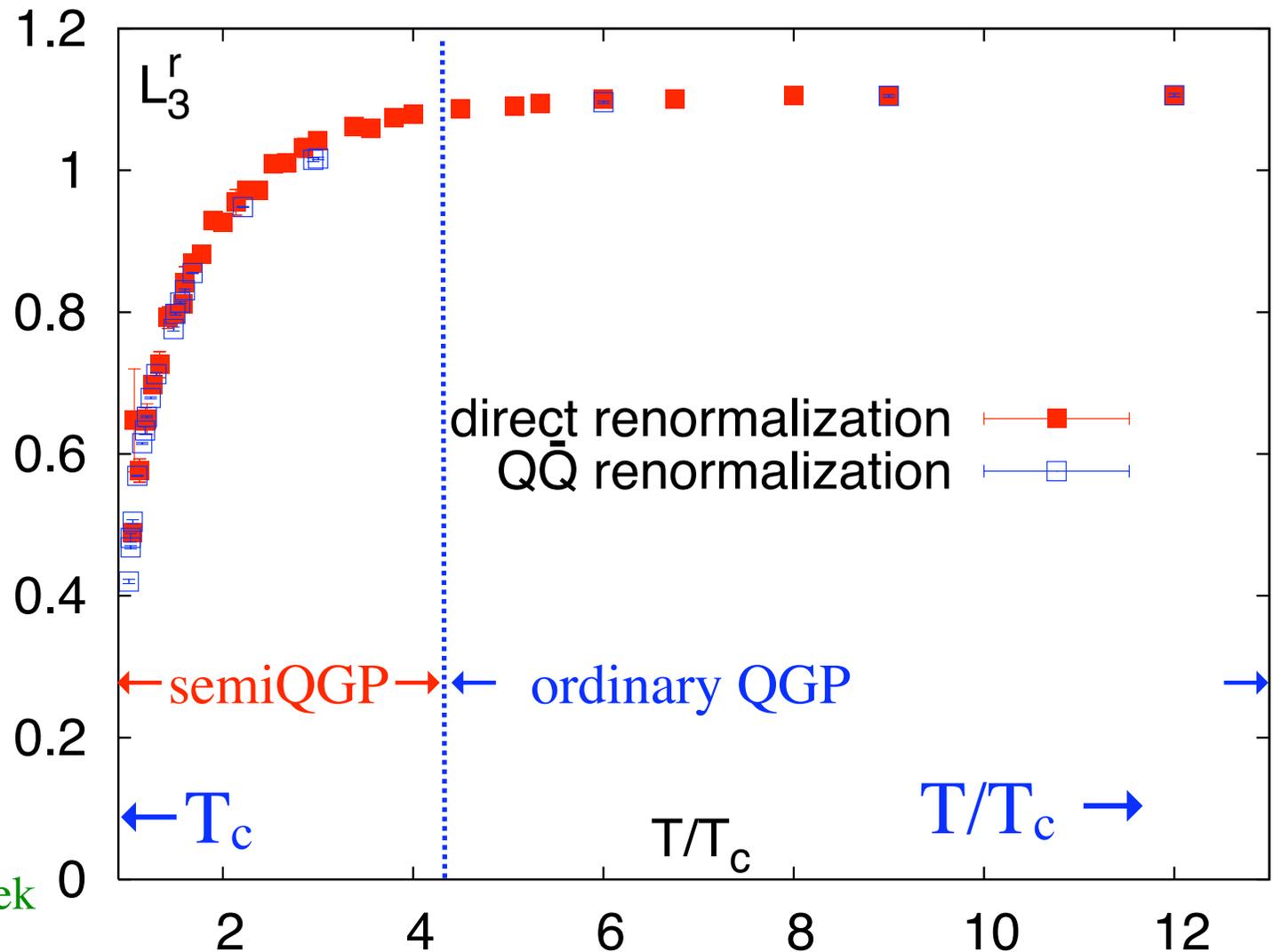
Lattice: SU(3) “Semi”-Gluonic Plasma

$T < T_c$: $\langle \text{Loop} \rangle = 0$, confined phase.

$\langle \text{Loop} \rangle(T_c^+) \sim 0.5$! $\langle \text{Loop} \rangle < 1$ for T_c to $\sim 4 T_c$: “semi”-Gluonic Plasma.

$T > \sim 4 T_c$: ordinary QGP . Polyakov Loop ~ 1 (+ pert. corr.'s)

(Ren.d)
Polyakov
Loop \uparrow



arXiv:0711.2251

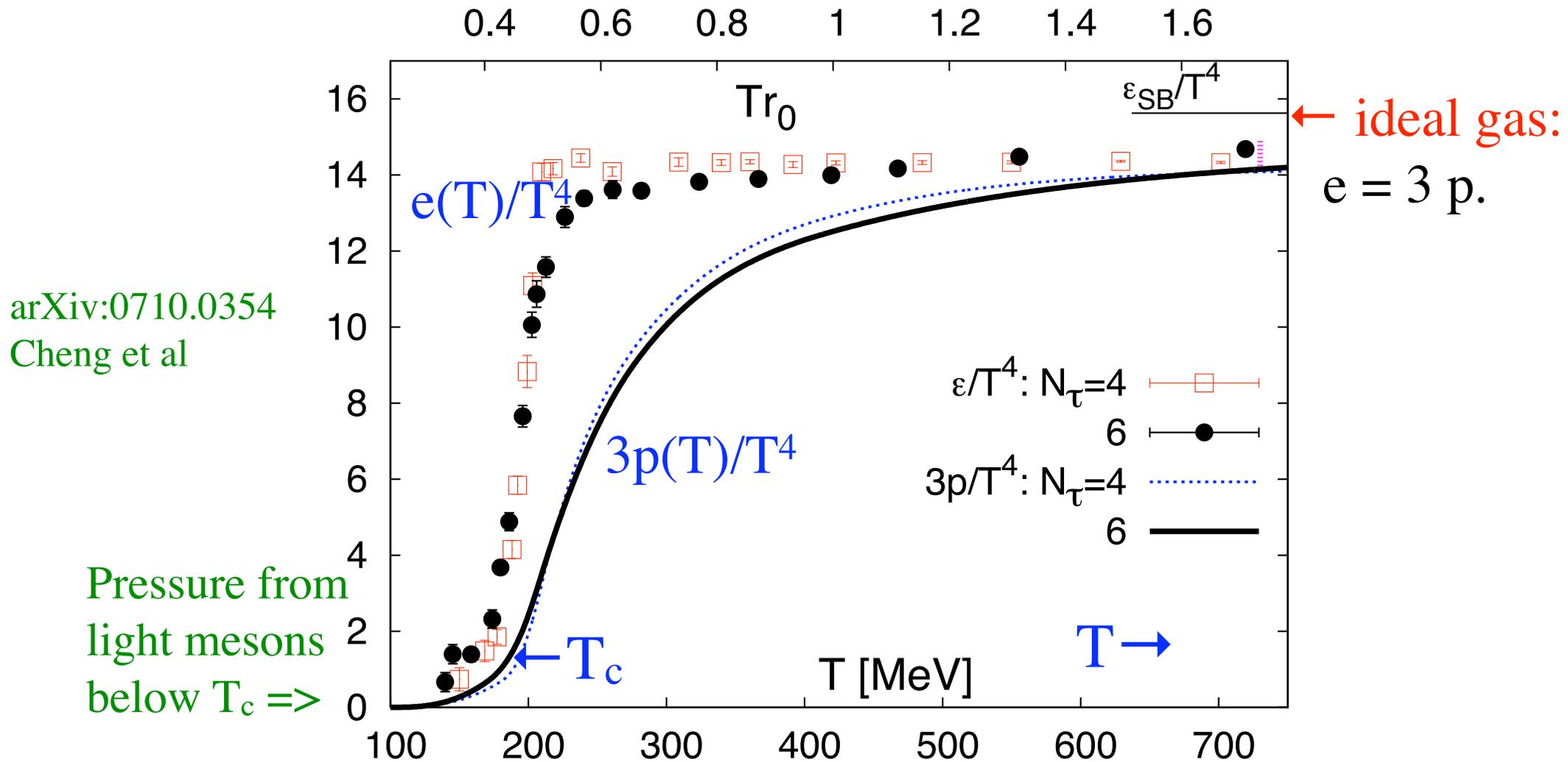
Gupta, Hubner, Kaczmarek

Lattice: Quark-Gluon Plasma, “2+1” flavors

$T_c \sim 190$ MeV. Crossover, no true phase transition. Huge increase in pressure:

$$p(T) = c_{pert} T^4 - B_{fuzzy} T^2 - B_{MIT}$$

Energy density $e(T)$, 3 x pressure(T), each/ T^4 . Pressure $\sim 90\%$ ideal by $\sim 3T_c$.



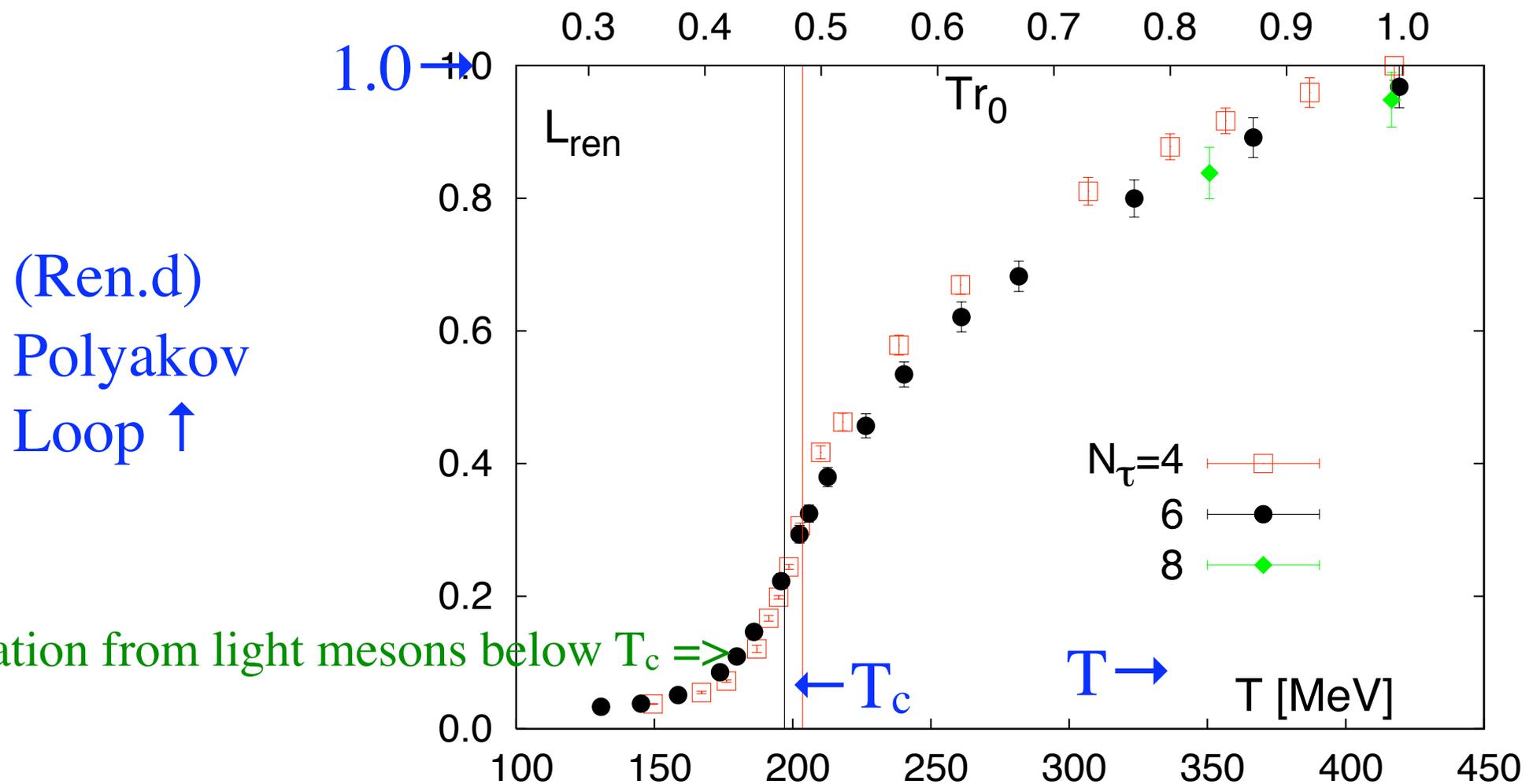
Lattice: "Semi"-QGP, 2+1 flavors

Quarks carry $Z(3)$ charge, and so partially ionize QGP, even *below* T_c

Lattice: *Moderate* ionization below T_c . Loop ~ 0 below $\sim 0.8 T_c$.

$\langle \text{Loop} \rangle(T_c) \sim 0.3$. $\langle \text{Loop} \rangle \sim 1.0$ at $2 T_c$. Semi-QGP for $0.8 \rightarrow 2.0 T_c$

arXiv:0710.0354, Cheng et al

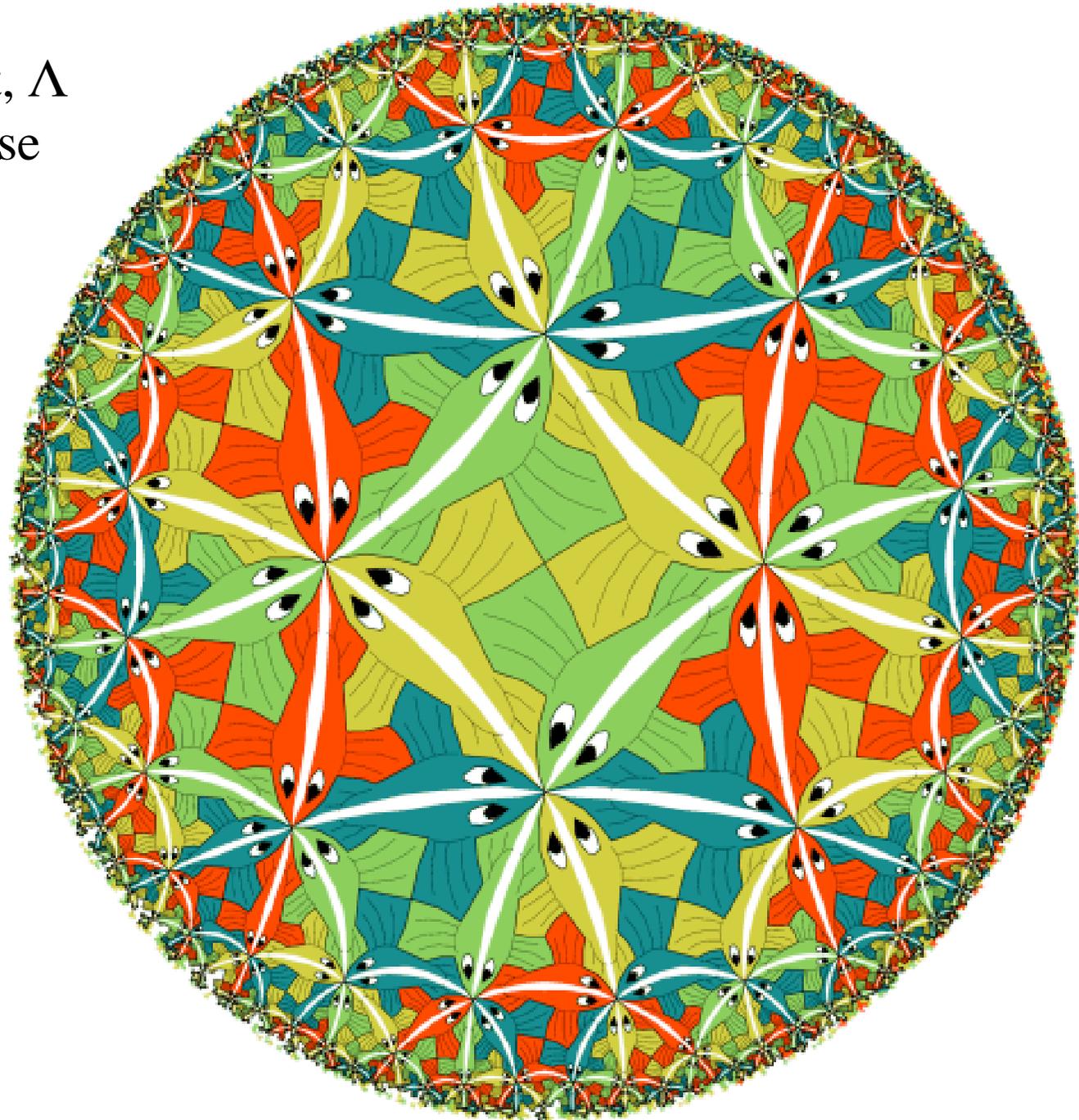
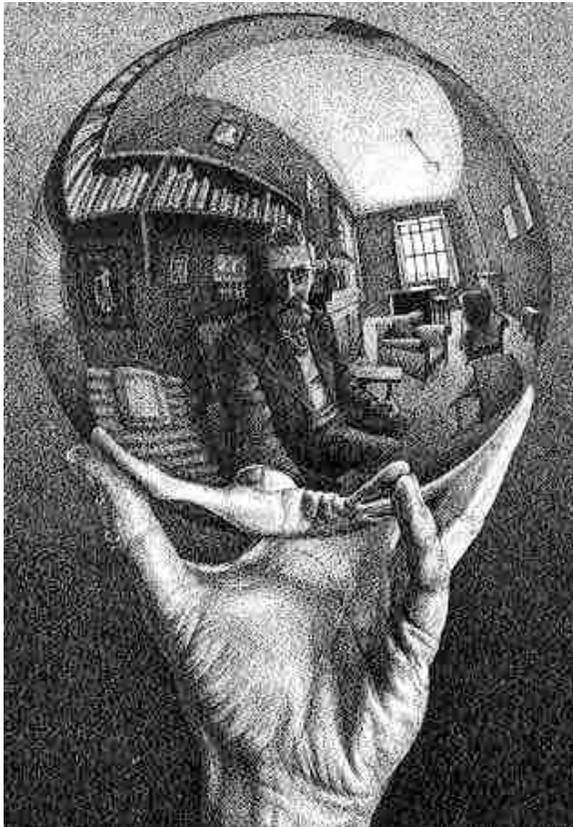


AdS: Anti de Sitter space

De Sitter space: Gravity with positive cosmological constant, Λ
Accelerated, expanding universe

Anti de Sitter: $\Lambda < 0$

Spatial cross section of AdS = hyperbolic space



M. K. Escher, courtesy of J. Maldacena

AdS/CFT Duality

Most supersymmetric QCD: “ $\mathcal{N}=4$ ” SUSY for SU(N) gluons: 4 supercharges.

Gluons (spin 1) + 4 spin 1/2 + 6 spin 0, all SU(N) matrices. No quarks.

One dimensionless coupling, α_s , but does *not* run! Extraordinary theory:

No mass scale, both classically and quantum mechanically!

Conformal Field Theory (CFT). Probably exactly soluble.

Maldacena’s Conjecture: $\mathcal{N}=4$ SU(∞) *dual* to string theory on AdS₅ x S⁵

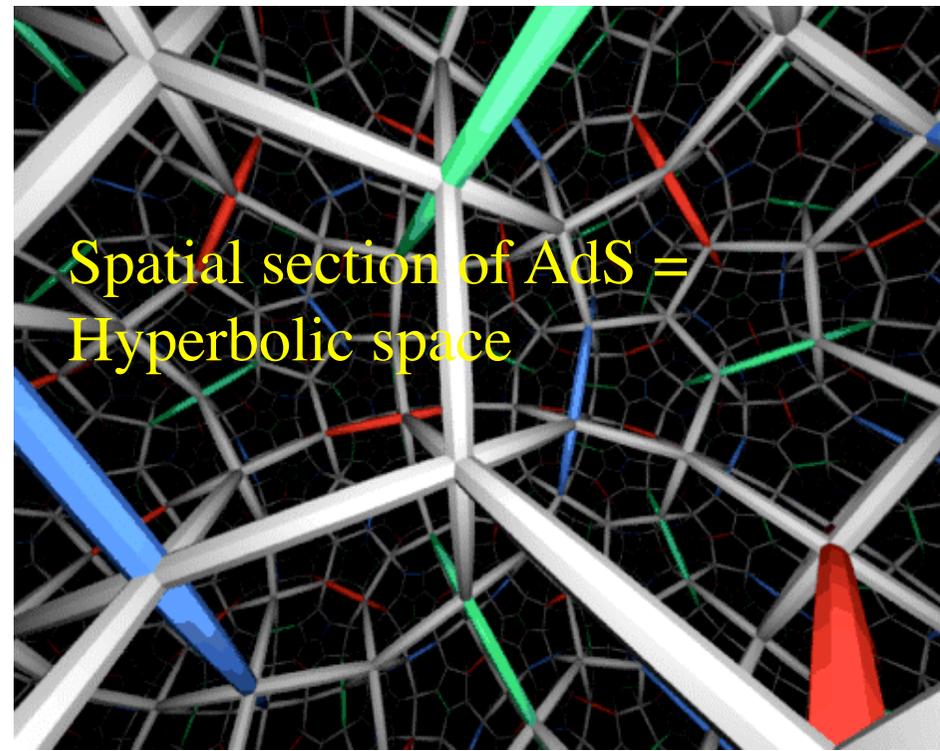
AdS₅ (AdS in 5 dimensions) + S⁵ (five sphere)

= Type IIB string in 10 dim.’s

AdS/CFT duality: Strong coupling in one theory is weak coupling in the other.

So *strong* coupling for $\mathcal{N}=4$ SU(∞) same as *weak* coupling on AdS₅ x S⁵.

Weak coupling string theory = *classical* supergravity!



SUSY QCD and “strong”-QGP

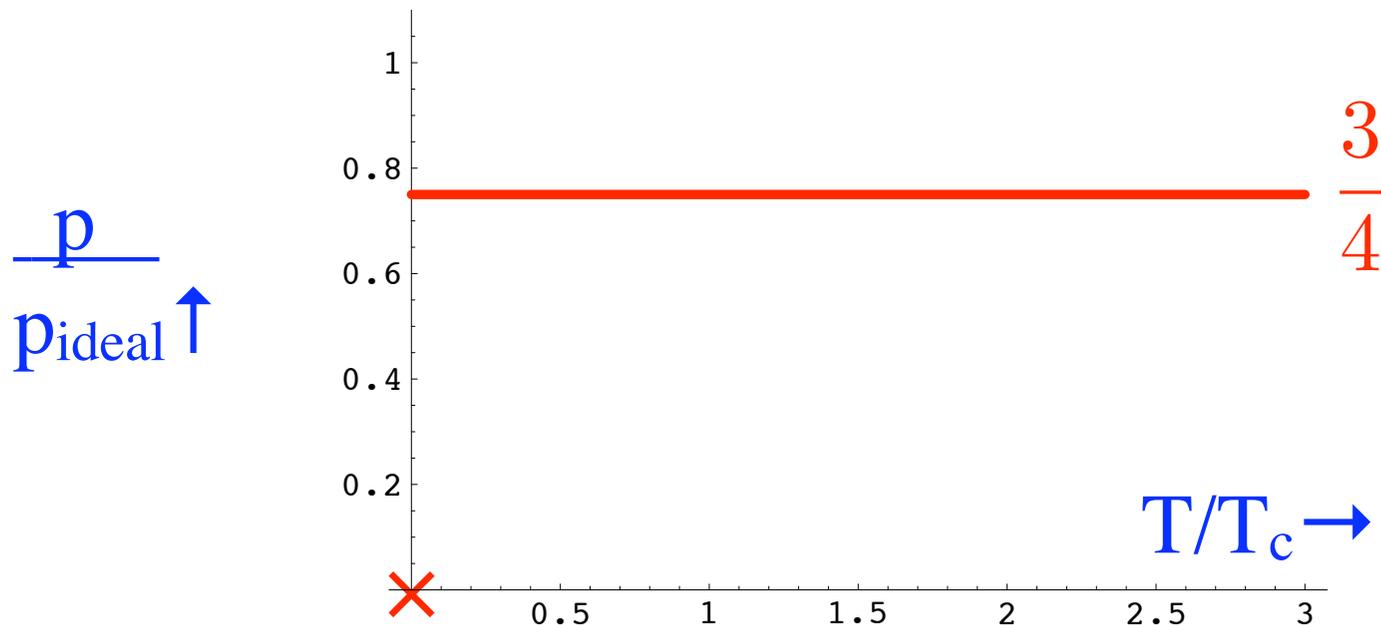
If one can compute with AdS/CFT, often easier for $\alpha_s = \infty$ than $\alpha_s \approx 0$!

Results for $\mathcal{N} = 4$ SU(∞), infinite α_s : pressure = 3/4 ideal.

Conformal field theory \Rightarrow p/T^4 flat with T.

Deconfined phase: heavy quark potential Coulombic. Polyakov Loop?

QCD (at 3 T_c): α_s is 300 x α_{em} in QED. So near T_c , take $\alpha_s = \infty$, “strong”-QGP



Shear viscosity in s-QGP

$\mathcal{N} = 4$ SU(∞), infinite α_s : Also transport properties: shear viscosity, η

$\eta/s = \text{viscosity/entropy} = 1/(4\pi)$. Universal bound?

arXiv:hep-th/0104066. Policastro, Son, Starinets

In non-relativistic QED plasmas,
 η has a *minimum* in a dense regime,
 “strongly coupled” plasma =>

Perhaps in strong-QGP?

nucl-th/0701002, Mrowczynski & Thomas

In weak coupling, $\eta \sim T^3/\alpha_s^2$ for $\alpha_s \approx 0$.

Semi-QGP:

η suppressed by $\langle \text{Loop} \rangle^2$ when $\langle \text{Loop} \rangle$ small.

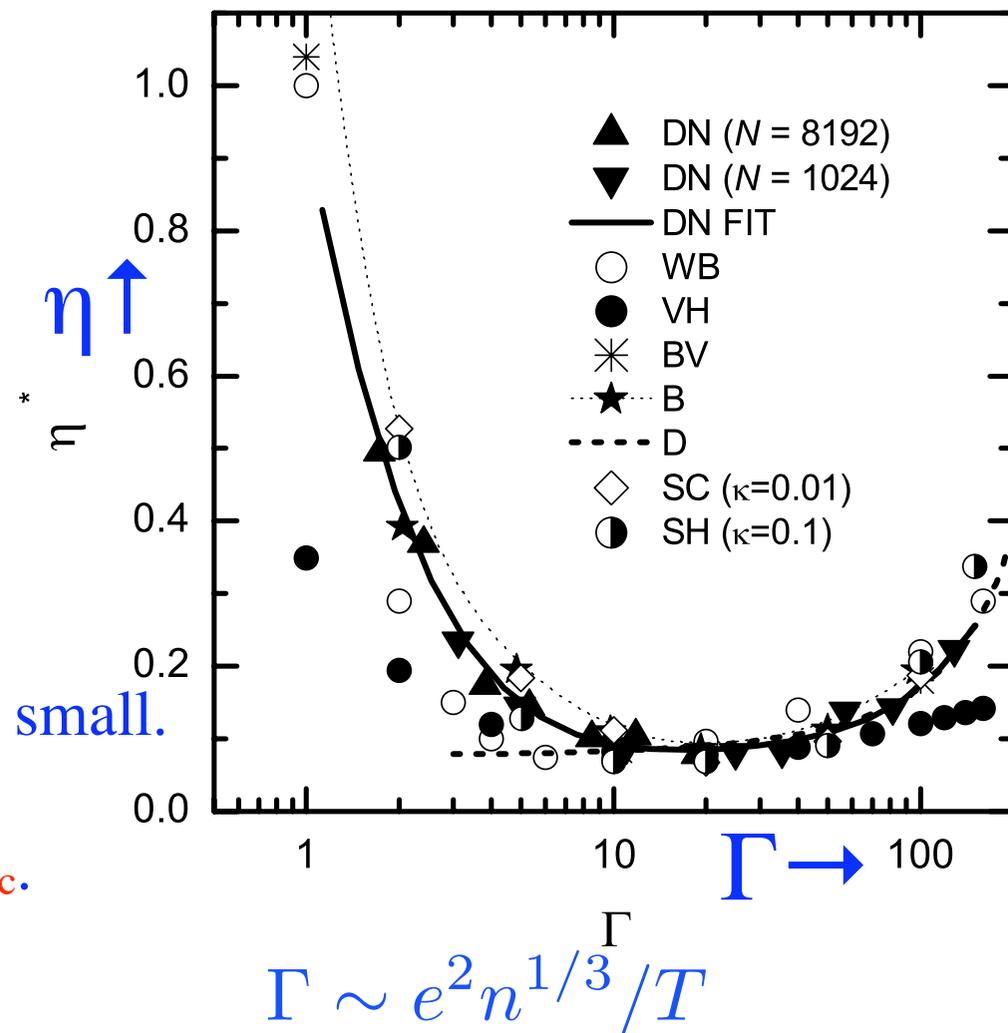
Minimum near T_c ?

=> Large increase in η/s as $T: T_c \Rightarrow 2 T_c$.

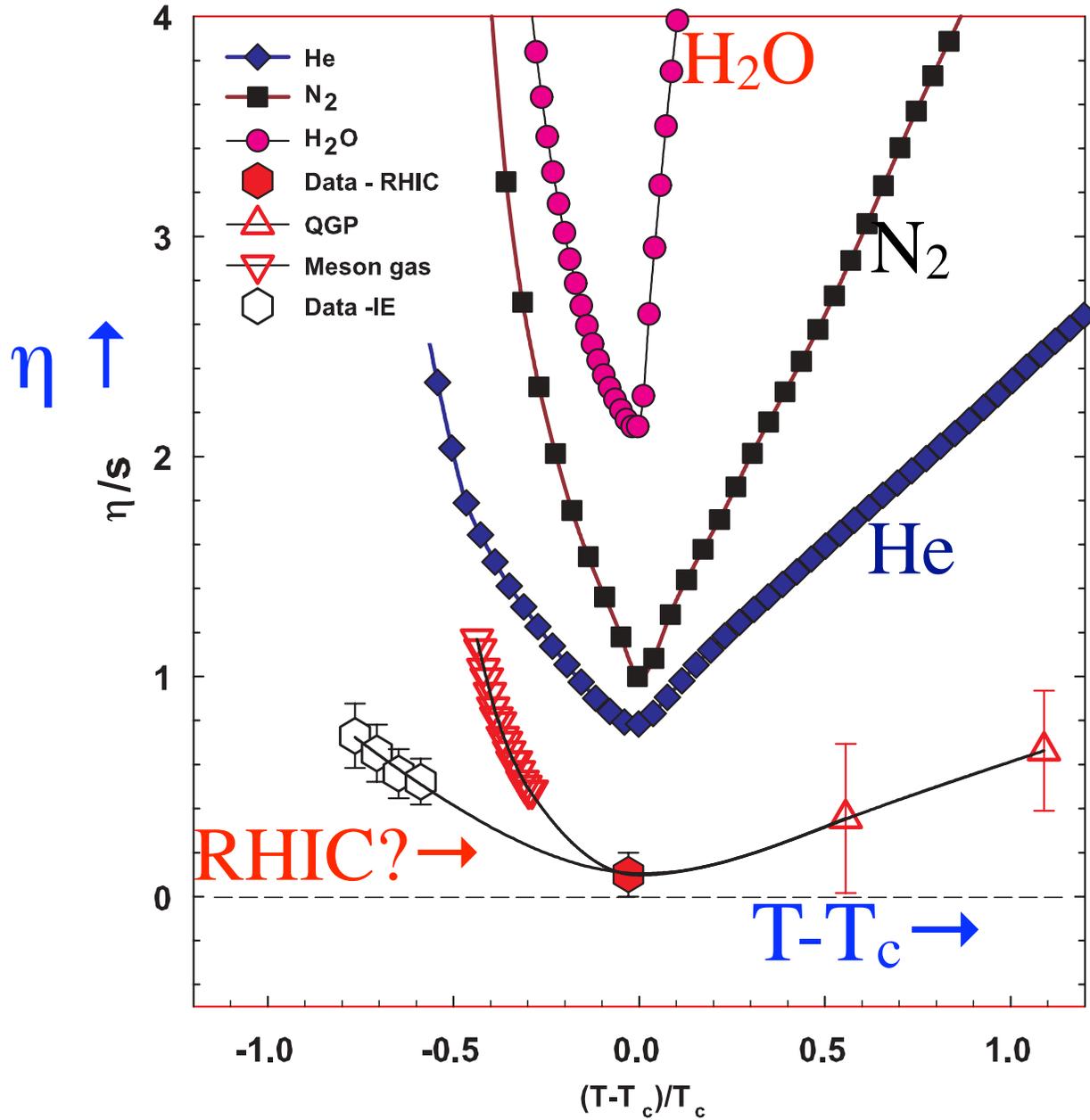
Yoshimasa Hidaka & RDP '08.

arXiv:0710.5229,

Donko, Hartmann, Kalman

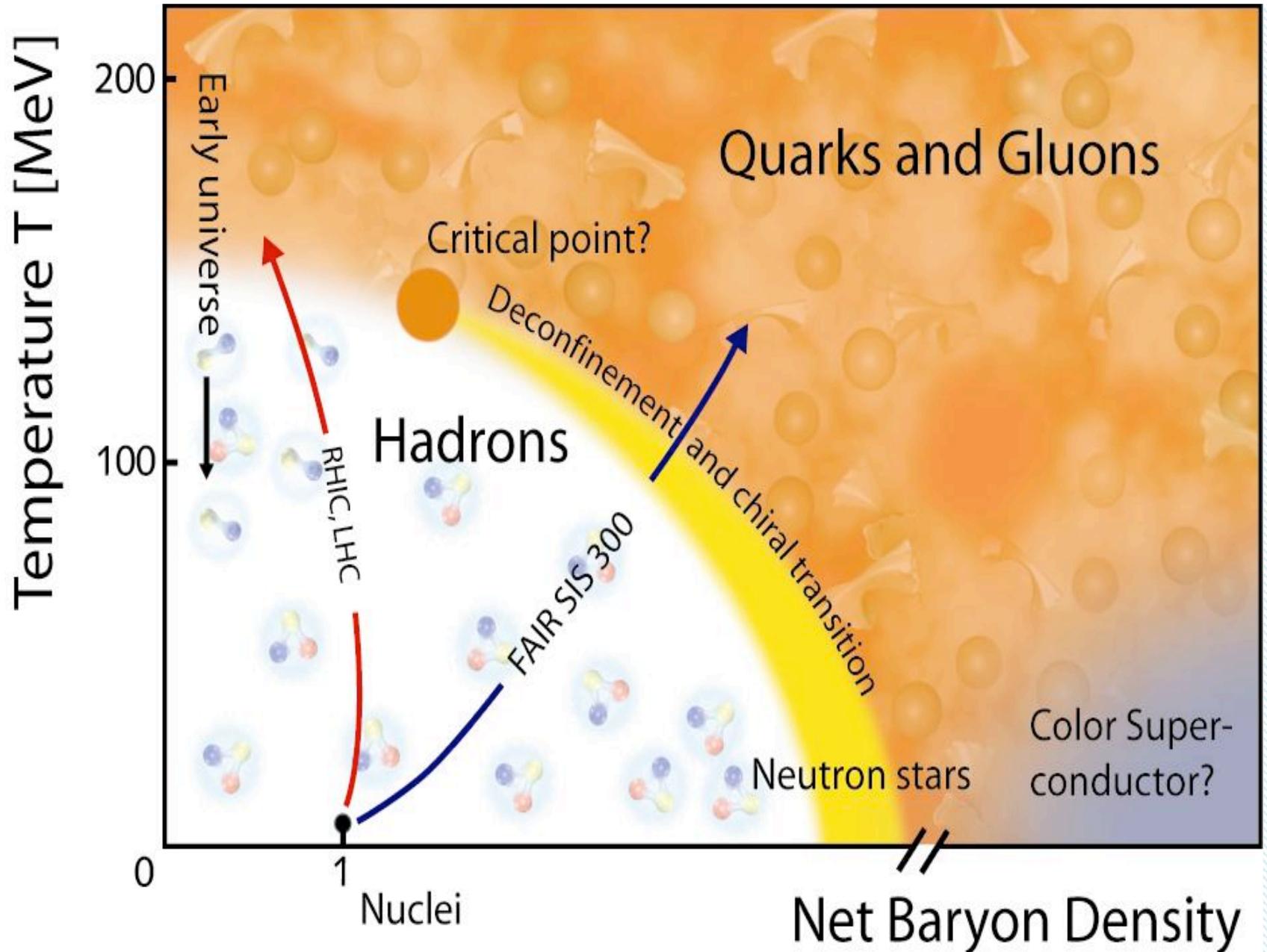


Minimum in η common



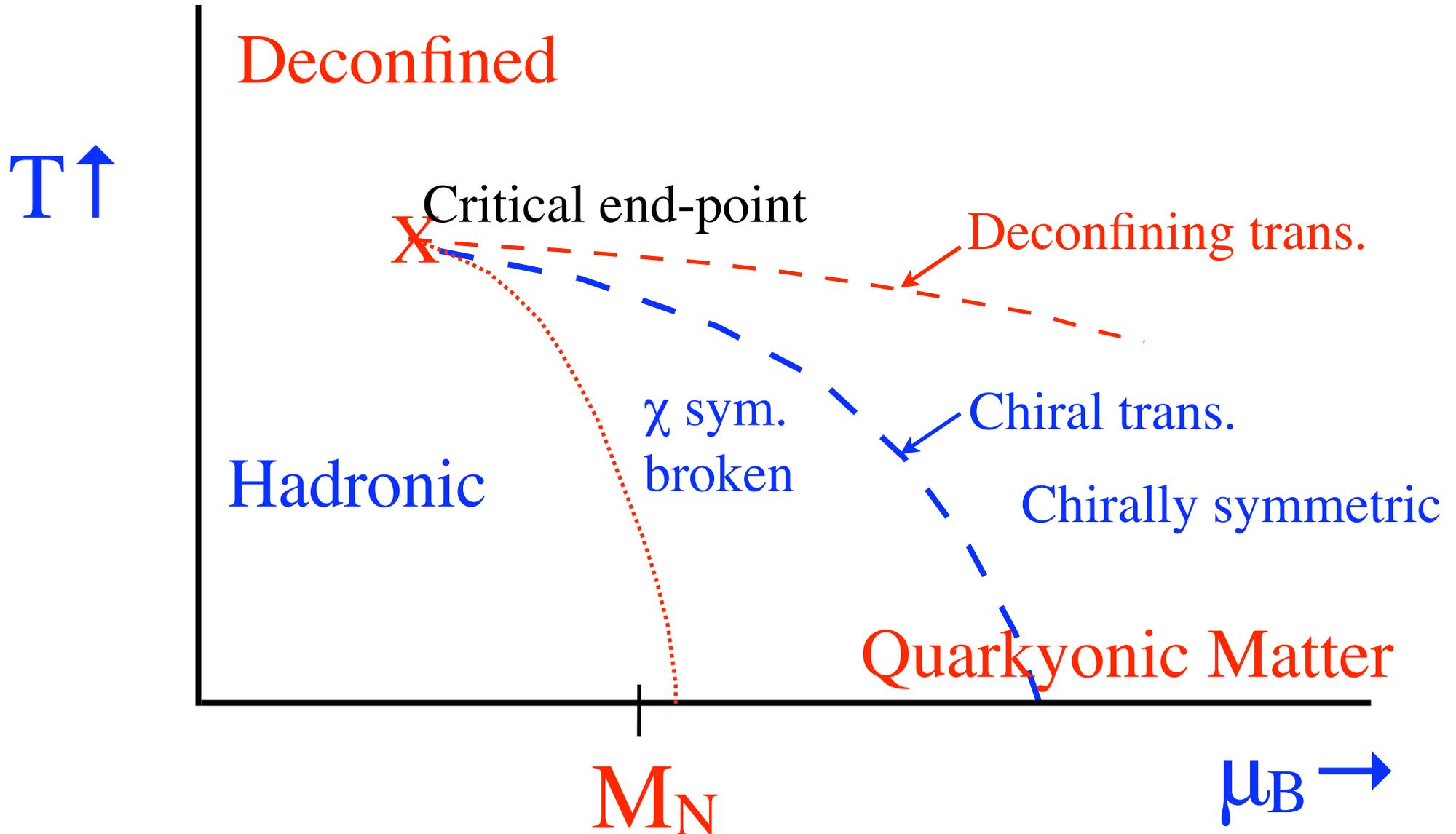
Usual phase diagram

In plane of T and baryon density: critical end point? Rajagopal, Shuryak, Stephanov



“Quarkyonic” matter

New phase diagram: novel behavior at low temperatures, high quark density:
“Quarkyonic” matter: quarks in Fermi sea, but baryons at Fermi surface. Valid
for infinite # colors - for QCD? McLerran and RDP '07.



Hunting for the “Unicorn” in Heavy Ion Collisions



Unicorn = QGP. Hunters = experimentalists. So: all theorists are dogs...

AA collisions at high energies

Collide: pp, protons on protons. Benchmark for “ordinary” QCD.

AA, nuclei on nuclei. Atomic # “A”: 60 => 200, Cu -> Au, Pb. “Hot” nuclei.

pA, proton (or deuteron) on nucleus. Another check: QCD in “cold” nuclei

Why AA? Baryons are like hard spheres, so for A: 60 - 200,
transverse size $A^{2/3}$: 15 - 35 \times proton. **Big nuclei are big!**

Total energy in the center of mass, $E_{\text{cm}} = \sqrt{s}$ (GeV); per nucleon, $\sqrt{s}/A = \sqrt{s_{\text{NN}}}$.

SPS @ CERN	5 => 17 GeV
RHIC @ BNL	20 => 200 GeV
LHC @ CERN	5500 GeV
SIS @ GSI	2 => 6 GeV

SPS, SIS Fixed Target
RHIC, LHC Colliders

LHC > '08

SIS @ GSI, Darmstadt > '12



Geometry of AA collisions at high energy

At *high energies*: nuclei Lorentz contracted along beam (15 fermi \Rightarrow 0.3 fermi)

AF \Rightarrow nuclei don't stop, pass through each other.

Collider: lab = center of mass frame

Momenta of produced particles: along beam, p_z ; transverse to beam, p_t

Baryons in original nuclei go down beam pipe, at large $\pm p_z$

For pp collisions: particles \sim constant for some range in p_z about $p_z = 0$:

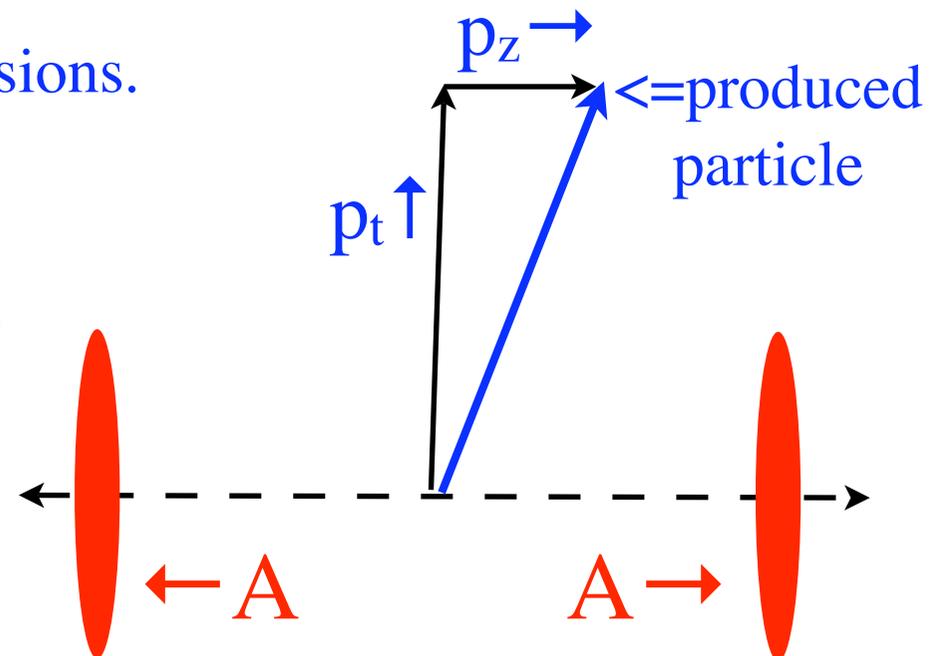
“rapidity plateau” (rapidity $\sim p_z$; boost invariance)

Bjorken: look at rapidity plateau in AA collisions.

Rapidity plateau \sim free of incident baryons.

\Rightarrow most likely to be at nonzero temperature,
zero (quark) density.

Collider: central plateau 90° to beam



Typical Au-Au collision @ RHIC

Experiments @ RHIC:

“Big”: ~ 400 people. STAR & PHENIX

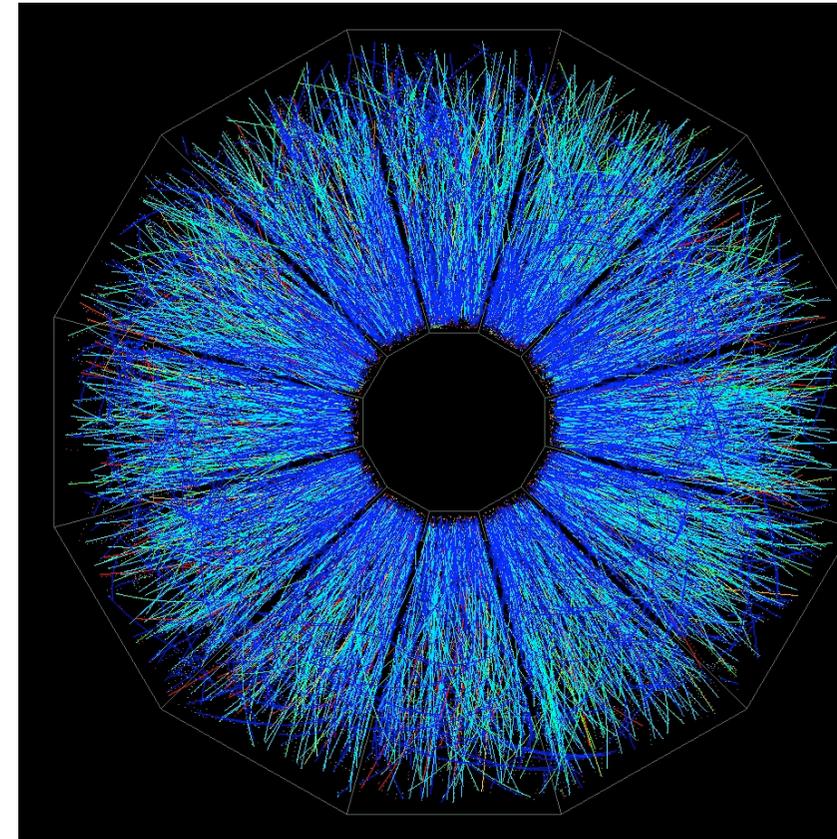
“Small”: ~ 50 people. PHOBOS & BRAHMS

No surprises in total multiplicity; ~ 1.3 A × pp:

total # particles ~ ~ log(total energy)

~total # experimentalists

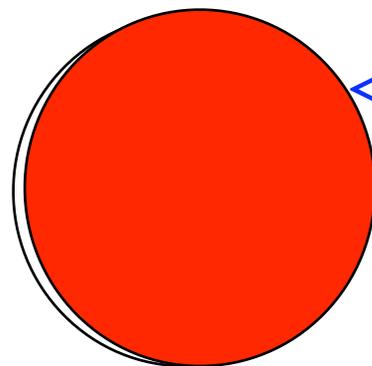
theorists ~ log(log(total energy)).



Need hunters more than dogs...

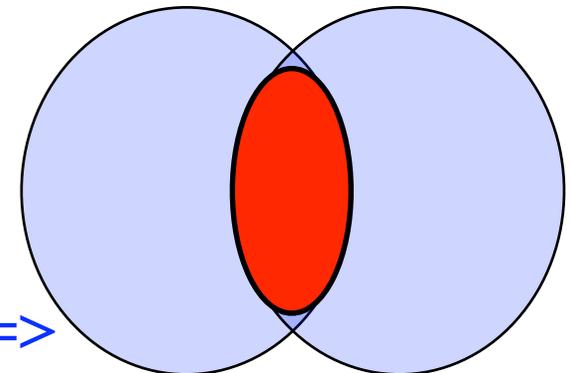
Total # particles/unit rapidity ~900 ↑

Also: can exp.'y measure how much nuclei overlap in transverse plane



←= central

peripheral =>



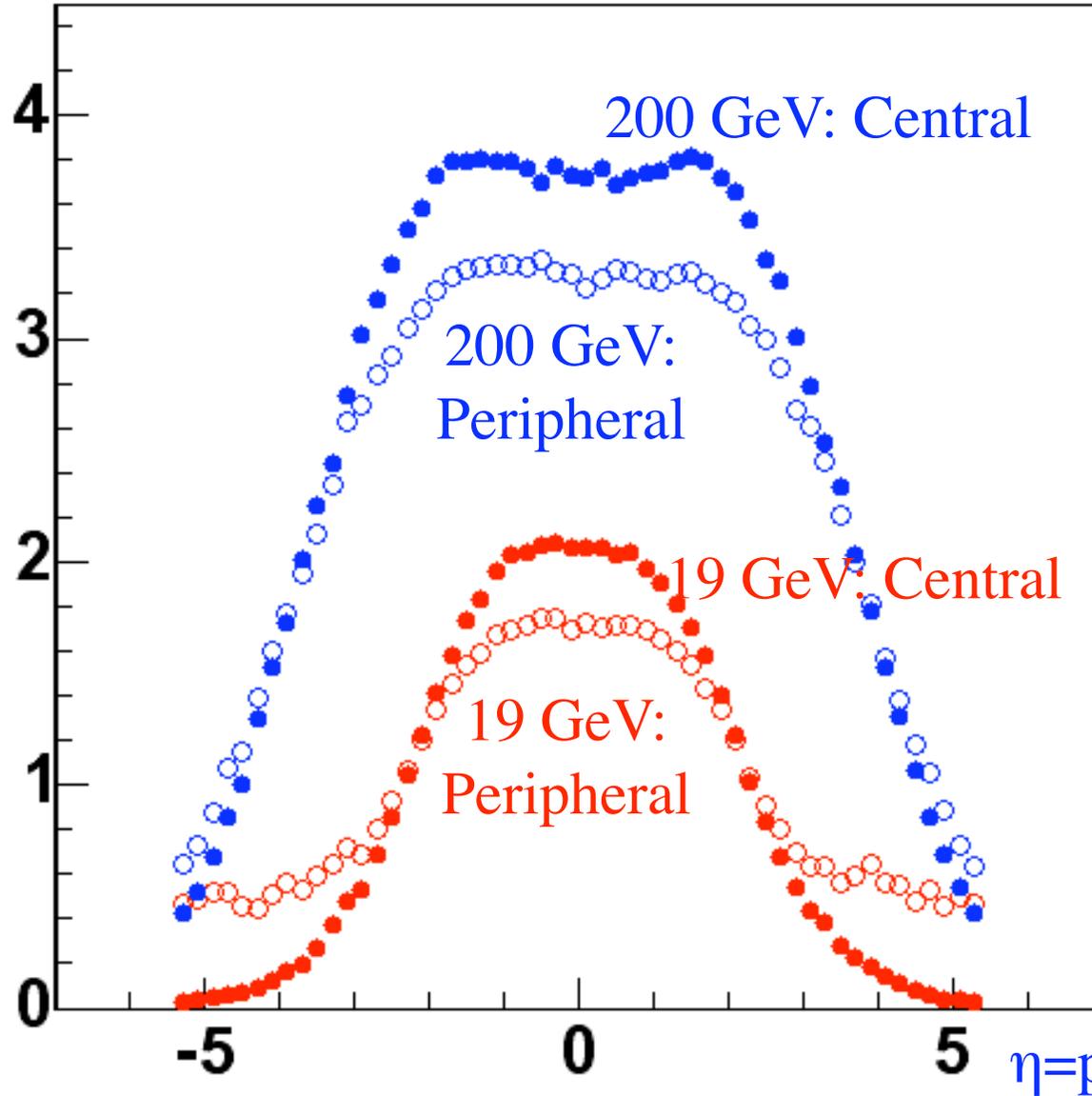
Overall multiplicity: slow growth

$dN/d\eta/ \uparrow$

$N = \#$ particles

$\eta =$ "pseudo"-
rapidity
(no particle ID)

/ by #
"participants"



200 GeV, RHIC
900 particles
/unit η

19 GeV, SPS
600 particles
/unit η

No big increases in multiplicity, as predicted by cascade models.

Rapidity plateau $\pm .5$ (out of ± 5.0) for dN/dy ($y =$ true rapidity)

The Tail Wags the (Dog) Unicorn

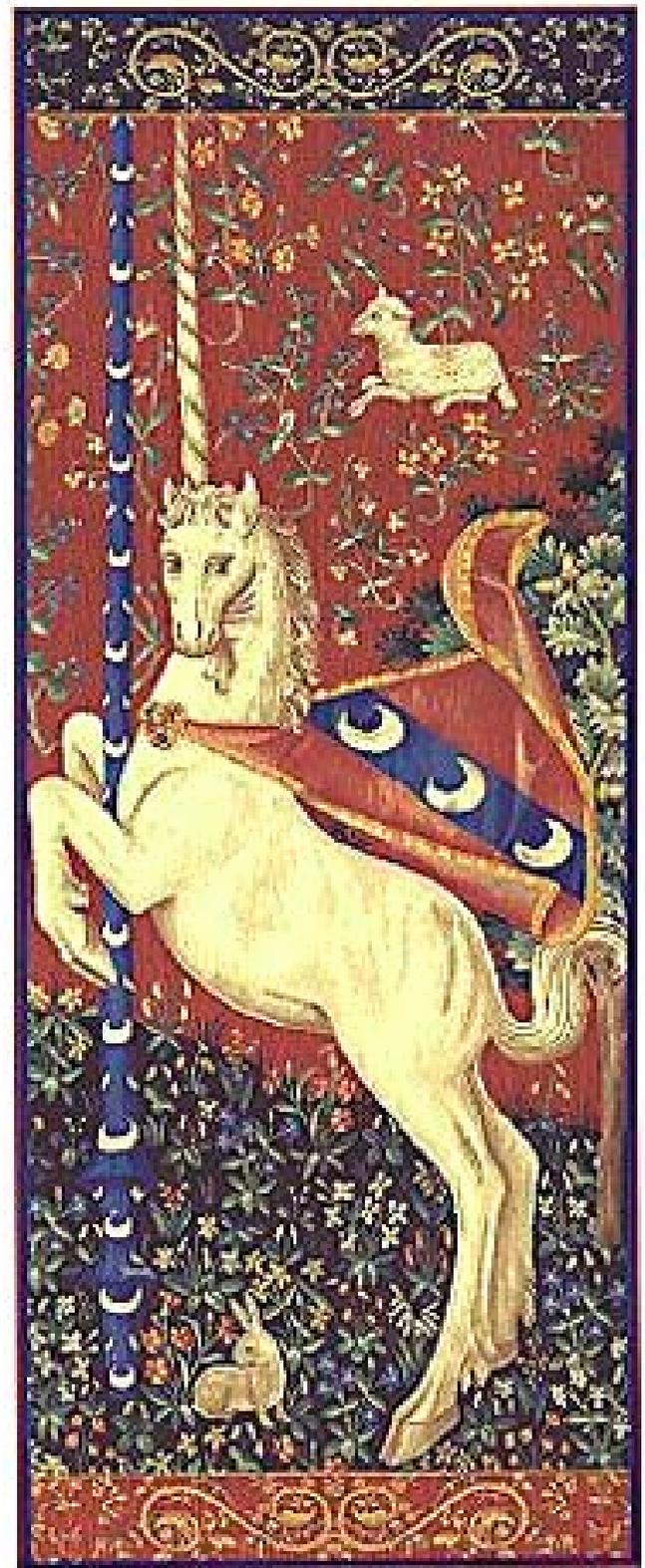
Body of the “Unicorn”:

For $T \sim 200$ MeV,
majority of particles at *small* momenta,
 $p_t < 2$ GeV.

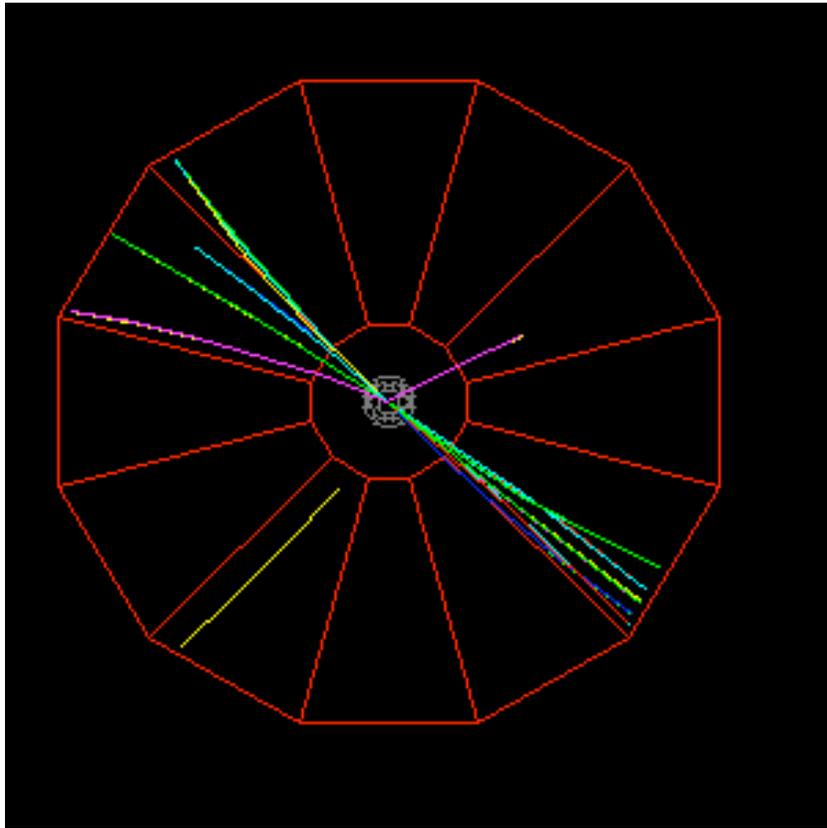
Tail of the “Unicorn”:

Look at particles at *high* momentum,
 $p_t > 2$ GeV, to probe the body.

Concentrate on zero rapidity, 90° to beam.



Jets: “seeing” quarks and gluons in QCD



At high transverse momentum (p_t), instead of hadrons, have \sim quarks & gluons: *jets*.

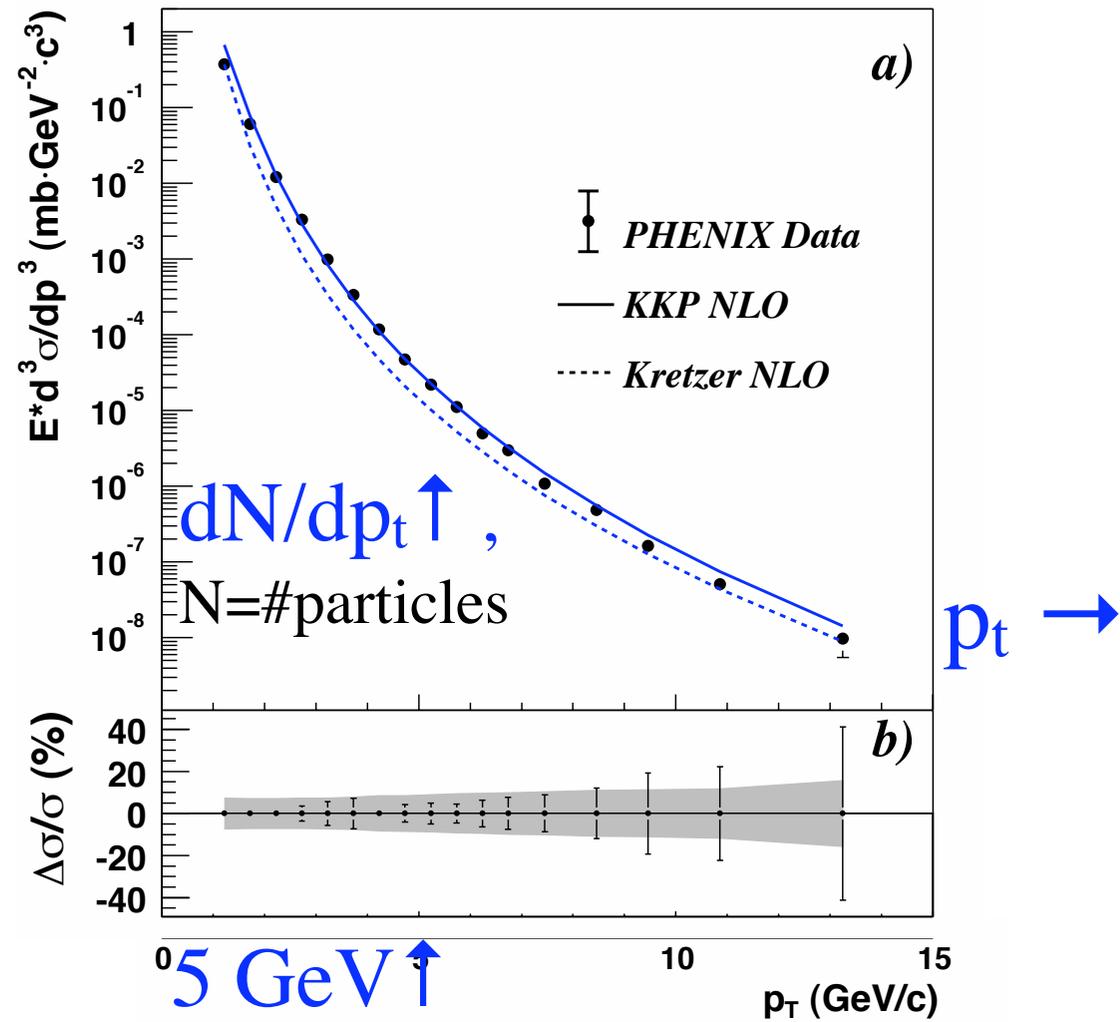
≤ 2 jets from pp collision at RHIC.

By momentum conservation, for each jet, there is a backward jet.

In pp coll.'s, jets can be computed for large p_t , down to $p_t \sim 1$ GeV.

Jets rare: particles at $p_t \sim 2$ GeV $\sim .1$ % of total!

Look at jets in AA



“ R_{AA} ”: *robust* signal of new physics

R_{AA} = for a given p_t , # particles central AA / ($A^{4/3}$ # particles pp)

For π^0 's, $p_t : 2 \rightarrow 20$ GeV, $R_{AA} \sim 0.2$. As if jets emitted *only* from surface!

Due to “energy loss” in thermal medium?

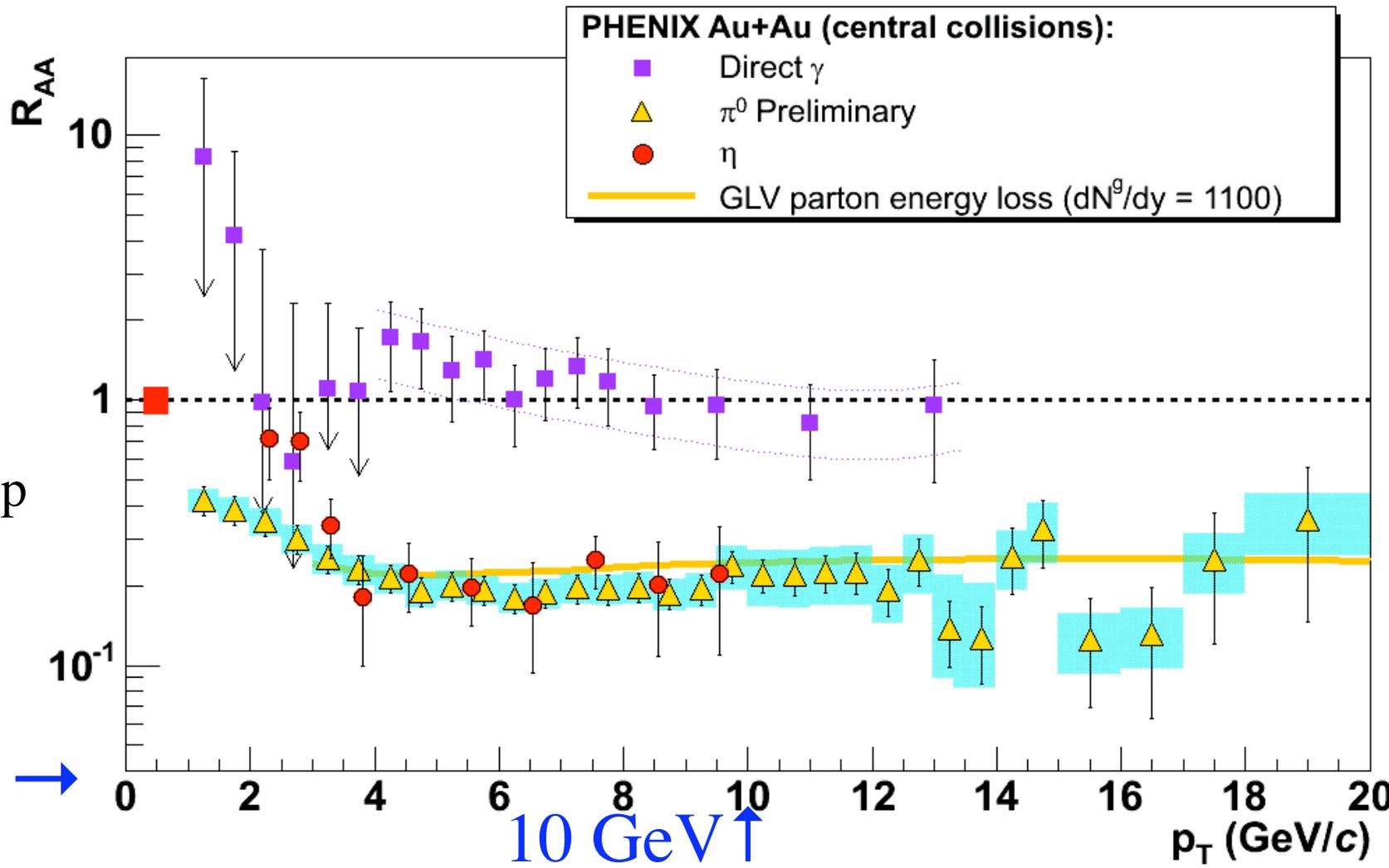
$A^{4/3}$: *experimentally*: for γ 's, $R_{AA} \sim 1.0$ π^0 's “eaten”, γ 's not

R_{AA} : \uparrow

particles
central AA/
particles pp

$A=200 \Rightarrow$

$p_t \rightarrow$

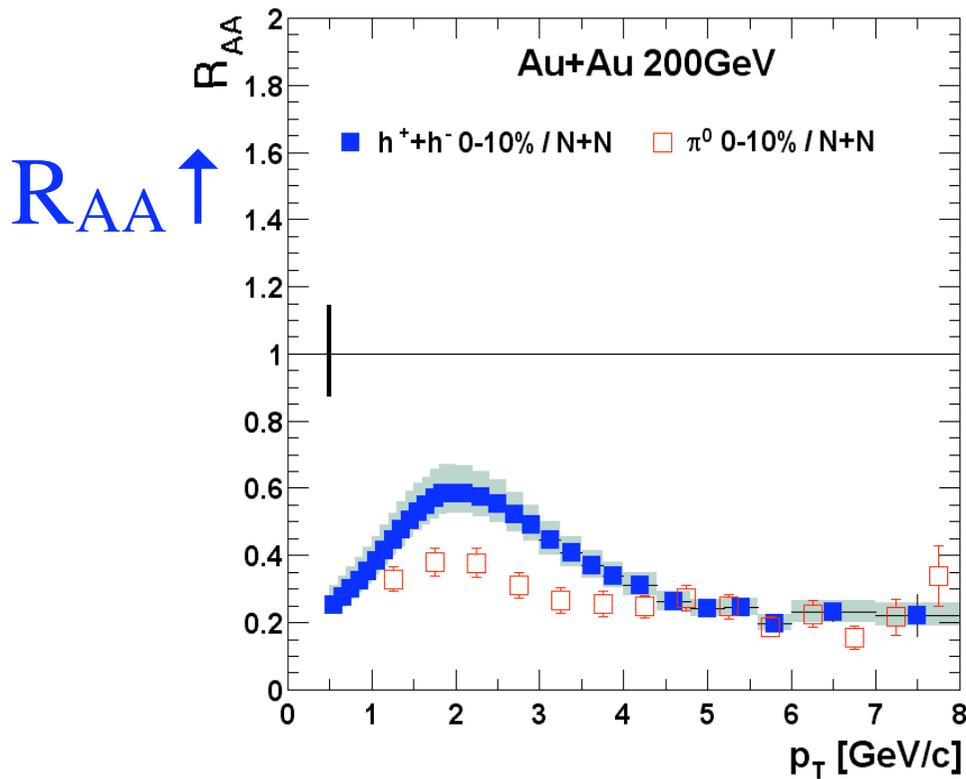


R_{AA} final state effect: *not* in R_{dA}

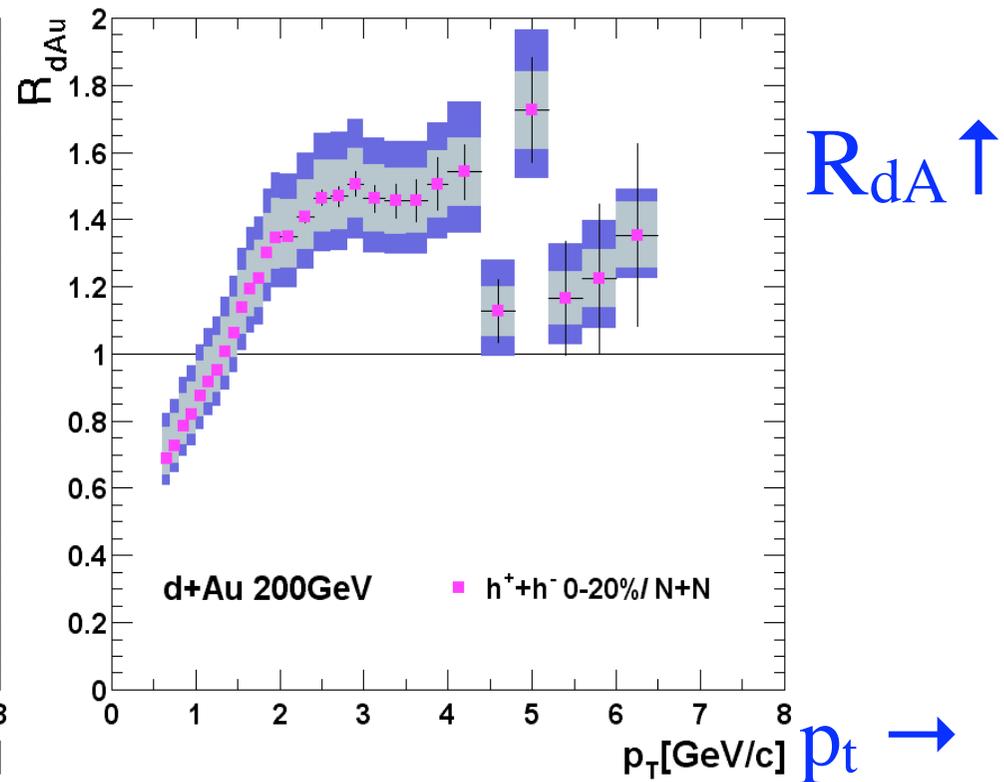
R_{dA} : like R_{AA} , but for dA coll.'s/pp coll.'s. At zero rapidity:

dA: *enhancement*, from initial state “Cronin” effect ($\Rightarrow 1$ @ $p_t > 8$ GeV)

AA: *suppression* \Rightarrow final state effect



Suppression in AA \uparrow
 $R_{AA} \sim 0.4$ @ 3 GeV

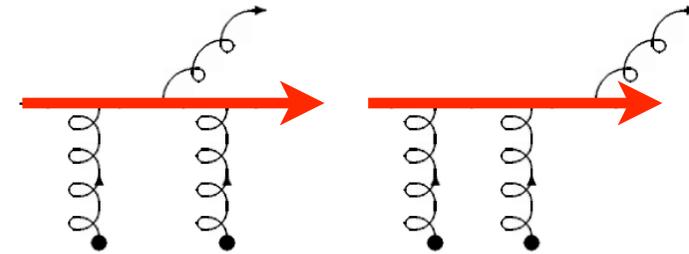


Enhancement in dA \uparrow
 $R_{dA} \sim 1.4$ @ 3 GeV

Explanations of R_{AA}

“Energy loss”: fast particle emits radiation,
scatters off of thermal bath.

Involves Landau-Pomeranchuk-Migdal effect



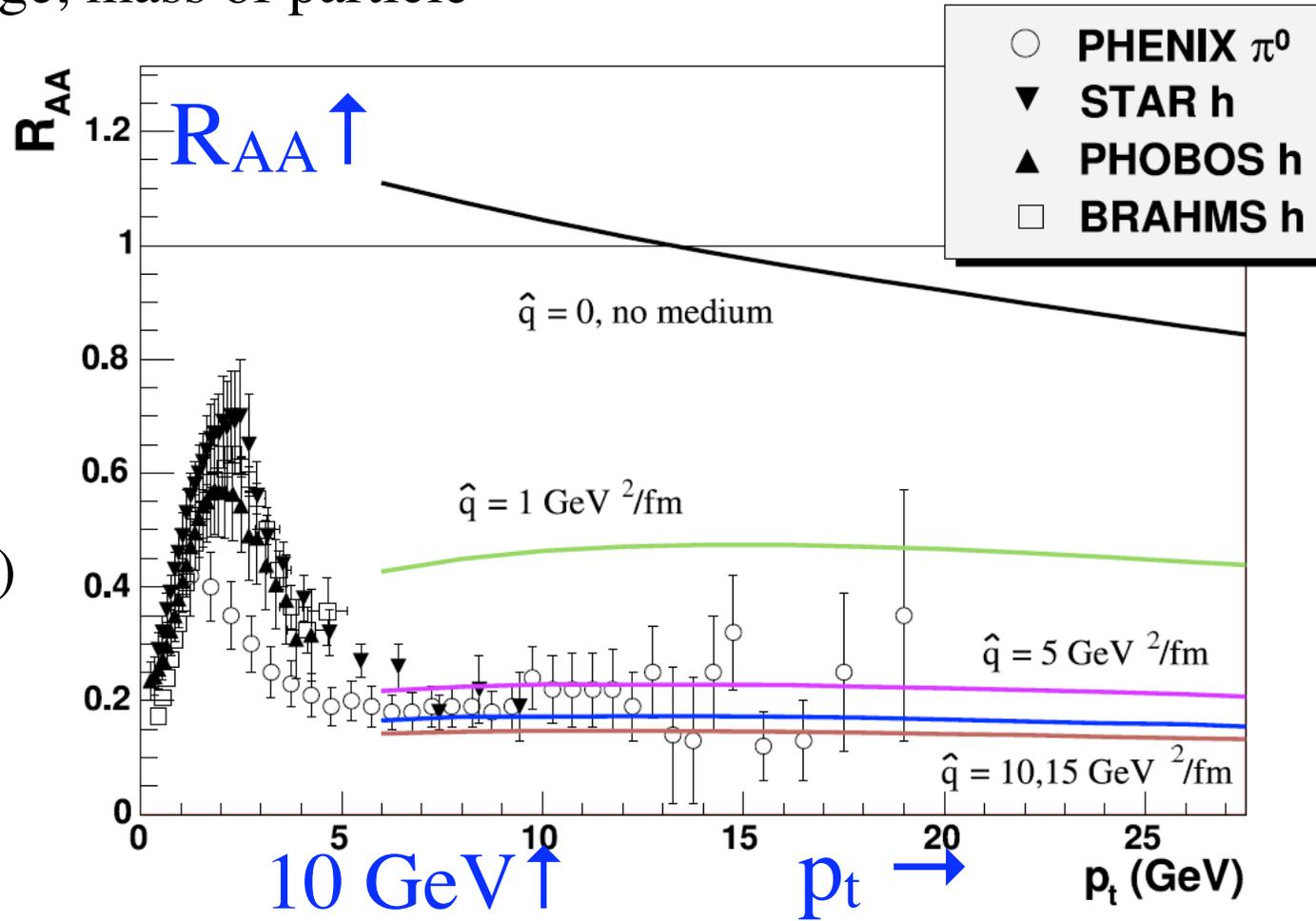
“strong”-QGP: large energy loss.

Details depend upon charge, mass of particle

“semi”-QGP:

propagator \sim loop,
small near T_c .

Details \sim independent of
charge, mass of particle
(assuming thermalization)



R_{AA} for heavy quarks: also suppressed!

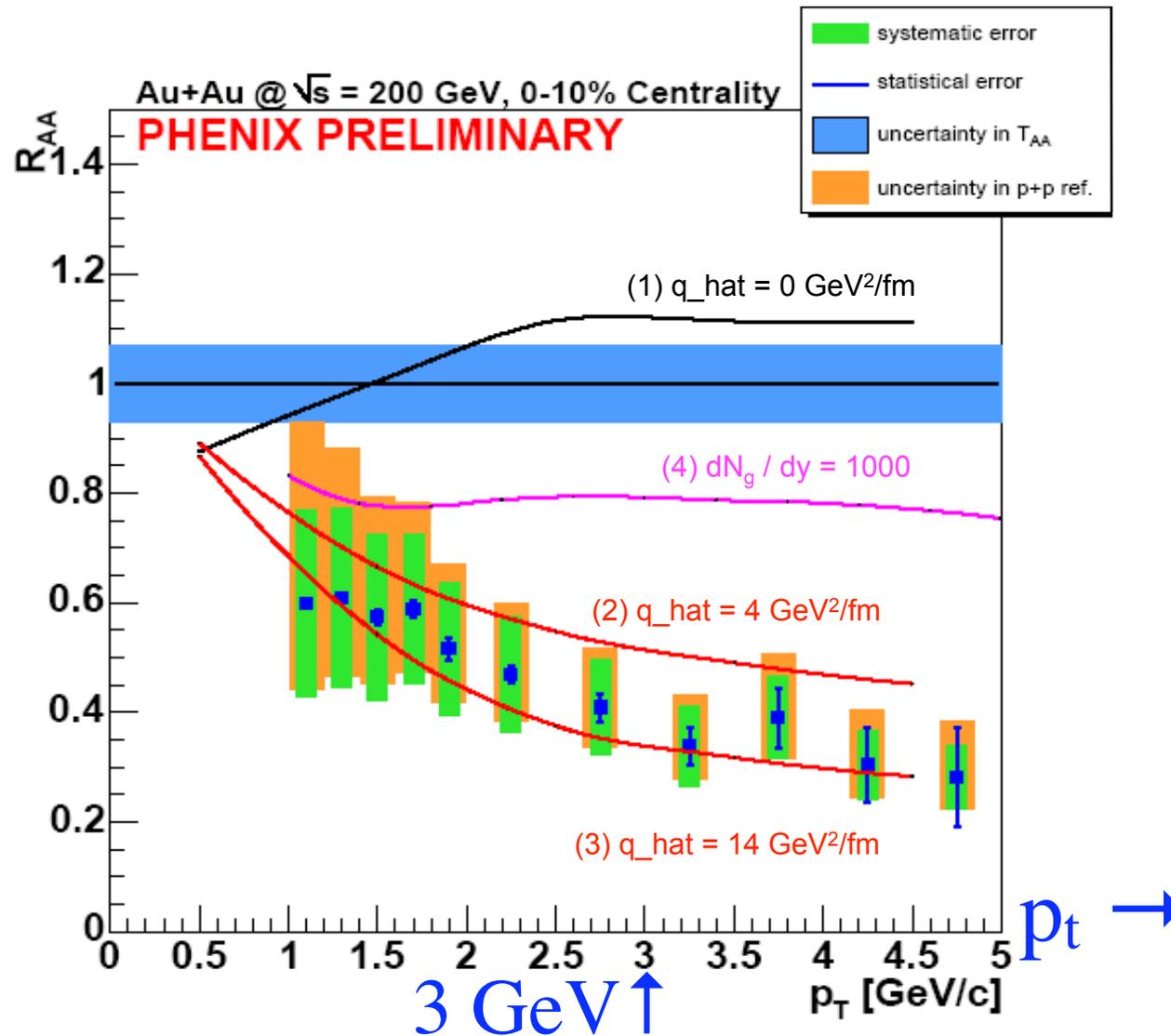
PHENIX: R_{AA} for charm quarks \sim light quarks!

Mass of charm quark $m_{\text{charm}} \sim 1.5 \text{ GeV}$; $T \sim 200 \text{ MeV}$.

Heavy quark less sensitive to medium by $T/m_{\text{charm}} \sim 1/8$. *No sign of that!*

Experimental evidence for “sQGP”: heavy quarks \sim same as light!

$R_{AA} \uparrow$



Central AA collisions “eat” jets!

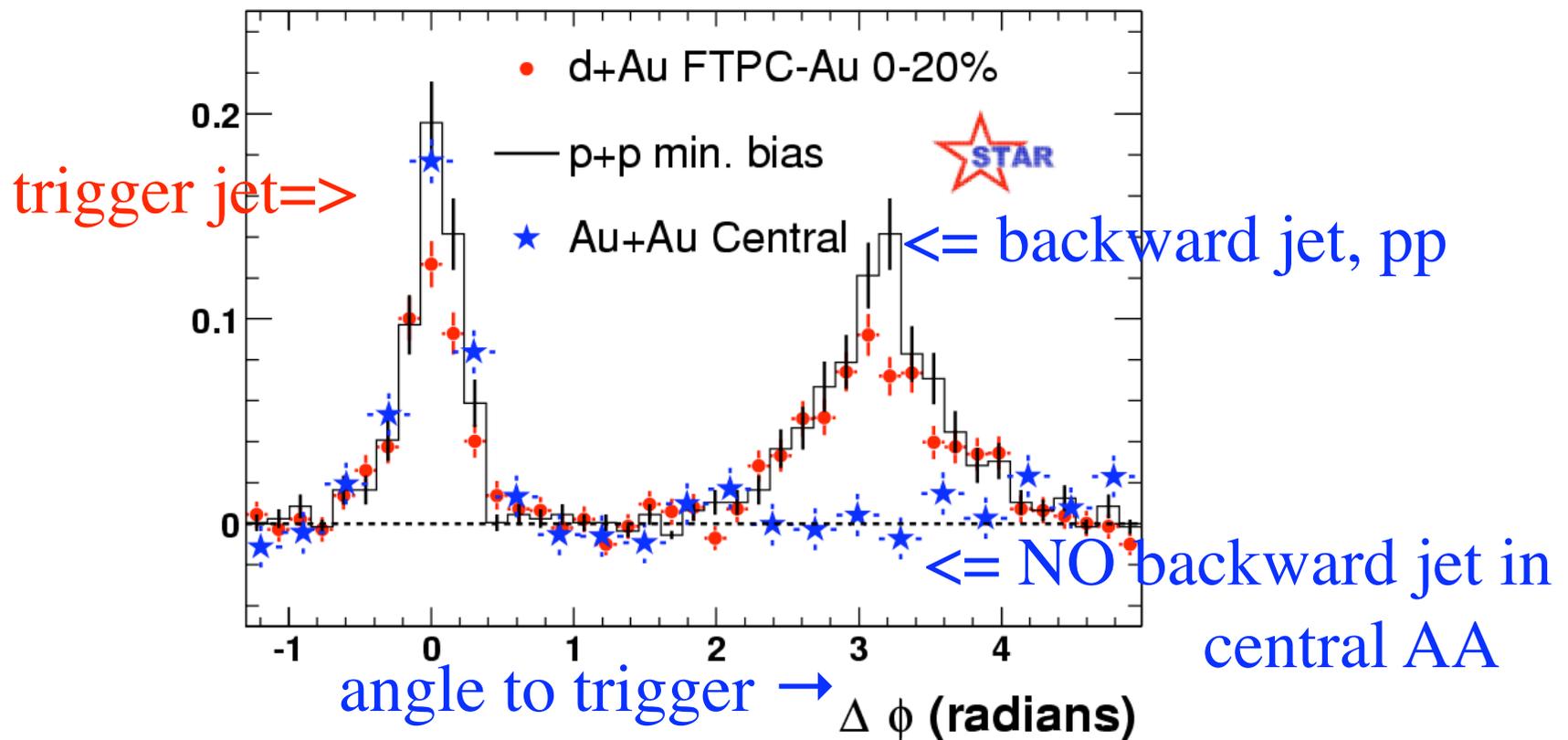
Unlike pp, in central AA, cannot trigger on individual jets. Can:

Trigger on hard “jet”, $p_t: 4 \rightarrow 6$ GeV. Look for backward “jet”, $p_t > 2$ GeV

In pp or dAu collisions, *clearly* see backward jet.

In central Au-Au, away side jet gone: “*stuff*” in central AA “*eats*” jets

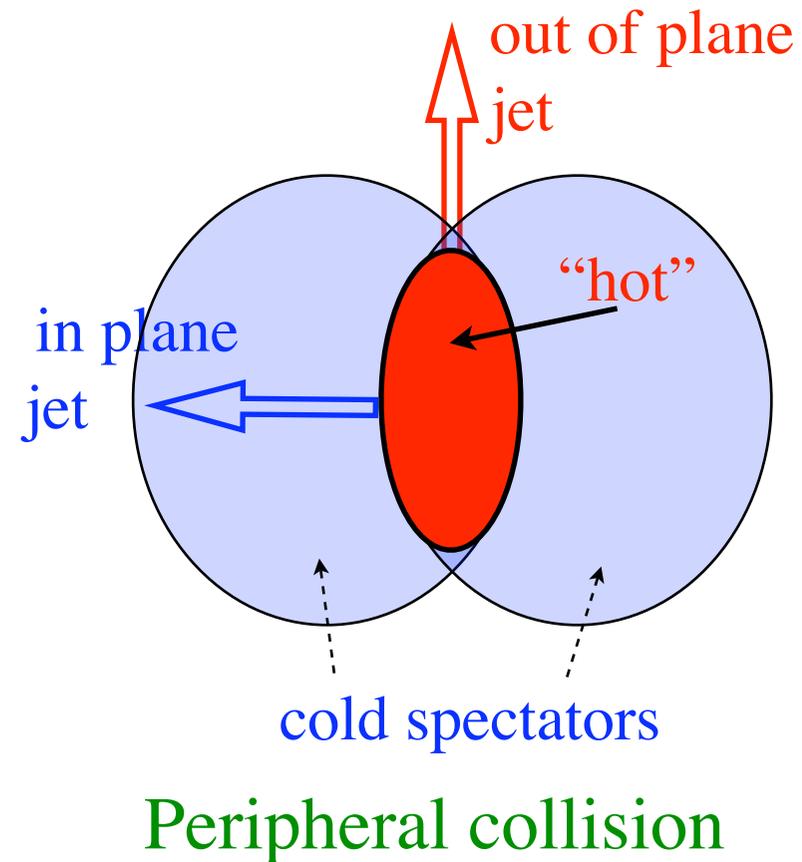
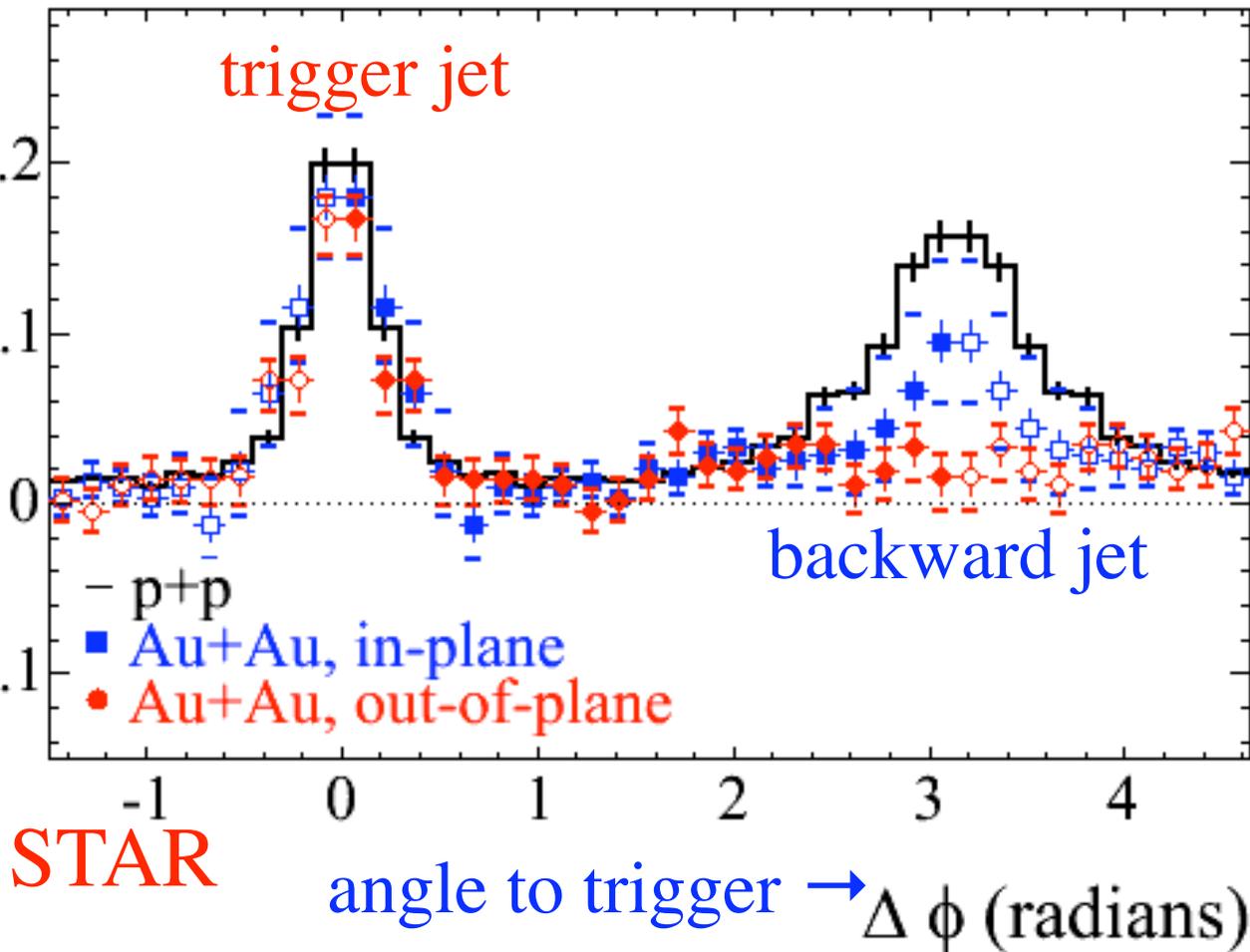
Adams *et al.*, Phys. Rev. Let. 91 (2003)



Peripheral coll.'s: more jets eaten *out* of plane

Peripheral collisions: “hot stuff” forms “almond”. In vs. out of reaction plane
Out: more “hot stuff”. *In*: less hot stuff, more cold nuclear matter

Exp.’y: backward jet more strongly suppressed *out* of plane than *in* plane =>
Geometrical test that central AA “eats” jets preferentially

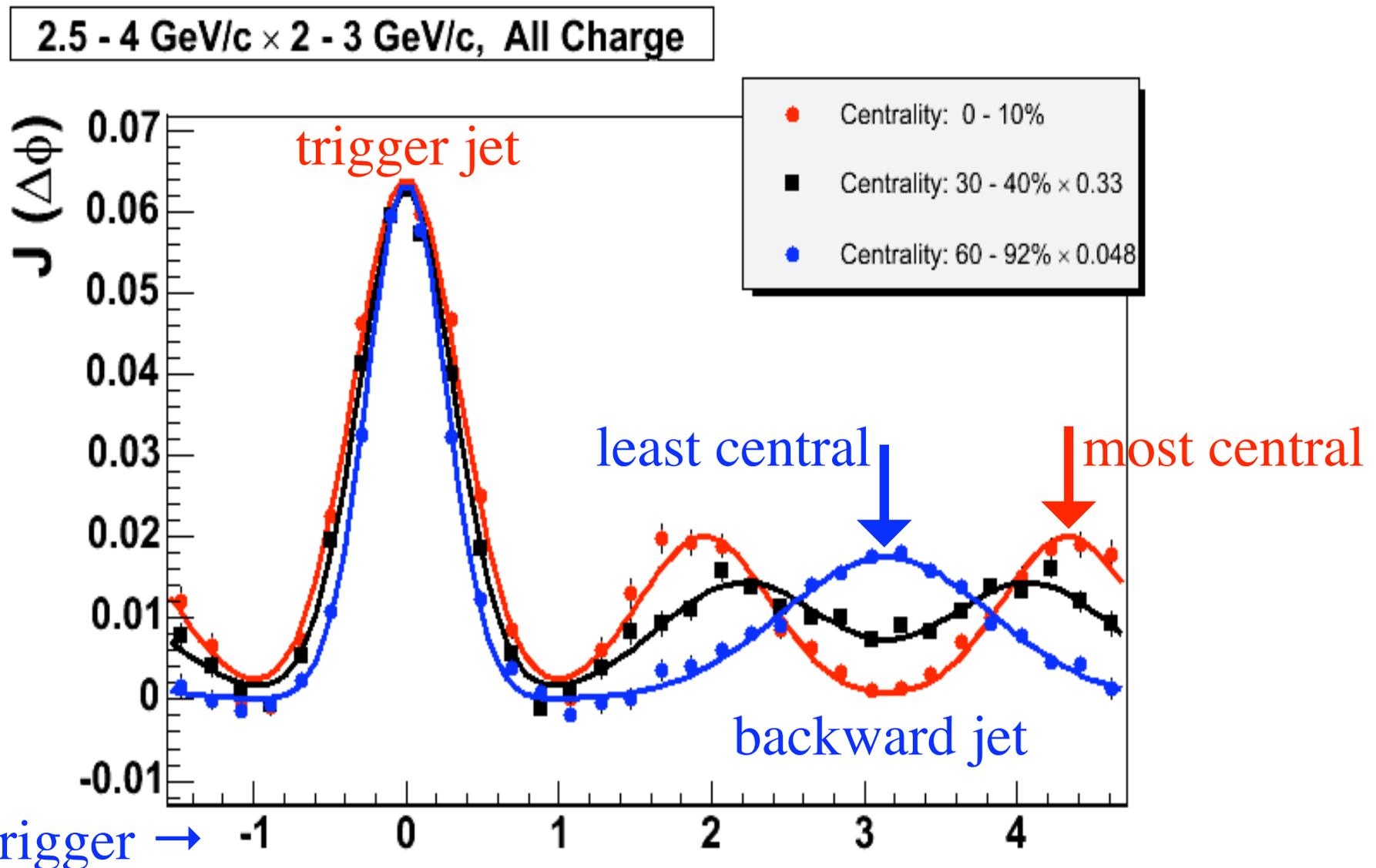


AA collisions modifies jet shapes

PHENIX: shape of away side jet appears to be modified by “stuff”:

Mach cone or Cerenkov radiation in strong-QGP? semi-QGP?

Need to test with 3-particle correlations; appears real.



The Body of the Unicorn: the sQGP

Particles peaked about zero transverse momentum

$T_c \sim 200$ MeV: expect thermal to $p_t \sim 2$ GeV.

Thousands of particles: use hydrodynamics? “Most Perfect Fluid on Earth”

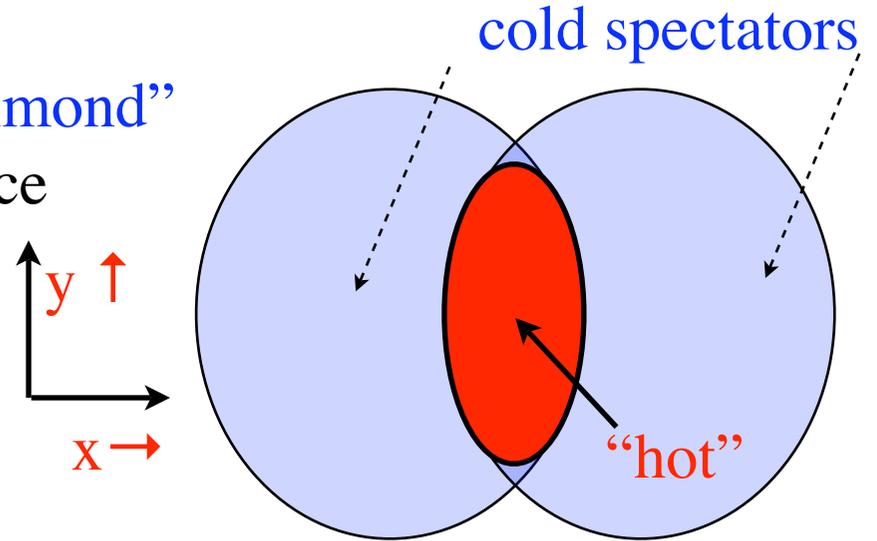


“Elliptic Flow”

For peripheral collisions, overlap region is “almond” in coordinate space, sphere in momentum space

So start with spatial anisotropy,

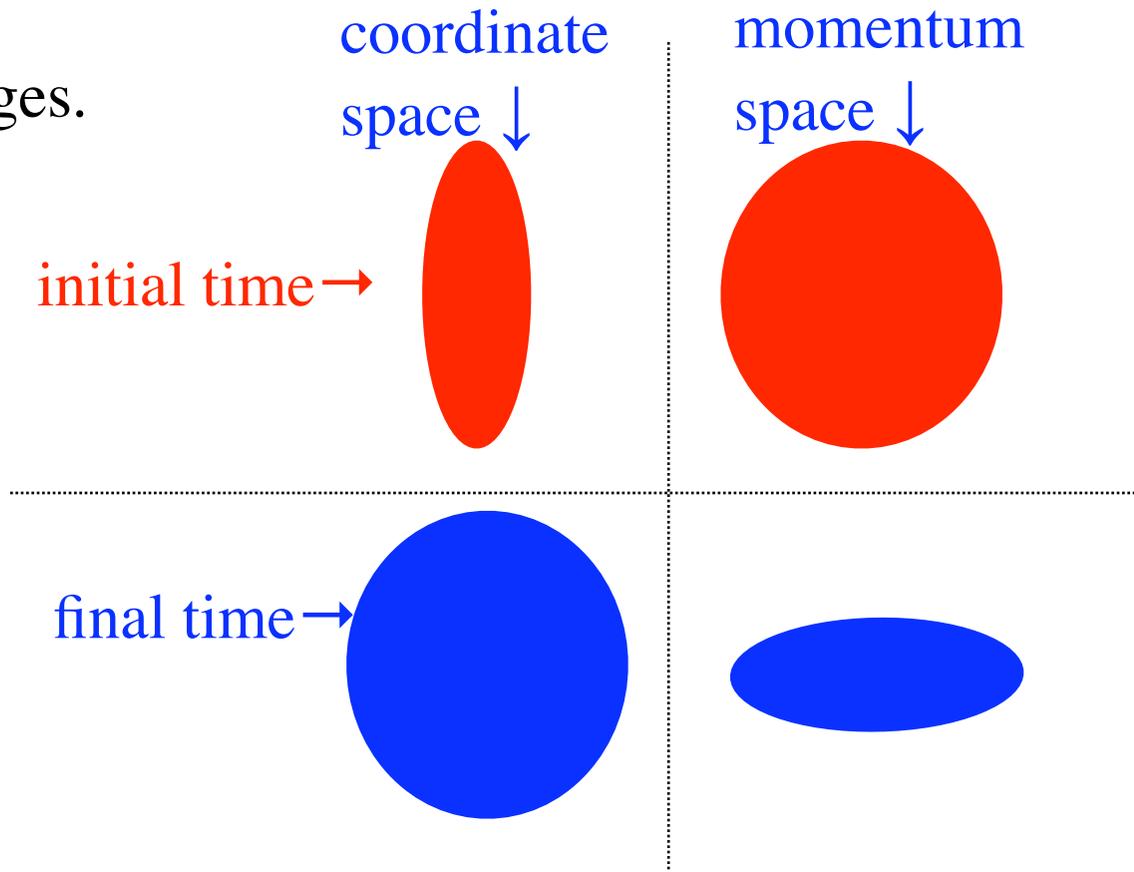
$$\epsilon = \frac{\langle y^2 - x^2 \rangle}{\langle x^2 + y^2 \rangle}$$



If particles free stream, nothing changes.

If collective effects present, end up with sphere in coordinate space, almond in momentum space: “elliptic flow”

$$v_2 = \frac{\langle p_y^2 - p_x^2 \rangle}{\langle p_x^2 + p_y^2 \rangle}$$



“Most Perfect Fluid on Earth”

Large # particles: try *ideal* hydrodynamics:

1. *Short* initial time (tune)
2. MIT Bag Equation of State
3. *Small* viscosity in QGP phase.
4. Hadronic “afterburner”: pions, K’s, p’s.

Good fit to π ’s, K’s, p’s.... for *both* single particle spectra *and* v_2 .

Need *small* viscosity.

Viscosity \sim 1/cross section:

small viscosity \Rightarrow

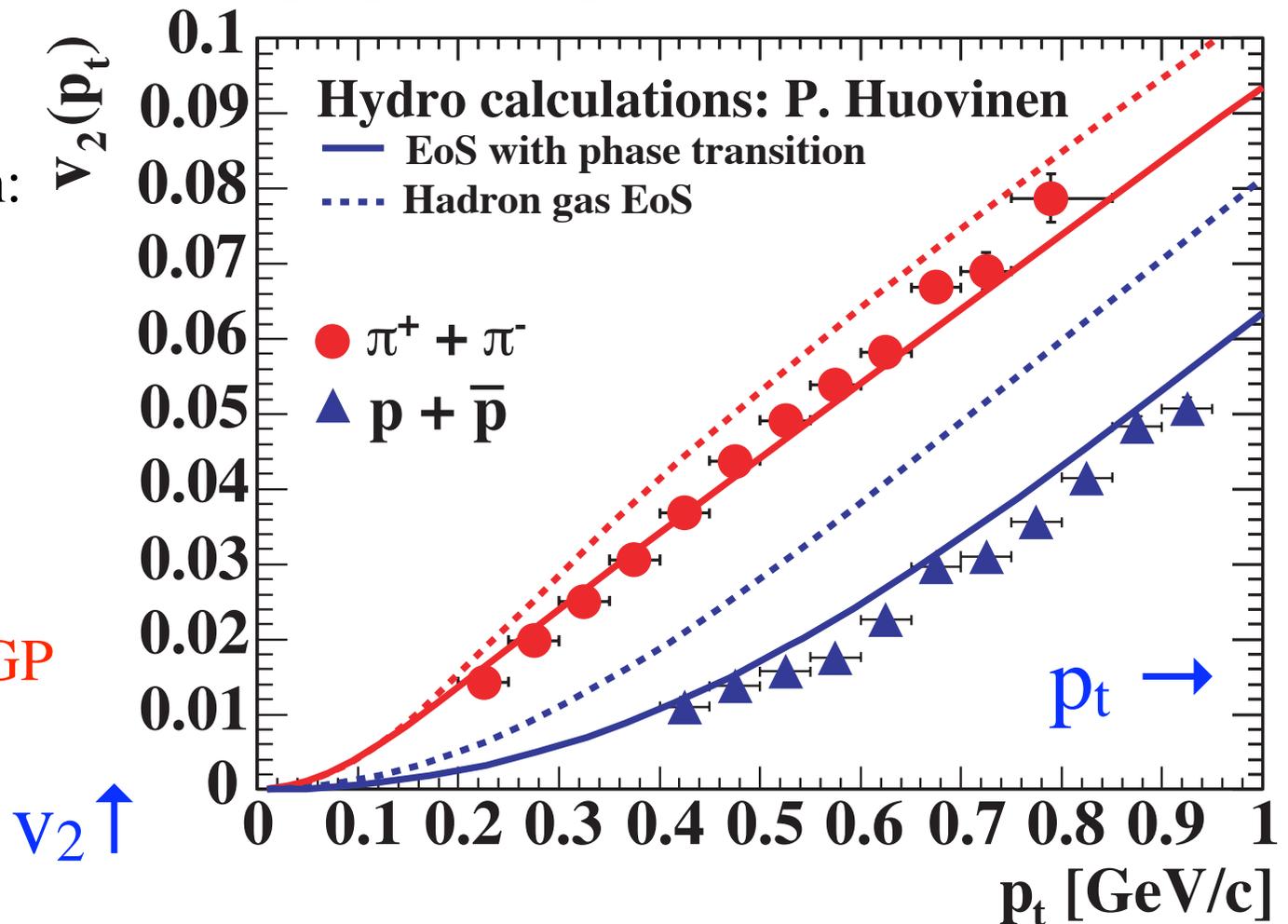
strong (coupling) QGP?

Gyulassy, Shuryak, Heinz,
Mrowczynski, Thoma...

Small Polyakov loops \Rightarrow

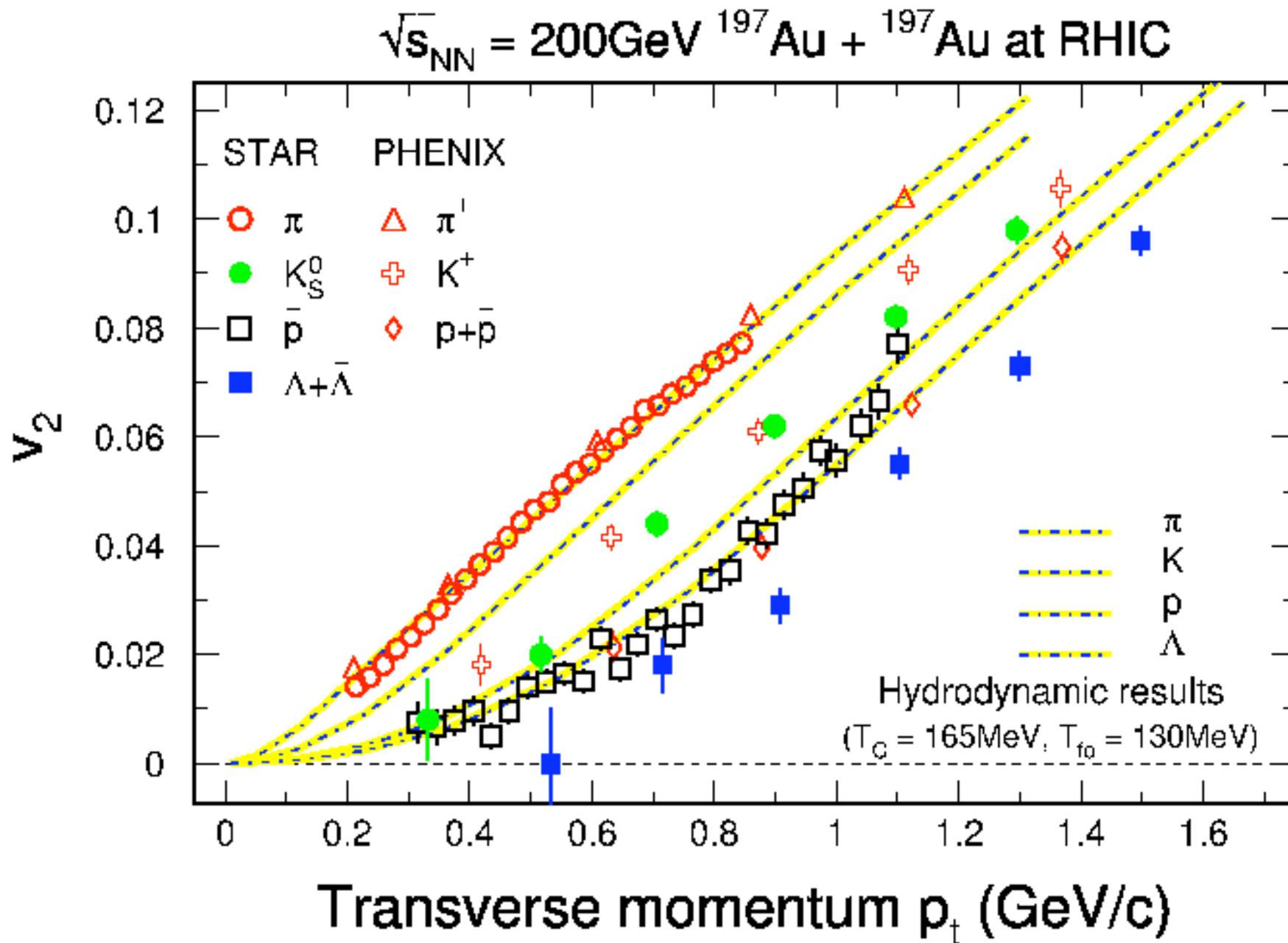
small viscosity in semi-QGP

Y. Hidaka & RDP ‘08



Strange Hydrodynamics, I

Hydrodynamic fits less dramatic for strange particles



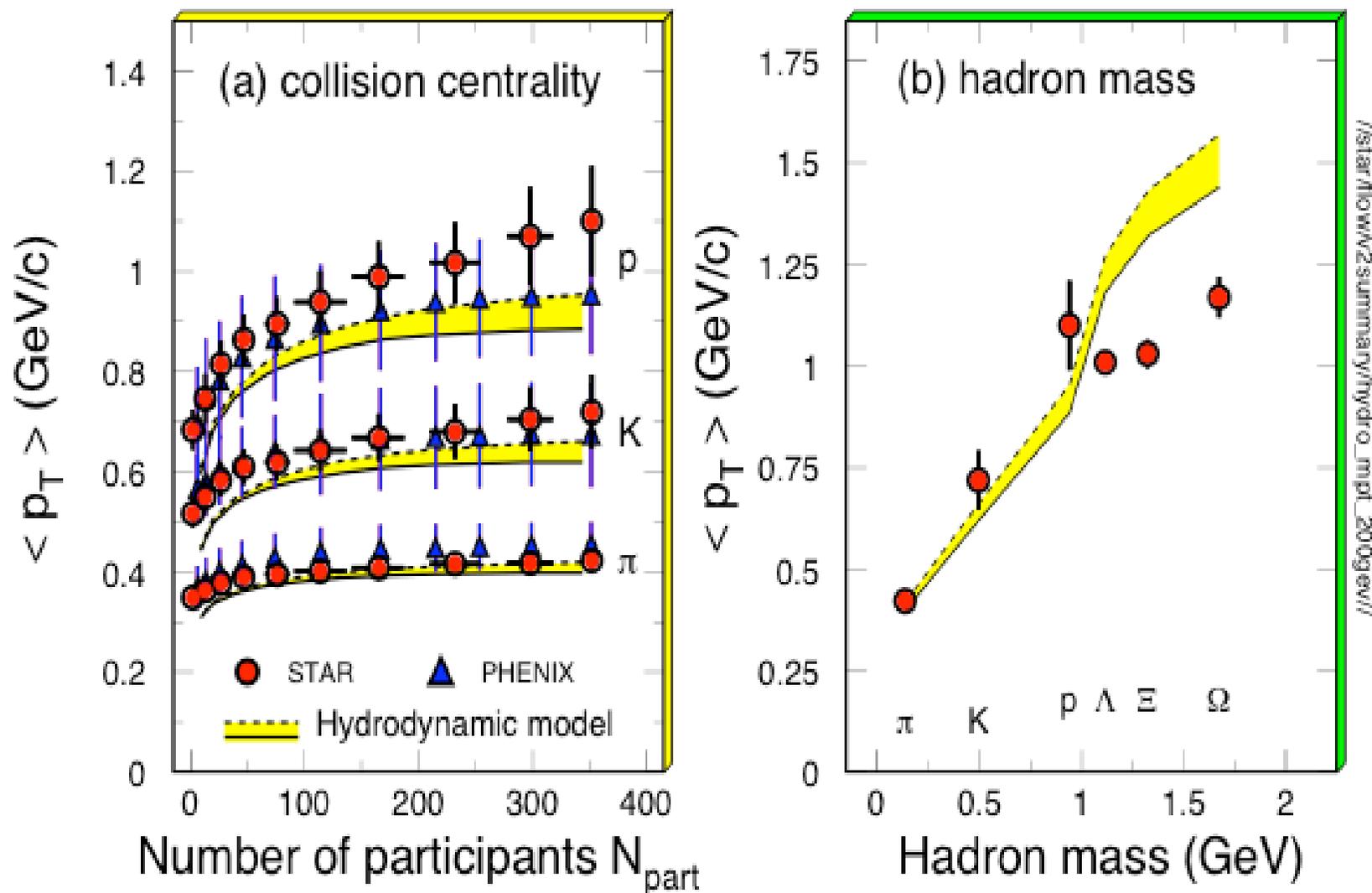
Strange Hydrodynamics, II

Strange particles? Hydro: $\langle p_t \rangle \sim \text{mass} \times (\text{velocity of medium})$

OK for π , K, p. But $\langle p_t \rangle \sim 1 \text{ GeV}$ for p , ϕ , Λ , Ξ , Ω !

Strange particles *don't* seem to flow with the medium. Why *constant* $\langle p_t \rangle$?

Au + Au Collisions at $\sqrt{s_{NN}} = 200 \text{ GeV}$

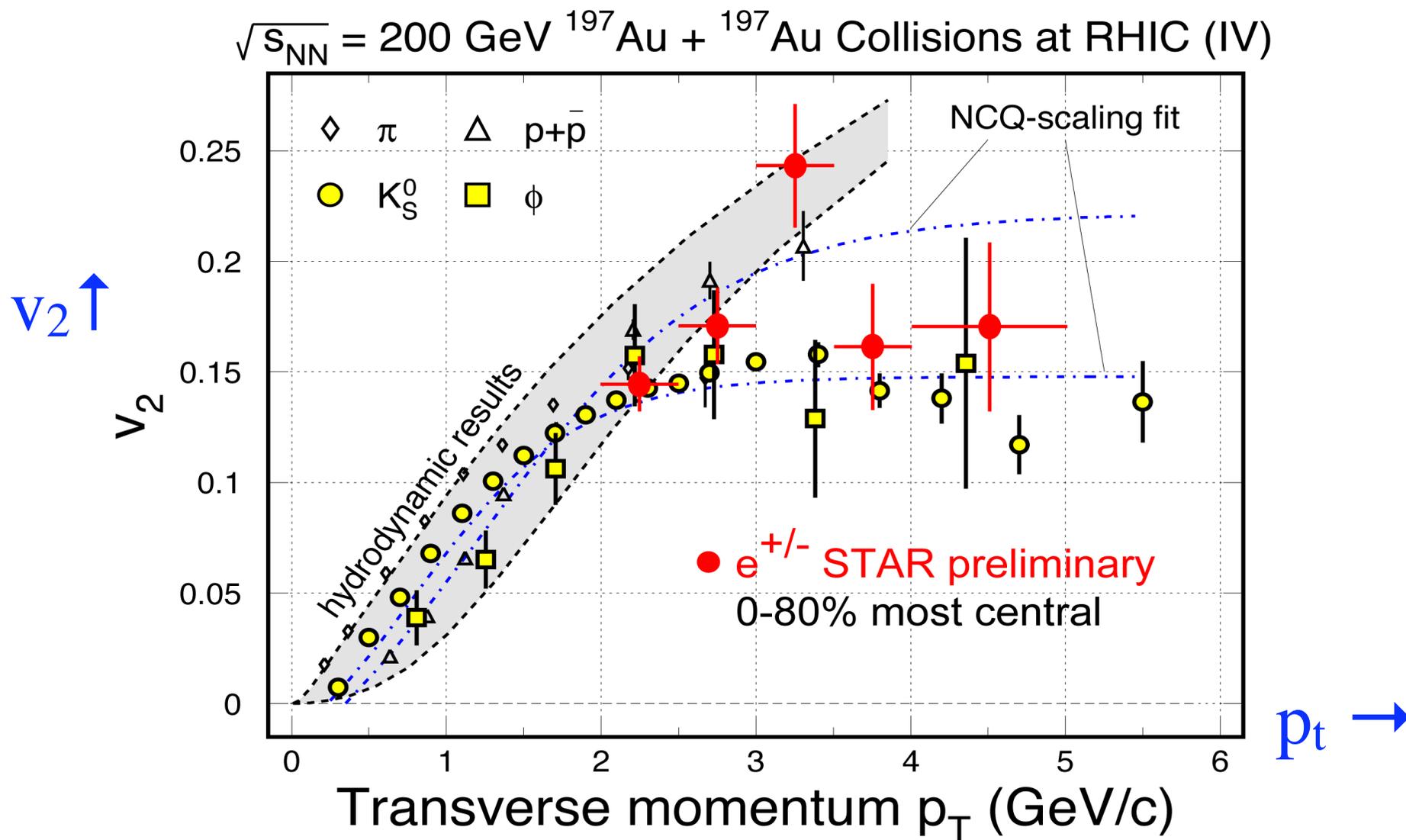


Even charm quarks flow

Look at charm quarks through single electrons.

See *large* elliptic flow: *no* suppression due to large mass.

Experimental definition of “sQGP”: heavy quarks affected *~same* as light quarks!

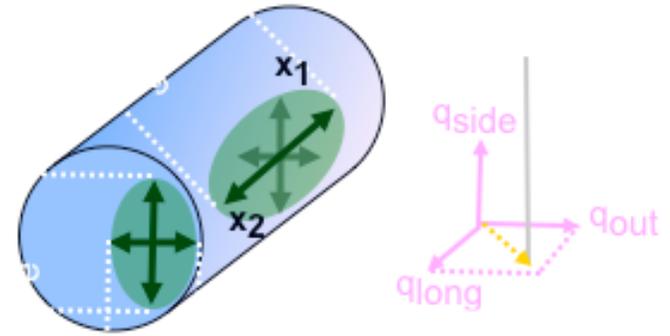


HBT radii: collisions “*explosive*”

Hanbury-Brown-Twiss: two-particle correlations of identical particles
 = *sizes at freezeout*. *Three* directions:
 along beam R_{long} , along line of sight R_{out} , perpendicular R_{side} .

$$C(p_1, p_2) = N(p_1, p_2) / (N(p_1)N(p_2))$$

$$= 1 + \lambda \exp(-R_{HBT}^2 (p_1 - p_2)^2)$$



Hydro.: $R_{\text{out}}/R_{\text{side}} > 1$, *increases* with p_t
 (“burning log”)

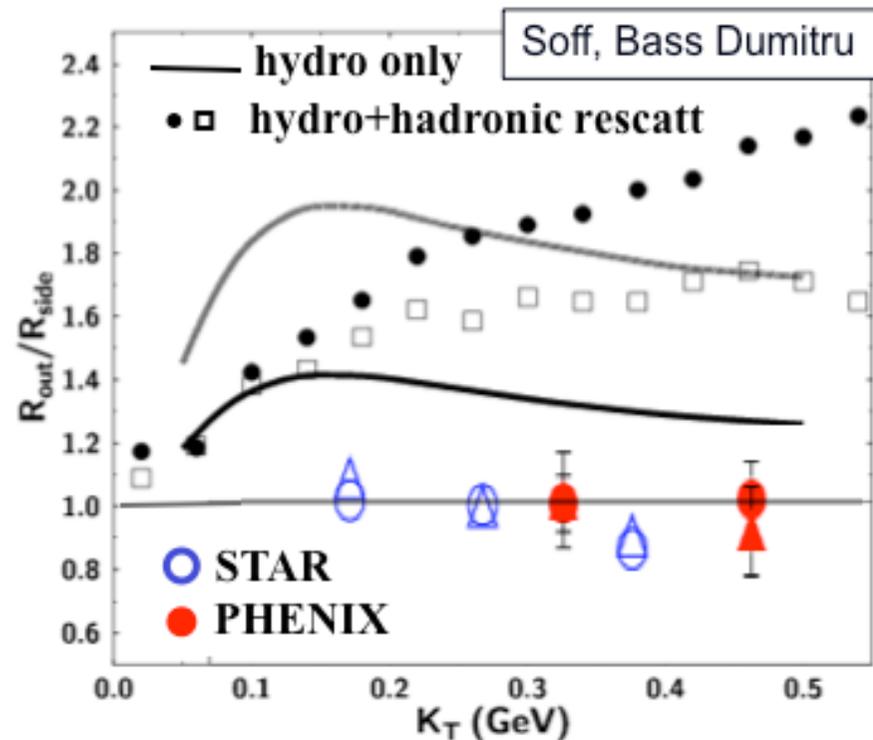
Exp.: $R_{\text{out}}/R_{\text{side}} \sim 1.0$, *flat* with p_t

Hydro. fails - *badly* - for HBT radii.

No big times from strong 1st order trans.!

HBT “*explosive*”: blast wave works:

Space-time history shell with
 lifetime $\sim 8\text{-}9$ fm/c, emission ~ 2 fm/c



HBT: p_t dependence same in pp, dA, AA!

sQGP: from RHIC to LHC

At RHIC, central AA \neq A (pp) collisions

“Tail wags the Unicorn”

Clearest signal from “high” pt:

R_{AA} , jet suppression...

Body of the Unicorn: “sQGP”

Assume RHIC near T_c , LHC well above.

“s” = strong. AdS/CFT rules!

No surprises from RHIC to LHC.

“s” = semi. Partial ionization of color.

Use lattice to construct effective theory

Qualitatively new behavior at LHC.

Nothing better than a good dog fight...





"A possible eureka."

From the SPS: NA60 and a “thermal” ρ

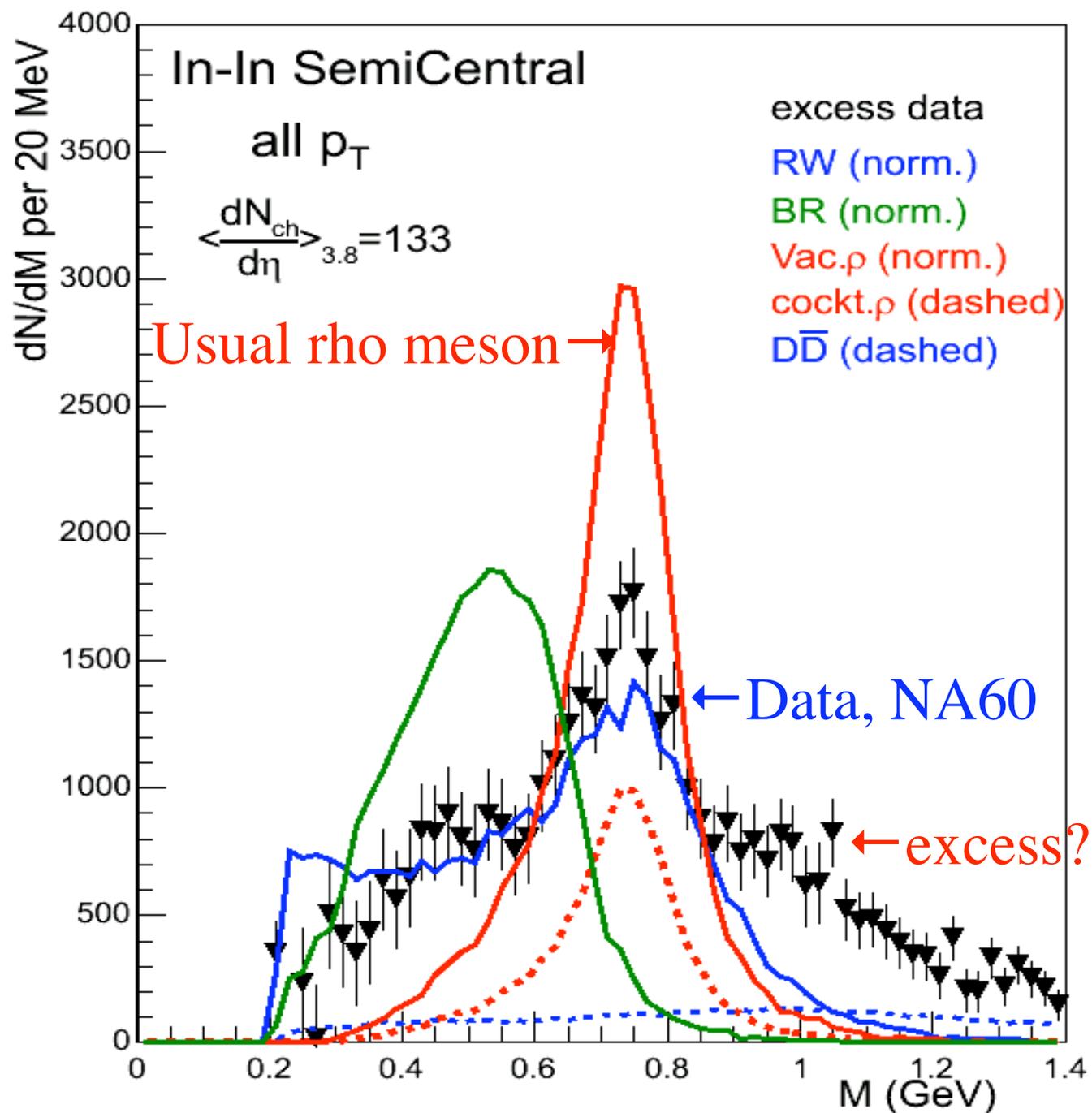
SPS = fixed target.
Kinematics more awkward, but can generate many more events.

Example: electromagnetic signals, such as e^+e^- pairs.

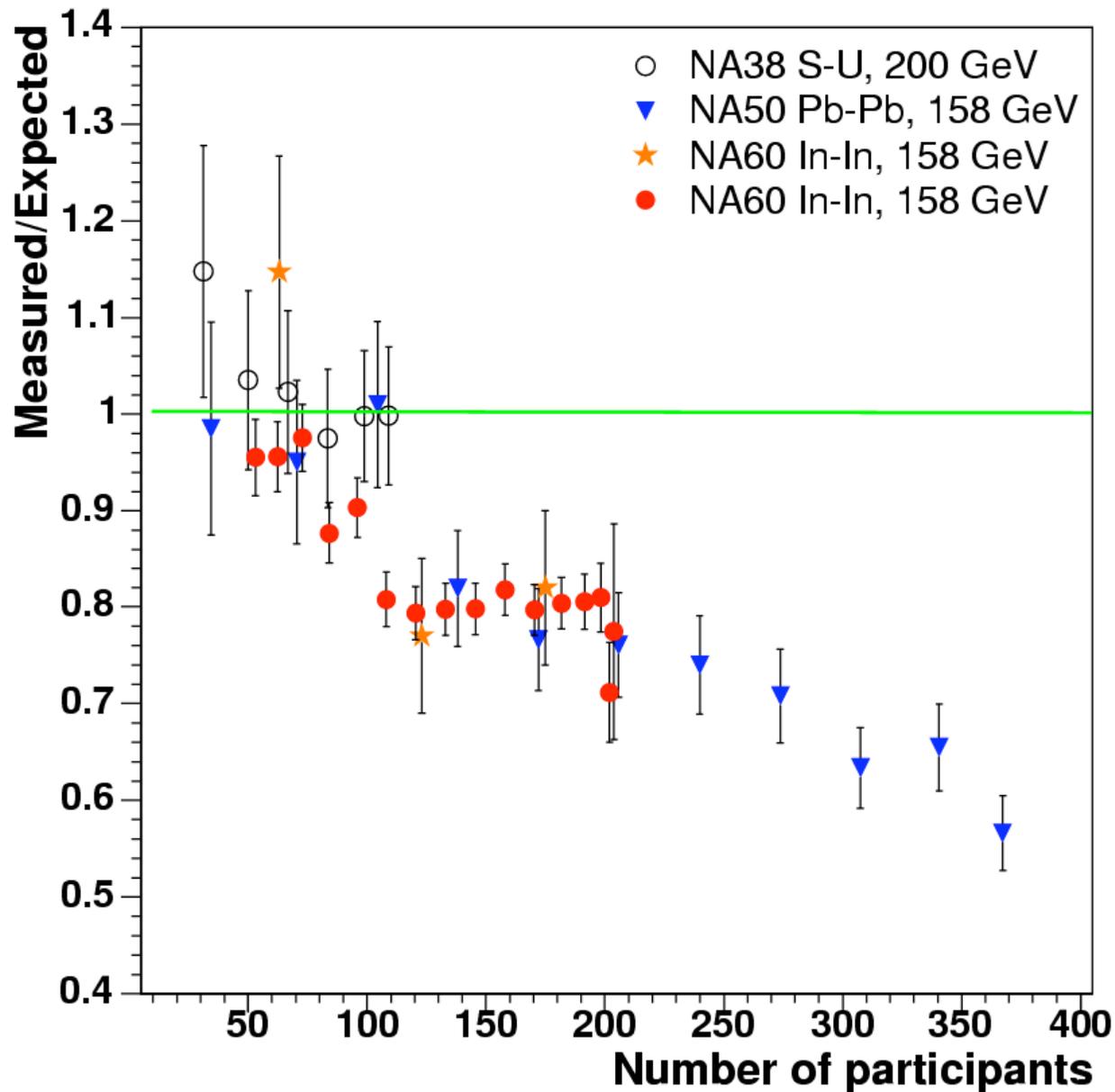
Look at “ ρ ” meson, mass ~ 770 MeV. Decays directly to e^+e^- .

Find “thermal” ρ : no shift in mass, thermal broadening.

Interesting excess above the ρ ?



SPS: NA50, NA60 J/Psi suppression



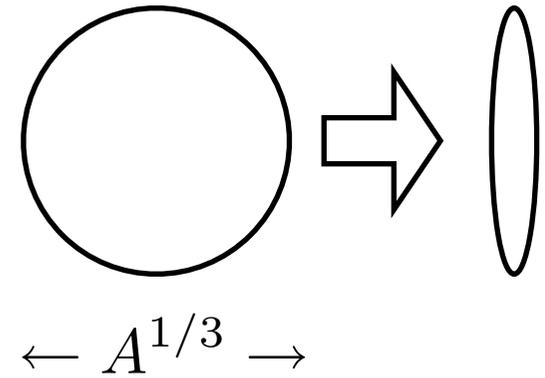
consistent, # participants good variable

Why do AA? “Saturation” as a Lorentz Boost

Incident nucleus Lorentz contracted at high energy

McLerran & Venugopalan: color charge bigger by $A^{1/3}$

$A \rightarrow \infty$: semi-classical methods, Color Glass



\Rightarrow Logarithmic growth in multiplicity:

$$\frac{dN}{dy} \sim \frac{1}{g^2(\sqrt{s}/A)} \sim \log(\sqrt{s}/A)$$

Expect at least same rise in $\langle p_t \rangle$.

Color Glass: “saturation momentum” function of energy, rapidity...

CG describes initial state. Final state?

Multiplicity, energy, & Color Glass

For example: compare central AuAu, 130 & 200 GeV:

All exp.'s: multiplicity increases by $\sim 14\% \pm 2\%$.

Kharzeev, Levin, Nardi...: Color Glass gives good $dN/d\eta$ with centrality....

But: STAR (alone) claims ratio of $\langle p_t \rangle = 1.02 \pm 2\%$: \sim SAME!

Color Glass, hydrodynamics... all predict $\langle p_t \rangle$ increases with multiplicity!

From initial to final: “parton hadron duality”: STAR, pp & AuAu, 200 GeV
one gluon \Rightarrow one pion

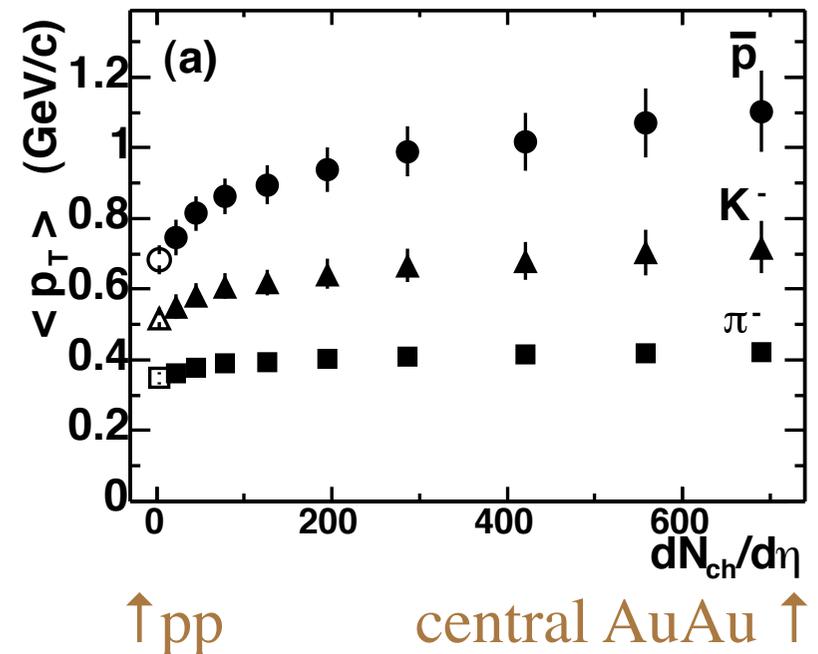
But from pp to central AuAu:

$\langle p_t \rangle \sim$ same for pions

$\langle p_t \rangle$ increases for K's, even more for p's!

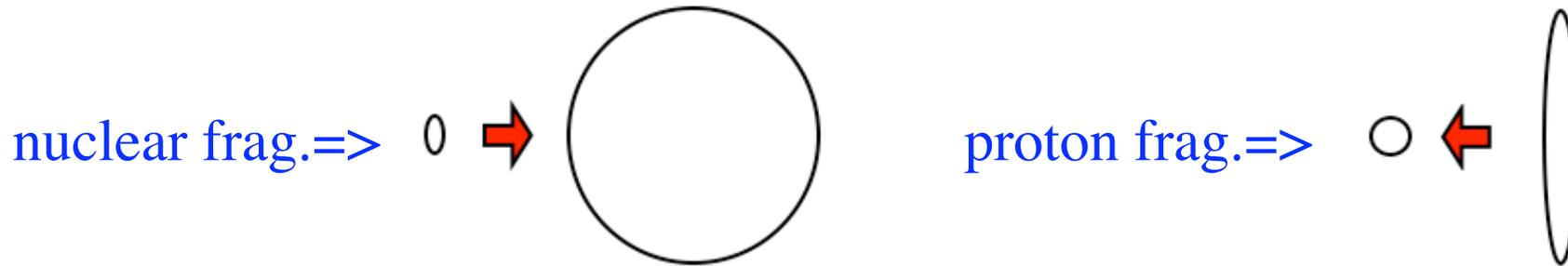
\Rightarrow CG not final state

Hydro: big “boost” velocity of medium.



Color Glass suppression: in dA, by the deuteron

Fragmentation region \sim rest frame. Incident projectile Lorentz contracted:



Nuclear fragmentation region: proton contracted. Study final state effects

Proton fragmentation region: study initial state effects

BRAHMS in dA:

enhancement @ zero rapidity
suppression @ proton frag. region.

Supports color glass initial state.

Need to study all rapidities.

R_{dA} :

