



A Unicorn at RHIC: the (s)QGP

At *very* high temperatures, expect hadrons (QCD) behaves qualitatively different:
new type of matter, Quark Gluon Plasma (QGP).

To see experimentally, collide heavy ions (AA) at high energies:

Relativistic Heavy Ion Collider, RHIC at Brookhaven; SPS & LHC at CERN.

Today:

Theory of the QGP, especially numerical “experiments” on a lattice

RHIC: Suppression of jets in AA collisions. *Robust* signal of new physics.

Non central collisions exhibit strong “flow”

Even heavy quarks are suppressed, flow!

Not a perturbative QGP, instead “sQGP” (s?)

“*Most Perfect Fluid on Earth*”

AdS/CFT duality: hot “Super QCD” at infinite coupling.

Theorists all agog about a *wealth* of experimental data.

Myths can come true, but maybe not the ones you expect...



QCD: Quantum ChromoDynamics

Gluons: spin-1, three colors, local SU(3) symmetry. Classically, *no* mass scale.
 One dimensionless coupling, α_s , which “runs” with momentum.

Theory *uniquely* specified by value of α_s at *one* (finite) momentum.

Asymptotically Free (AF): α_s *vanishes* at *infinite* momenta

$$\alpha_s(p) \sim \# / \log(p) , p \rightarrow \infty$$

AF => *ideal gas* at *infinite* temperature.

Pressure: $p(T \rightarrow \infty) \sim \# T^4$, $\# = 8 \times$ photon!

“Perturbative”: > 1 GeV (= 1000 MeV)

Powers of $\alpha_s \sim 1/\log(p)$...

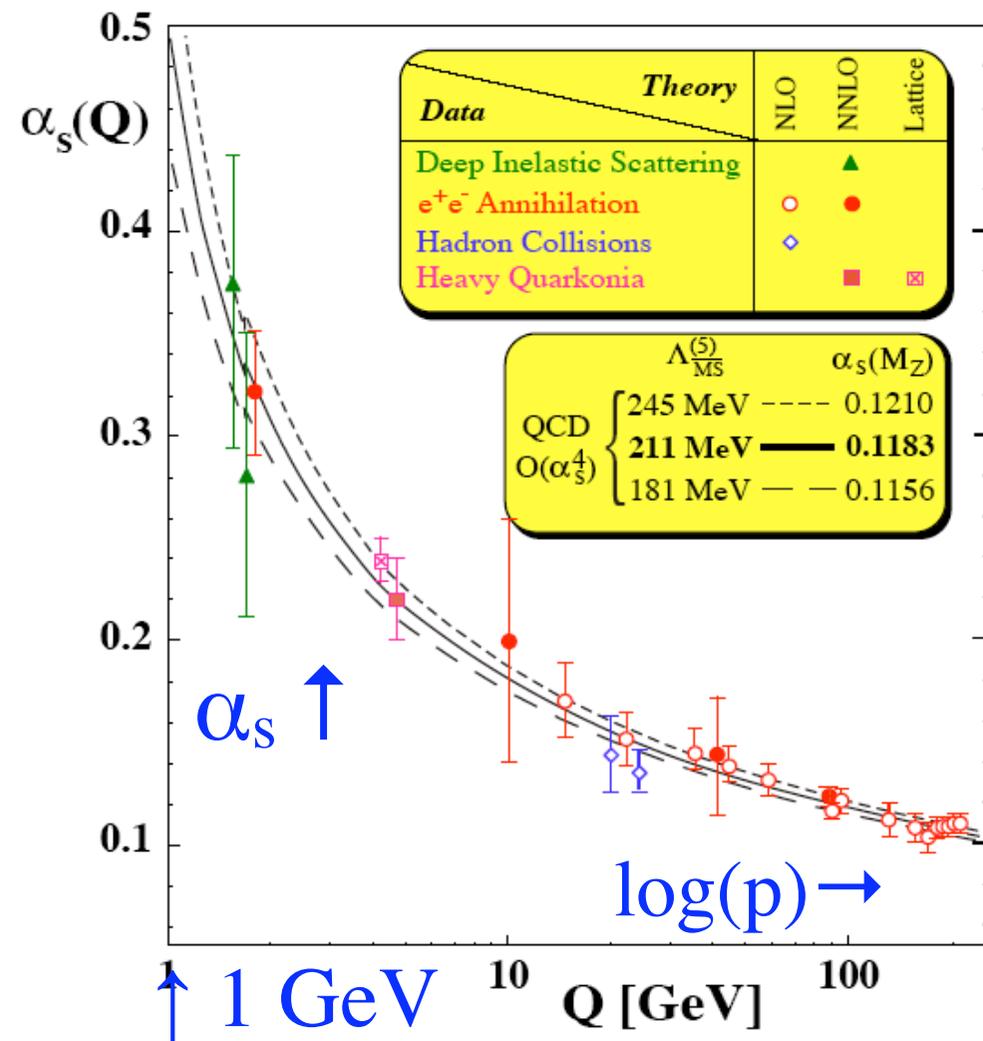
NON-perturbative: < 1 GeV

Powers of $(non-pert. \text{ mass})^2/p^2$

Quarks: six flavors, coupled to glue

“2+1” light: up (u), down (d) + strange (s)

Heavy: charm (c), bottom, top.



QCD & Confinement

Non-perturbative effects => confinement. *Only color singlets observable.*

Quark + anti-quark = meson.

Very light: pions ($\bar{u}d$...), mass = 140 MeV. kaons, ($\bar{u}s$...), 490 MeV
= Spin waves of global (“chiral”) symmetry for light quarks, u, d, & s.

Three quarks = baryon.

Neutrons, protons (ddu, uud), 940 MeV; also strange (uds ...), 1-2 GeV

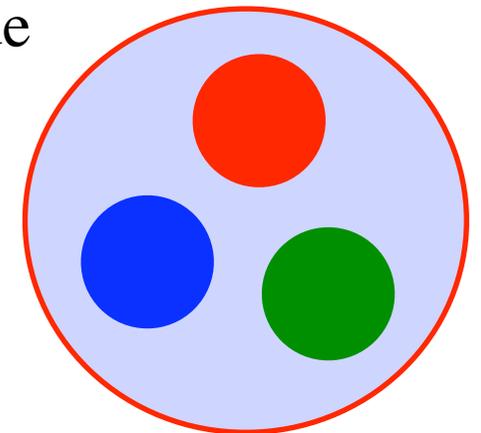
Charm mesons, J/Psi, ($\bar{c}c$) 3.1 GeV. *Glueballs? Where’s the glue?*

One model: MIT bag. Confine quarks & gluons inside bag =>

Costs Bag constant $B \sim (200 \text{ MeV})^4$; \sim perturbative inside

Dynamics of bag? *Surface has zero width*

Spectrum OK, \sim non-relativistic quark model



Large N Expansion

Instead of SU(3) color, consider SU(N). Claim: N=3 is “near” infinite N!

Gluons: SU(N) *matrix*, # = N^2-1 . Quarks: SU(N) *vector*, so # $\sim N$.
=> large N expansion is *dominated* by glue.

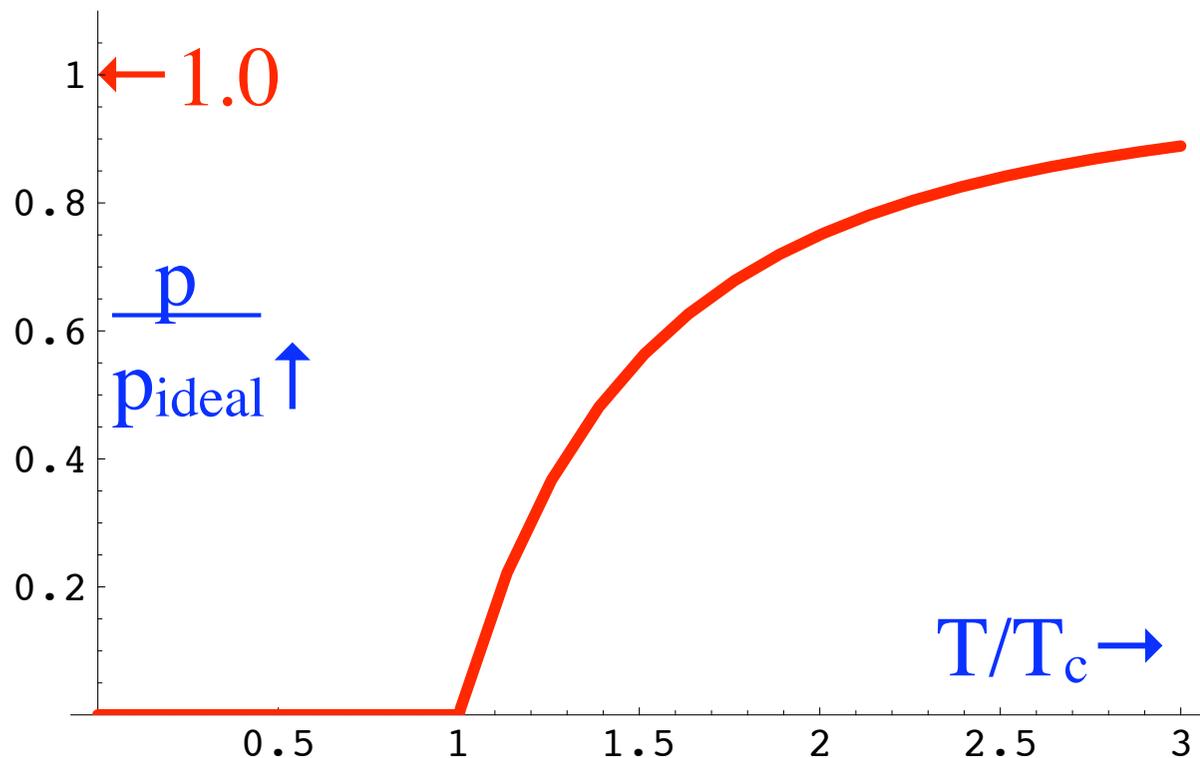
AF => *ideal* gas of $\sim N^2$ gluons at $T = \infty$. Pressure $p(T) \sim N^2$ for $T > “T_c”$

Confinement => only color singlets, # ~ 1 . So $p(T) \sim 1$ for $T < “T_c”$

At infinite N:
pressure/(ideal gas pressure)
is an *order parameter*:

Ratio $\neq 0$ above T_c , $\rightarrow 1$ as $T \rightarrow \infty$.
= 0 below T_c ($\sim 1/N^2$)

Possible form for this ratio =>
Like magnetization, in “reverse”



Numerical “experiments”: the lattice

How to compute non-perturbatively? Put QCD on a lattice, with spacing “ a ”.

In QCD, AF + renormalization group \Rightarrow *unique* answer as $a \rightarrow 0$, continuum limit

Approximate ∞ -dim. integral by finite dim.: “Monte Carlo”

In practice: need to “improve” to control lattice effects $\sim a^2$, etc.

For pure gauge (*no* quarks) near continuum limit from mid-’90’s!

Compute equilibrium thermodynamics: pressure(temperature) = $p(T)$

Lattices: 6-8 steps in time, 24-32 in space $\Rightarrow 8 \cdot 32^3 \cdot 8 = 10^6$ dim. integral

With dynamical quarks, much harder: today, *not* near continuum for light quarks

Quarks non-local, difficult to treat (global!) “chiral” symmetry of light quarks

Today: only $T \neq 0$, *not* dense quarks at $T=0$

With quarks: *orders* of magnitude more difficult

BIG calculations done on BIG machines...

“QCDOC”, QCD on a chip

RIKEN/BNL: 12,000 nodes, 20 teraflops...



Lattice: SU(3) gluons deconfine, $T_c \sim 270 \text{ MeV}$

Confined phase at $T=0$: only “glueballs”. *Lightest* glueball 1.5 GeV

Deconfinement at $T_c \sim 270 \text{ MeV} \pm 10\%$ *Small* relative to glueball masses.

Pressure *very* small below T_c : *very* much like large N!

MIT bag model does *not* work $p_{\text{MIT}}(T) = \#T^4 - B$

“Fuzzy” bag does:

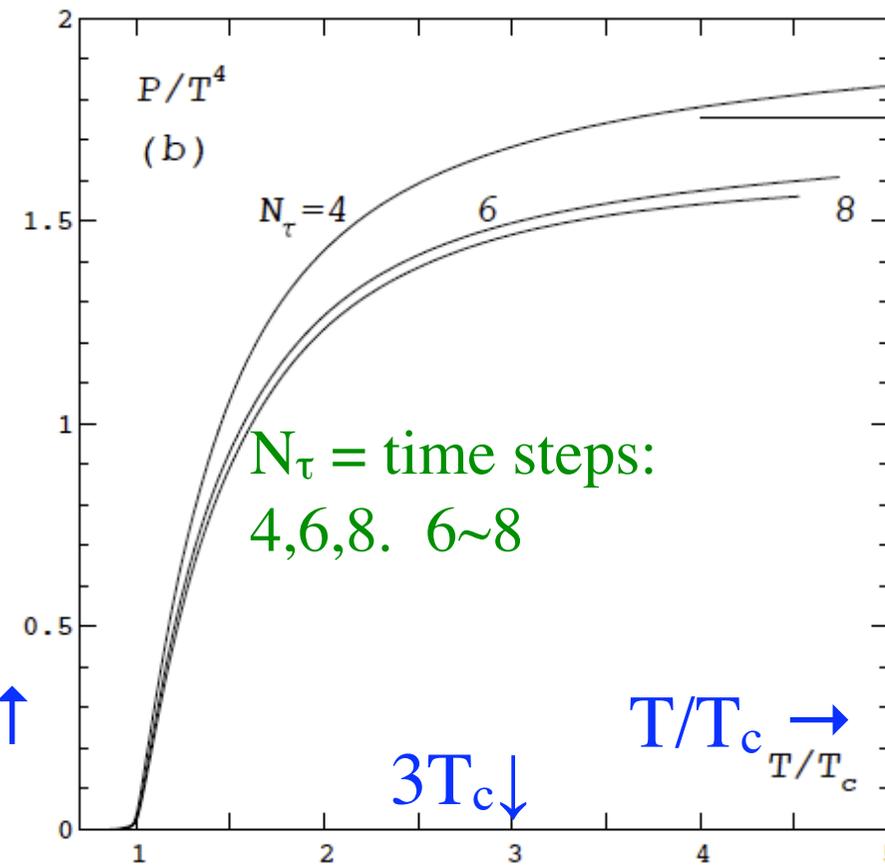
$$p_{\text{fuzzy}}(T) = \#(T^4 - T_c^2 T^2)$$

$\# = 90\%$ of ideal gas $\Rightarrow T^4 \sim$ perturbative

Term $\sim T^2$ non-pert. power law correction

Transition weakly first order
 $\sim Z(3)$ Potts model

$p(T)/T^4 \uparrow$



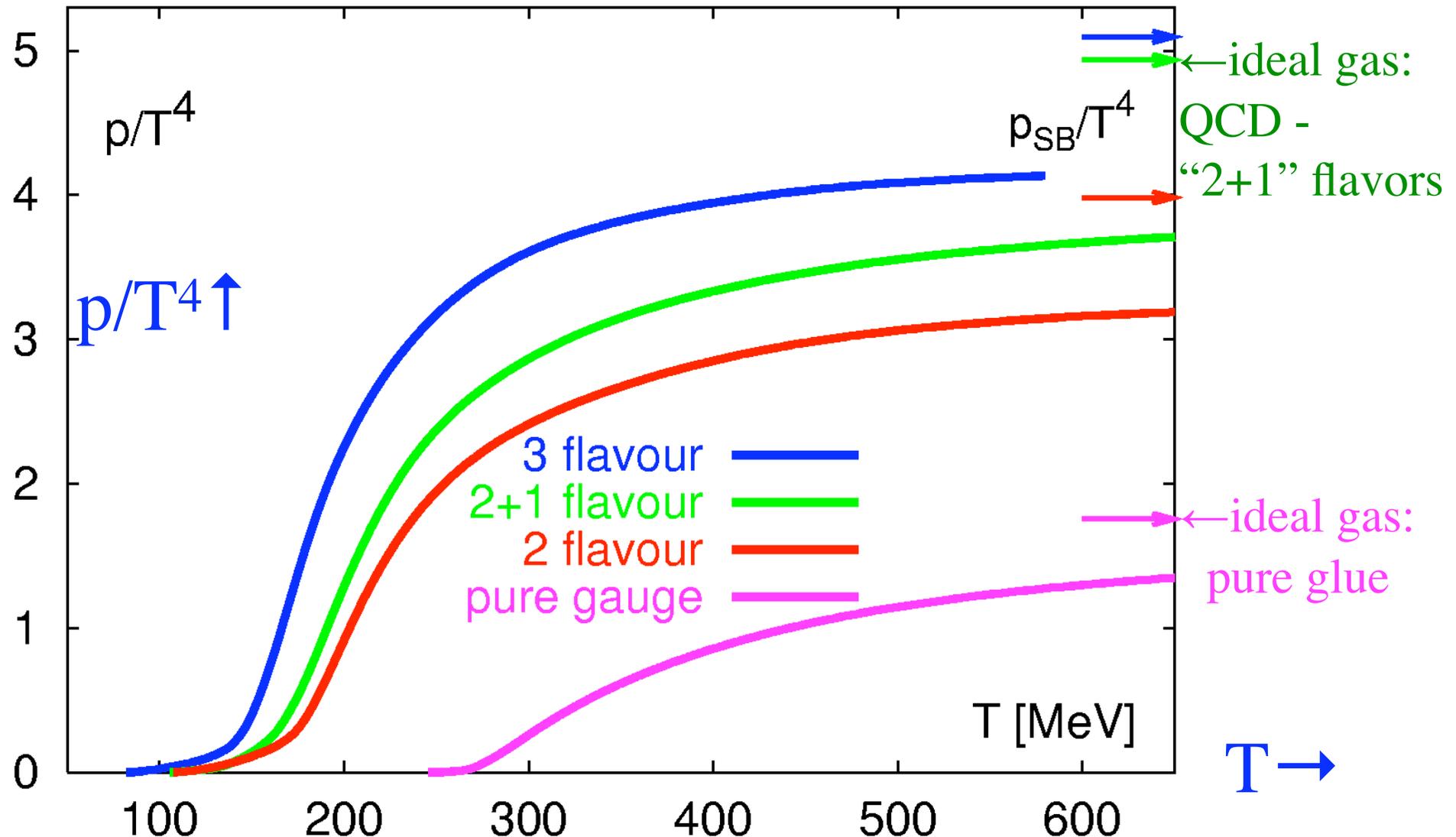
Lattice: QCD, with quarks: $T_c \sim$ “175” MeV

Adding quarks: ideal gas value increases by ~ 3

“ T_c ” decreases. Today: ~ 175 MeV \pm ? 1st order transition washed out, *crossover*.

Pressure nonzero $< T_c$: from pions, kaons..., masses $\sim T_c$

Pressure within 20% of ideal by $\sim 3 \times T_c$. Different behavior?



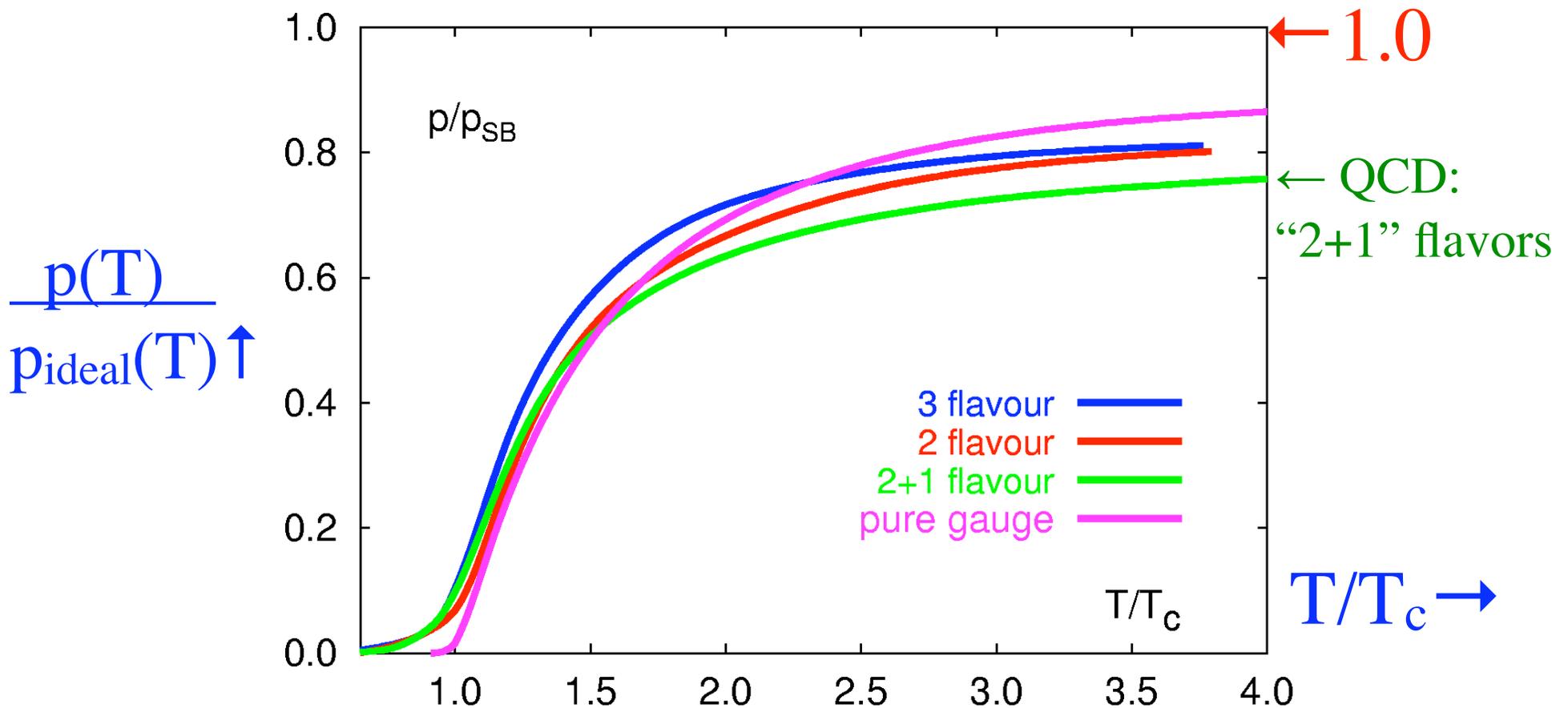
Lattice thermo.: QCD & “flavor independence”

Bielefeld: properly scaled, \approx *universal* pressure
Transition dominated by deconfinement

$$\frac{p}{p_{ideal}} \left(\frac{T}{T_c} \right) \approx \text{const.}$$

In *all* cases so far, chiral symmetry restored at *same* T_c as deconfinement: *why?*

Is transition from confined phase, directly to nearly ideal QGP?

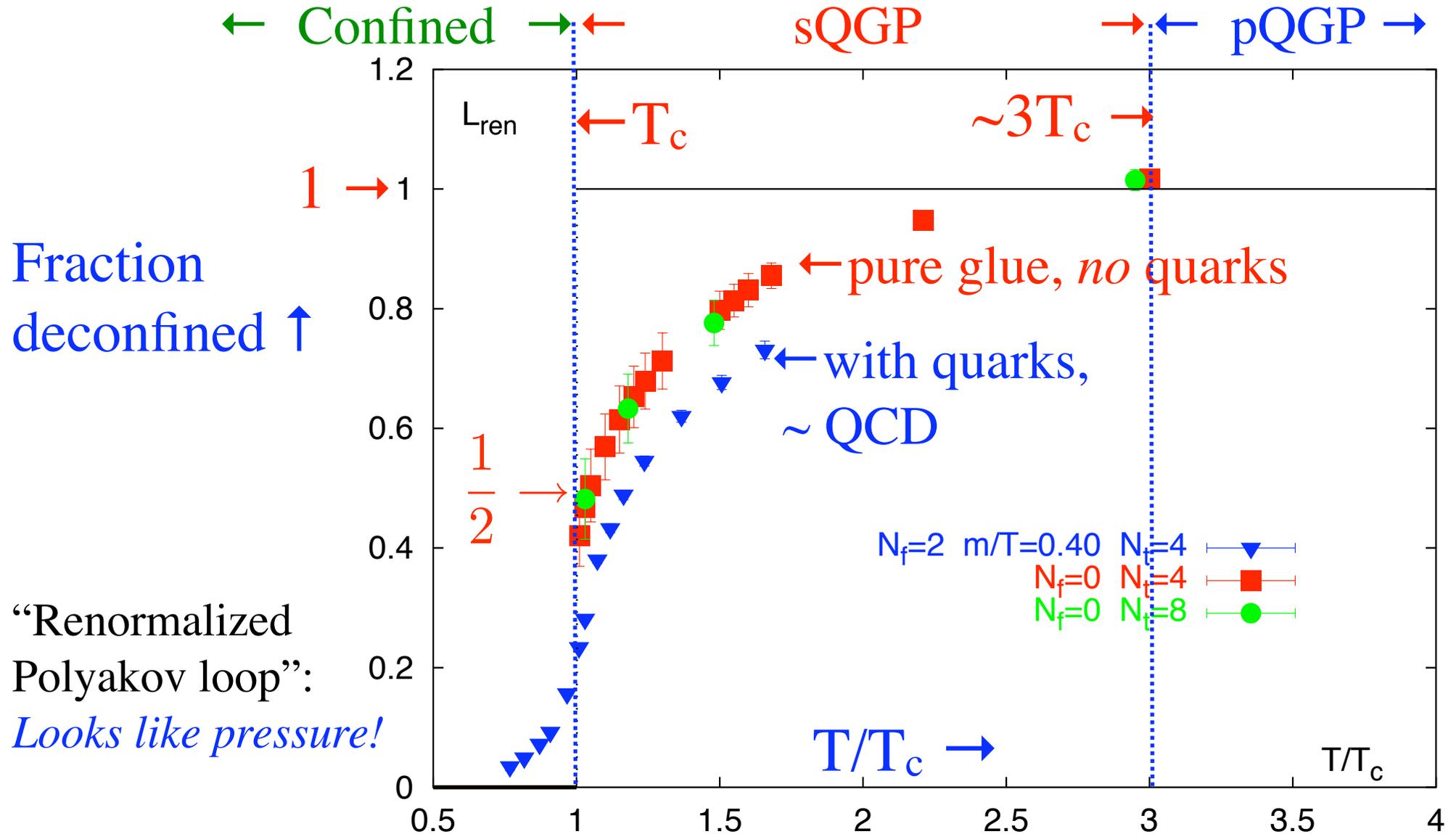


Partially deconfined, T_c to $\sim 3 T_c$: “sQGP”

On lattice, can measure *fraction of deconfinement*

$T > 3 T_c$: fraction ~ 1 , \Rightarrow \sim perturbative QGP, “pQGP”

$T: T_c \rightarrow 3 T_c$: fraction < 1 , only *partially* deconfined: “sQGP” (s?)



The “sQGP”: *not* strong coupling

“Helsinki Dogma”: resum pert. thy. to 4 loop order. *Fails* below $\sim 3 T_c$

Sure. Coupling big below $3 T_c \sim 500$ MeV?

No! While deconfinement only partial at T_c , coupling *not* big, even at T_c

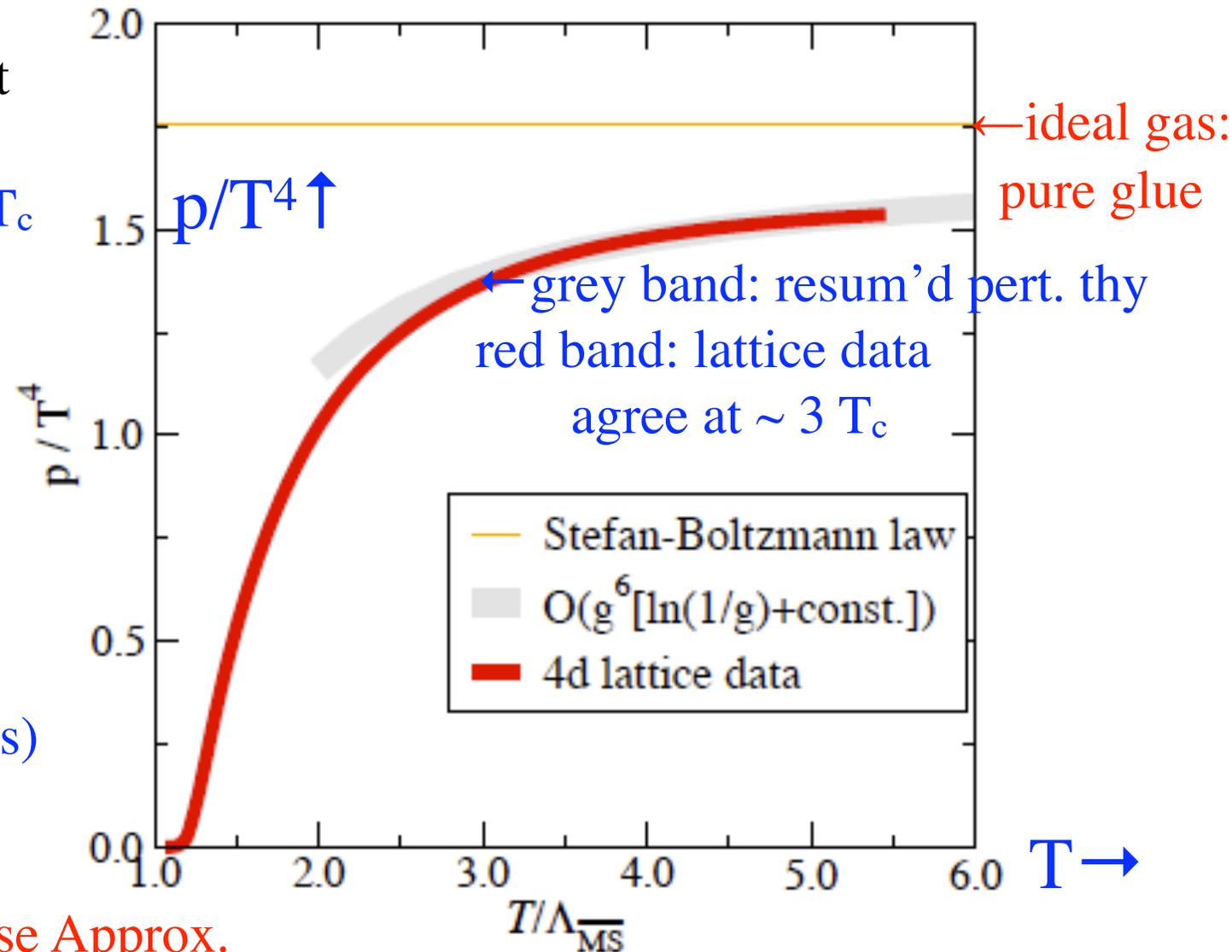
$\alpha_s^{\text{eff}}(T_c) \sim \alpha_s(1.6 \text{ GeV})!$

So what *is* “sQGP”, T_c to $\sim 3 T_c$?

Semiclassical, dominated by “solitons” (Wilson lines)

Eventually - NaRPA:

Non-abelian Random Phase Approx.



Hunting for the “Unicorn” in Heavy Ion Collisions



Unicorn = QGP. Hunters = experimentalists. So: all theorists are dogs...

AA collisions at high energies

Collide: pp, protons on protons. Benchmark for “ordinary” QCD.

AA, nuclei on nuclei. Atomic # “A”: 60 => 200, Cu -> Au, Pb. “Hot” nuclei.

pA, proton (or deuteron) on nucleus. Another check: QCD in “cold” nuclei

Why AA? Baryons are like hard spheres, so for A: 60 - 200,
transverse size $A^{2/3}$: 15 - 35 \times proton. **Big nuclei are big!**

Total energy in the center of mass, $E_{\text{cm}} = \sqrt{s}$ (GeV); per nucleon, $\sqrt{s}/A = \sqrt{s_{\text{NN}}}$.

SPS @ CERN	5 => 17 GeV
RHIC @ BNL	20 => 200 GeV
LHC @ CERN	5500 GeV
SIS @ GSI	2 => 6 GeV

SPS, SIS Fixed Target

RHIC, LHC Colliders

LHC > '08

SIS @ GSI, Darmstadt > '12



Geometry of AA collisions at high energy

At *high energies*: nuclei Lorentz contracted along beam (15 fermi \Rightarrow 0.3 fermi)

AF \Rightarrow nuclei don't stop, pass through each other.

Collider: lab = center of mass frame

Momenta of produced particles: along beam, p_z ; transverse to beam, p_t

Baryons in original nuclei go down beam pipe, at large $\pm p_z$

For pp collisions: particles \sim constant for some range in p_z about $p_z = 0$:

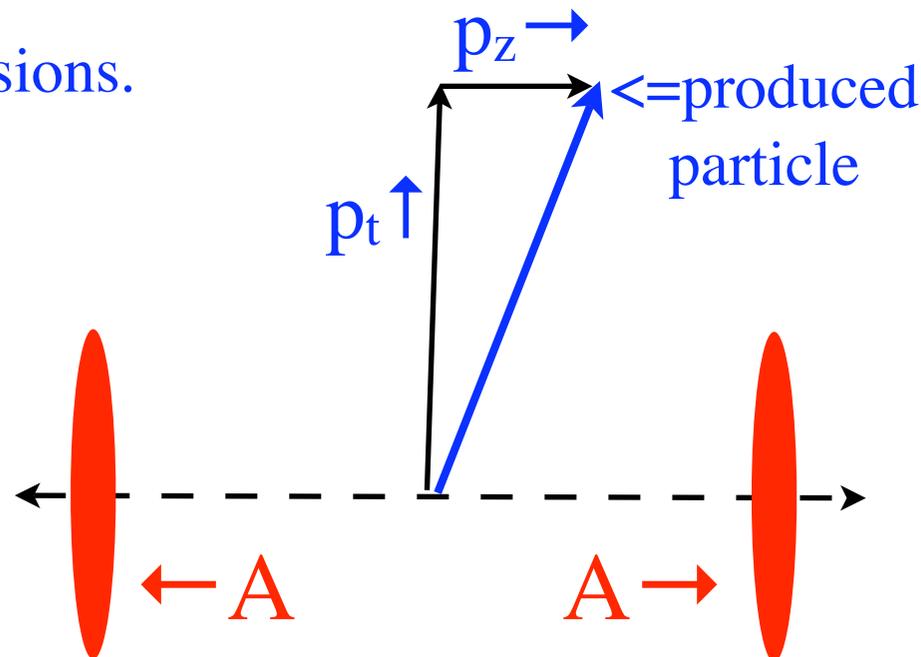
“rapidity plateau” (rapidity $\sim p_z$; boost invariance)

Bjorken: look at rapidity plateau in AA collisions.

Rapidity plateau \sim free of incident baryons.

\Rightarrow most likely to be at nonzero temperature,
zero (quark) density.

Collider: central plateau 90° to beam



Typical Au-Au collision @ RHIC

Experiments @ RHIC:

“Big”: ~ 400 people. STAR & PHENIX

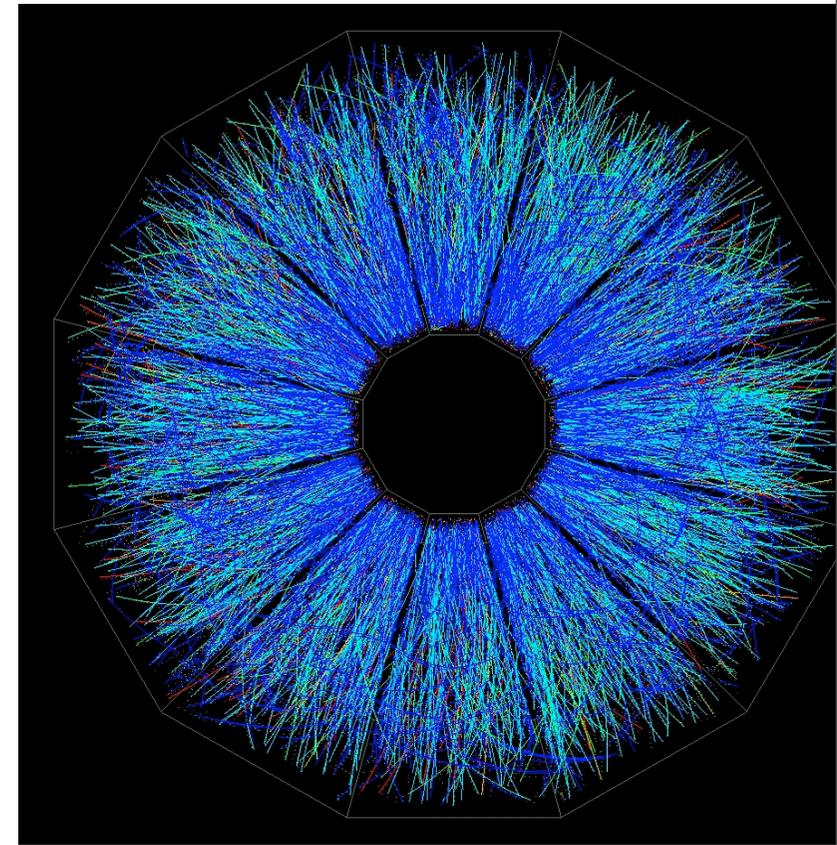
“Small”: ~ 50 people. PHOBOS & BRAHMS

No surprises in total multiplicity; $\sim 1.3 A \times pp$:

total # particles $\sim \sim \log(\text{total energy})$

\sim total # experimentalists

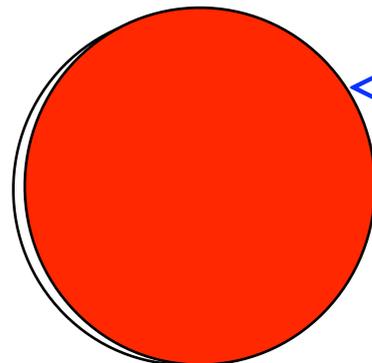
theorists $\sim \log(\log(\text{total energy}))$.



Need hunters more than dogs...

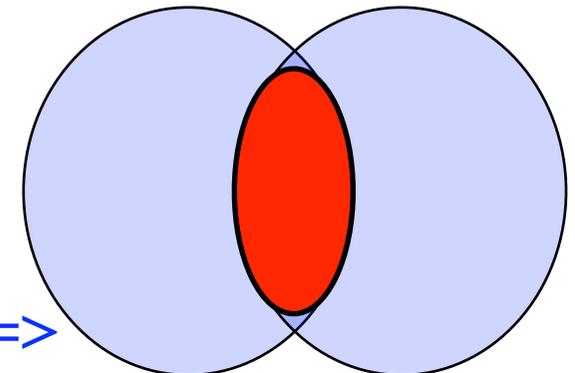
Total # particles/unit rapidity $\sim 900 \uparrow$

Also: can exp.'y measure how much nuclei overlap in transverse plane



\leftarrow central

peripheral \Rightarrow



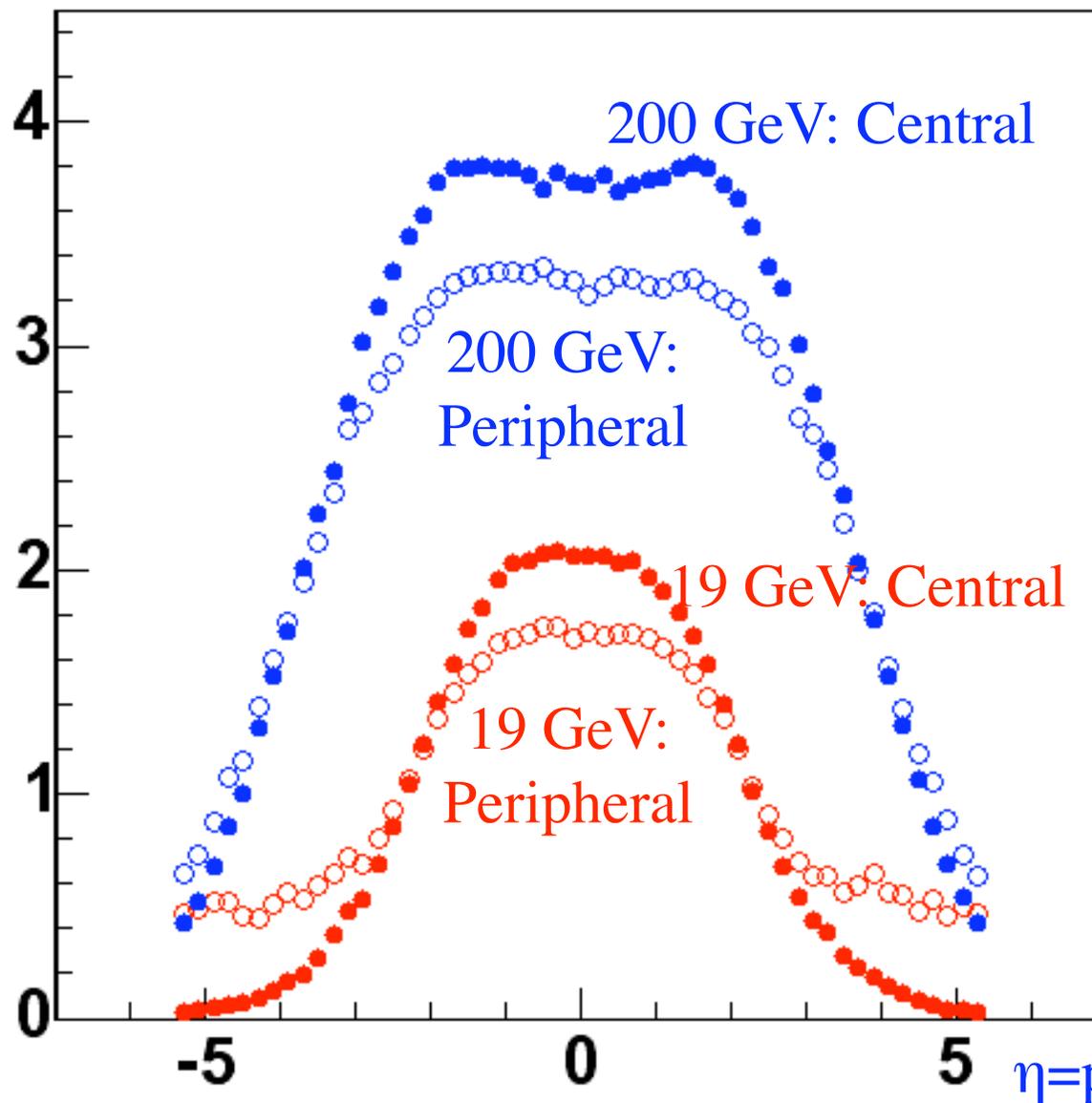
Overall multiplicity: slow growth

$dN/d\eta/\uparrow$

$N = \#$ particles

$\eta =$ "pseudo"-
rapidity
(no particle ID)

/ by #
"participants"



200 GeV, RHIC
900 particles
/unit η

19 GeV, SPS
600 particles
/unit η

No big increases in multiplicity, as predicted by cascade models.

Rapidity plateau $\pm .5$ (out of ± 5.0) for dN/dy ($y =$ true rapidity)

The Tail Wags the (Dog) Unicorn

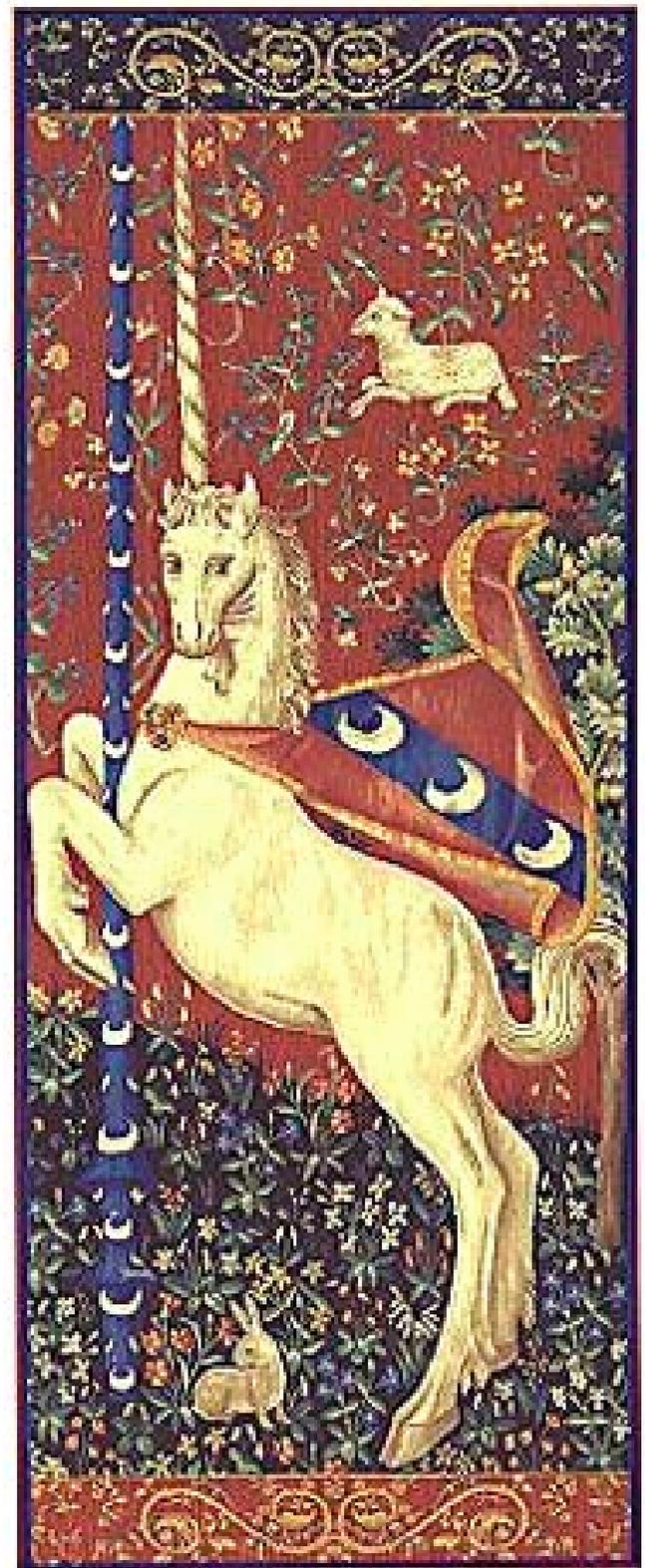
Body of the “Unicorn”:

For $T \sim 200$ MeV,
majority of particles at *small* momenta,
 $p_t < 2$ GeV.

Tail of the “Unicorn”:

Look at particles at *high* momentum,
 $p_t > 2$ GeV, to probe the body.

Concentrate on zero rapidity, 90° to beam.

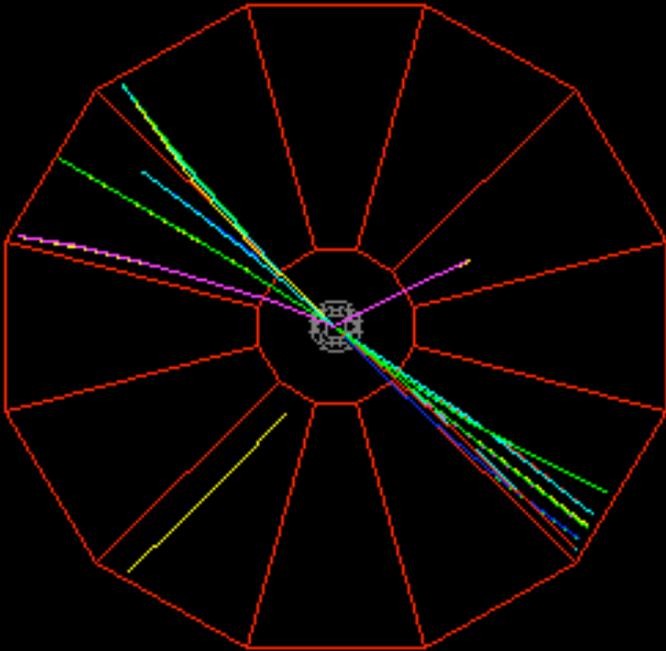


Jets: “seeing” quarks and gluons in QCD

At high transverse momentum (p_t), instead of hadrons, have \sim quarks & gluons: *jets*.

≤ 2 jets from pp collision at RHIC.

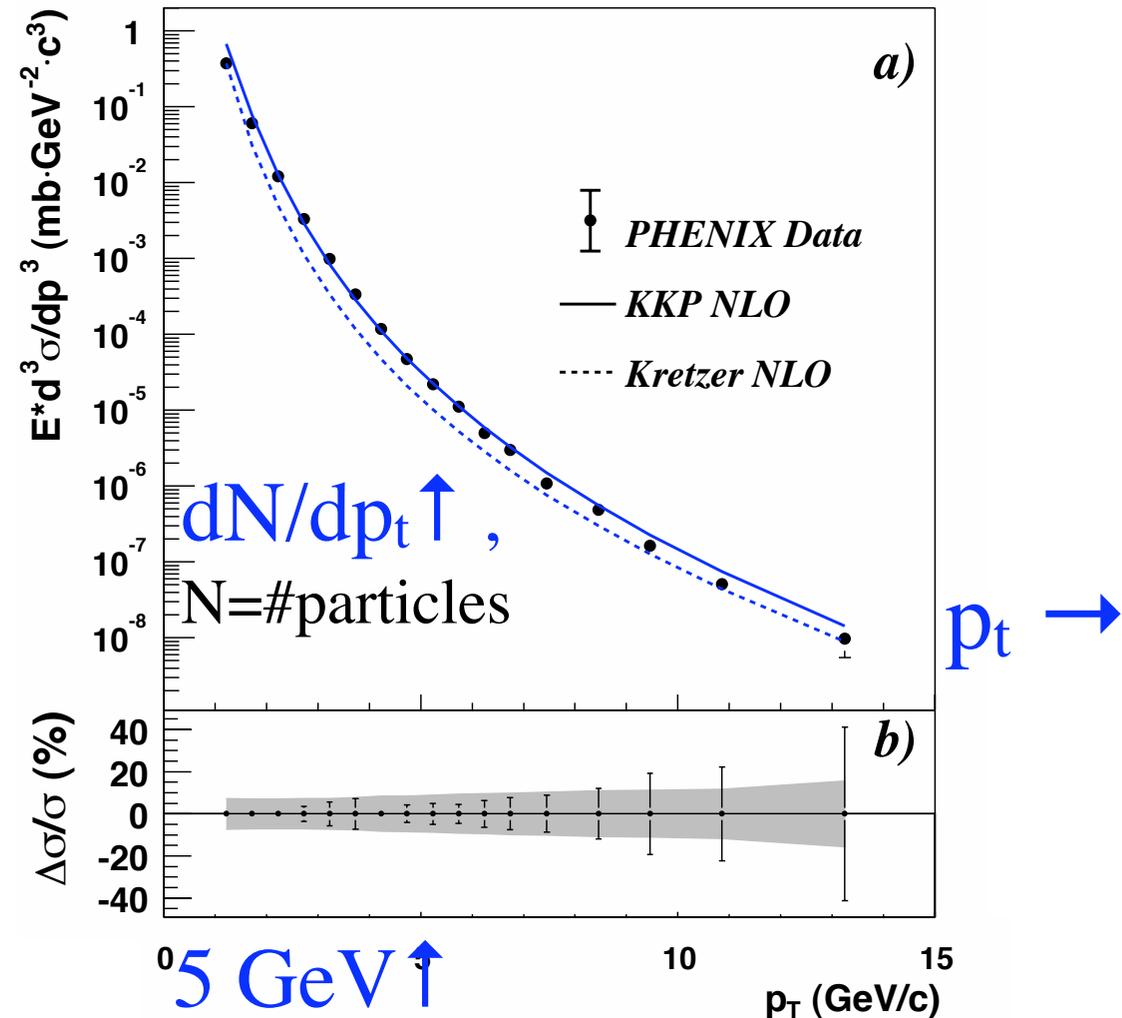
By momentum conservation, for each jet, there is a backward jet.



In pp coll.'s, jets can be computed for large p_t , down to $p_t \sim 1$ GeV.

Jets *rare*: particles at $p_t \sim 2$ GeV \sim .1 % of total!

Look at jets in AA



“ R_{AA} ”: *robust* signal of new physics

R_{AA} = for a given p_t , # particles central AA / ($A^{4/3}$ # particles pp)

For π^0 's, $p_t : 2 \rightarrow 20$ GeV, $R_{AA} \sim 0.2$. As if jets emitted *only* from surface!

Due to “energy loss” in thermal medium?

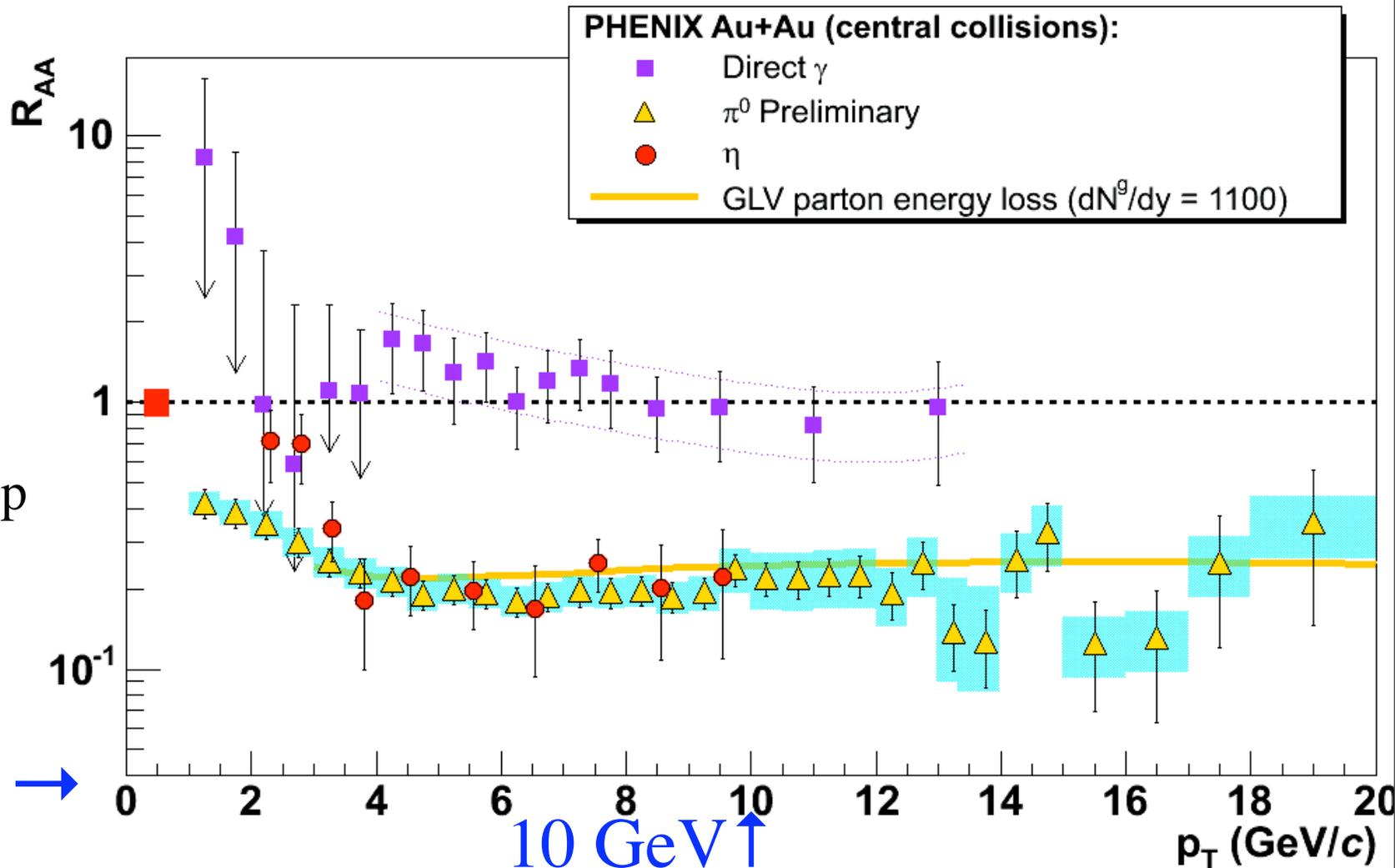
$A^{4/3}$: *experimentally*: for γ 's, $R_{AA} \sim 1.0$ π^0 's “eaten”, γ 's not

R_{AA} : \uparrow

particles
central AA/
particles pp

$A=200 \Rightarrow$

$p_t \rightarrow$

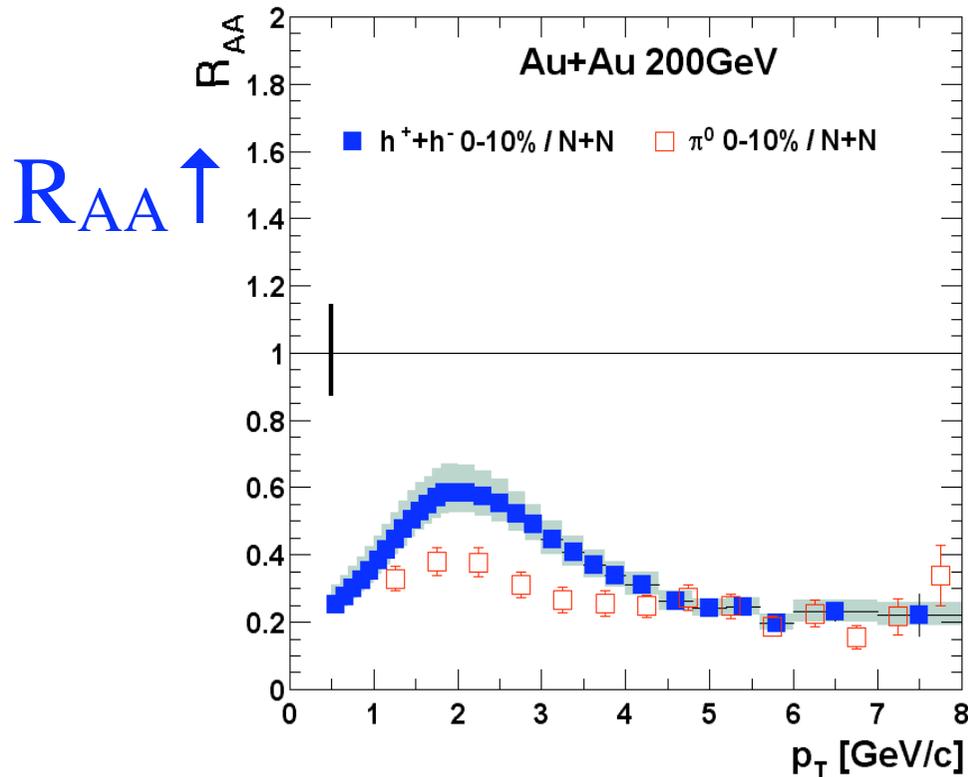


R_{AA} final state effect: *not* in R_{dA}

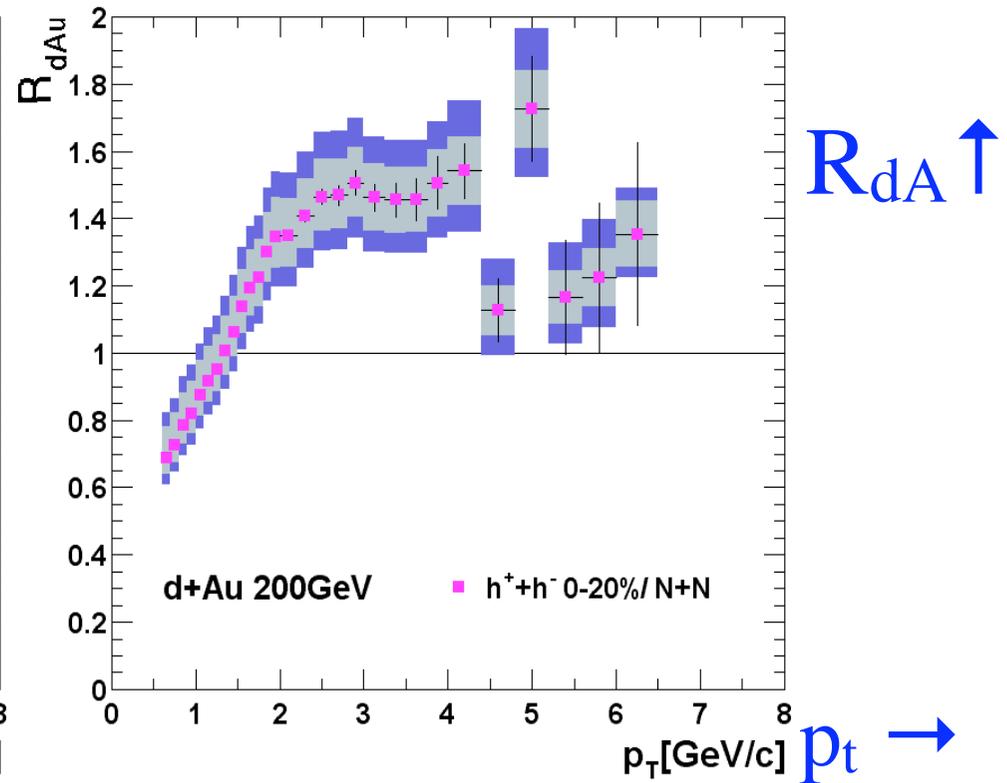
R_{dA} : like R_{AA} , but for dA coll.'s/pp coll.'s. At zero rapidity:

dA: *enhancement*, from initial state “Cronin” effect ($\Rightarrow 1$ @ $p_t > 8$ GeV)

AA: *suppression* \Rightarrow final state effect



Suppression in AA \uparrow
 $R_{AA} \sim 0.4$ @ 3 GeV



Enhancement in dA \uparrow
 $R_{dA} \sim 1.4$ @ 3 GeV

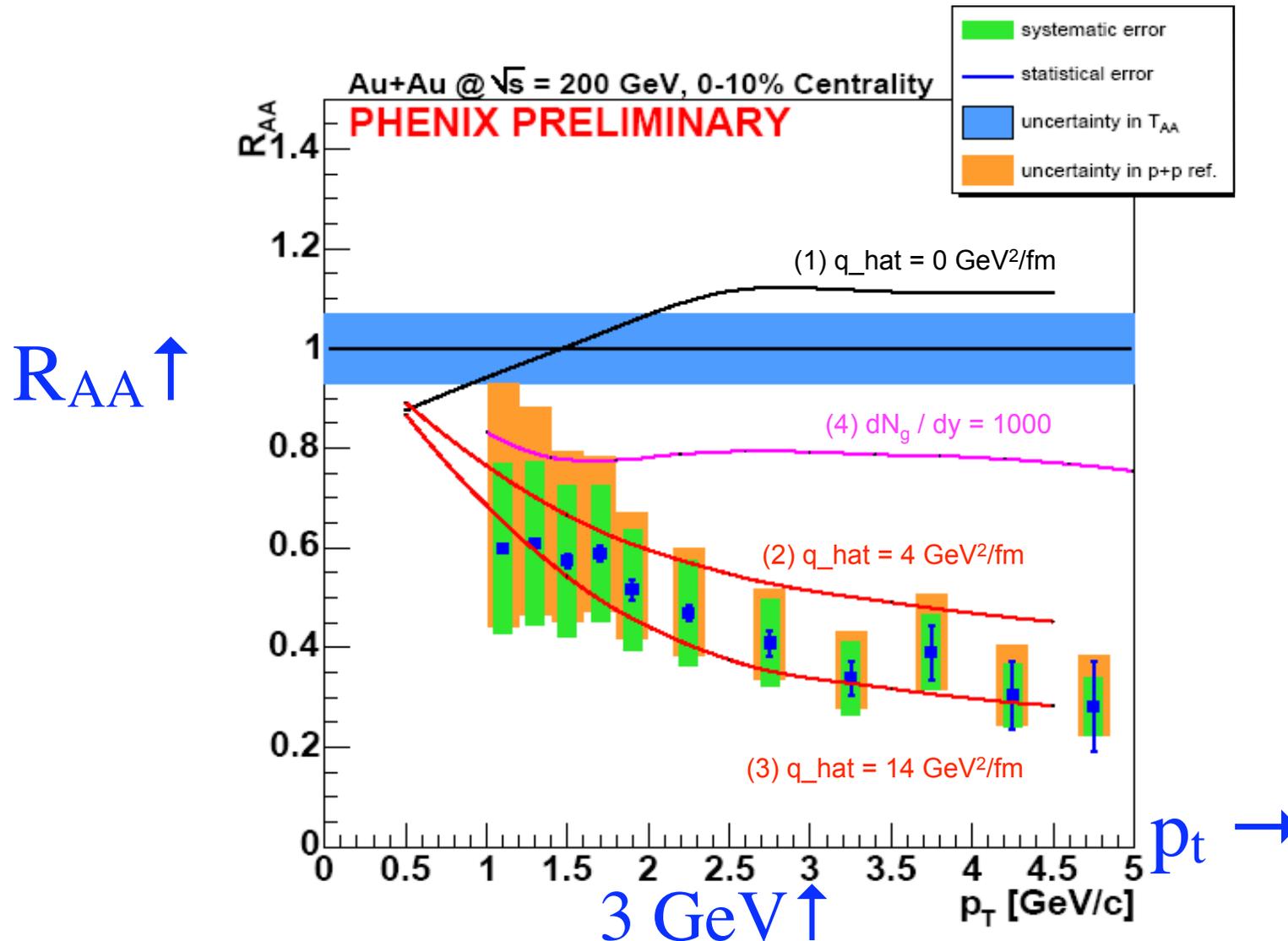
R_{AA} for heavy quarks: also suppressed!

PHENIX: R_{AA} for charm quarks \sim light quarks!

Mass of charm quark $m_{\text{charm}} \sim 1.5 \text{ GeV}$; $T \sim 200 \text{ MeV}$.

Heavy quark less sensitive to medium by $T/m_{\text{charm}} \sim 1/8$. *No sign of that!*

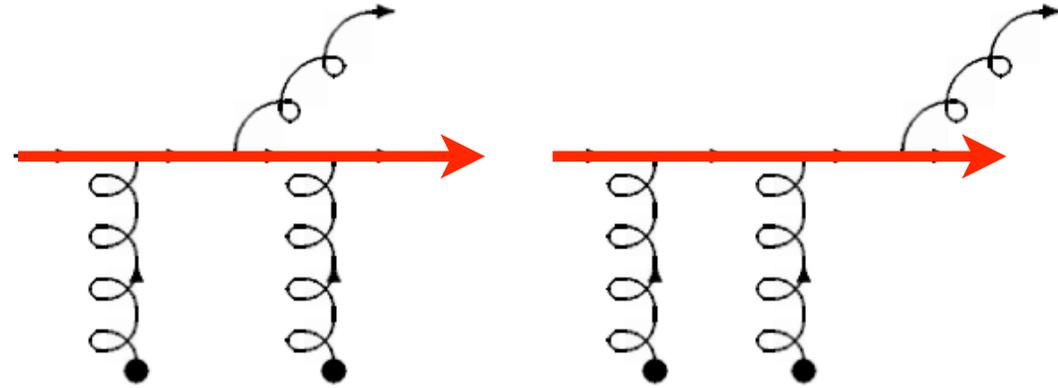
Experimental evidence for “sQGP”: heavy quarks \sim same as light!



Theory of R_{AA} : energy loss

Fast quark (or gluon) emits radiation, scatters off of thermal bath.

Like QED, but qualitatively different in QCD: emitted gluon emits more gluons, etc.



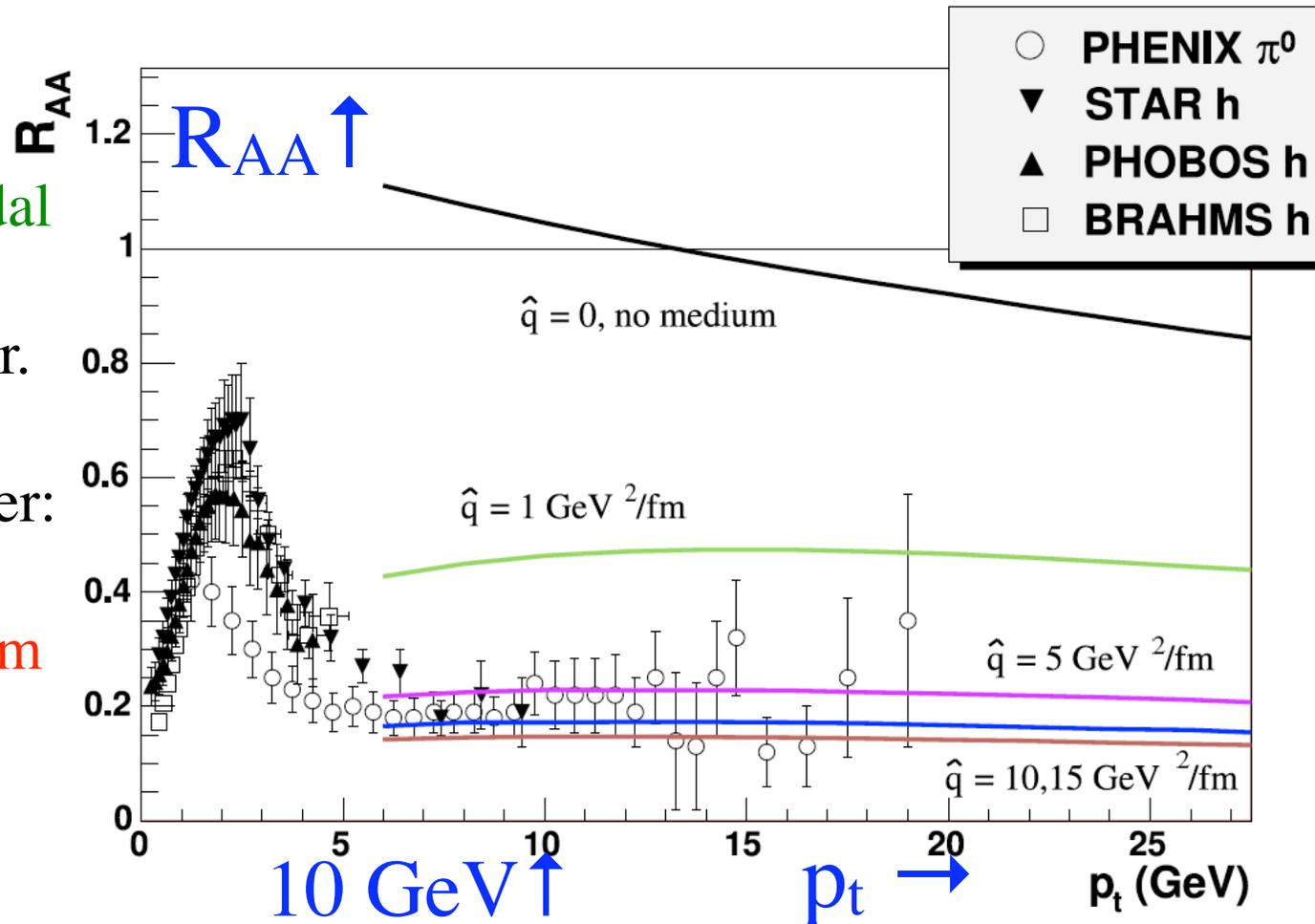
Involves LPM:
Landau-Pomeranchuk-Migdal

Parametrized by one number.

Theorists disagree on number:

“weak” coupling $\sim 2 \text{ GeV}^2/\text{fm}$

or “strong” $\sim 15 \text{ GeV}^2/\text{fm}$?



Central AA collisions “eat” jets!

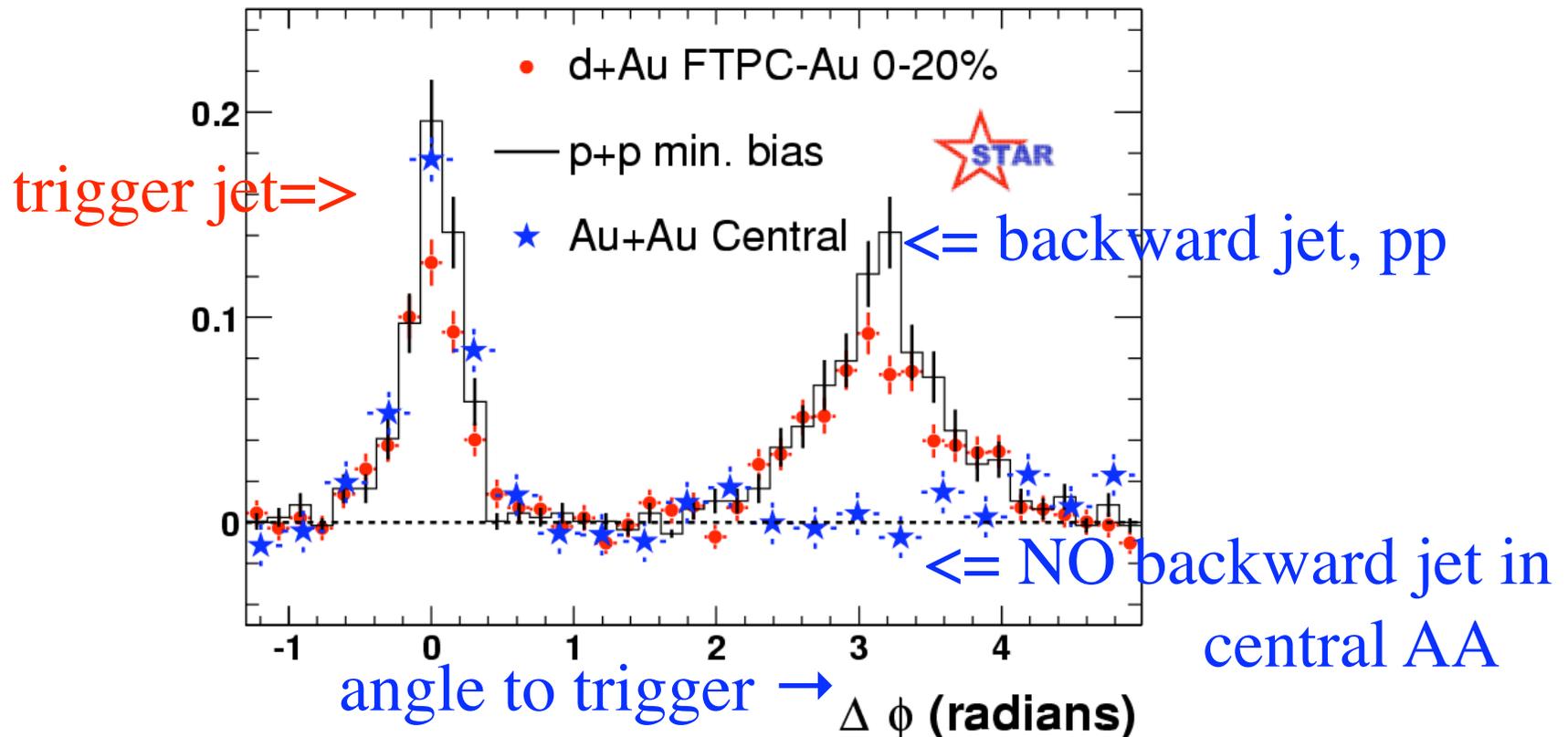
Unlike pp, in central AA, cannot trigger on individual jets. Can:

Trigger on hard “jet”, $p_t: 4 \rightarrow 6$ GeV. Look for backward “jet”, $p_t > 2$ GeV

In pp or dAu collisions, *clearly* see backward jet.

In central Au-Au, away side jet gone: “*stuff*” in central AA “*eats*” jets

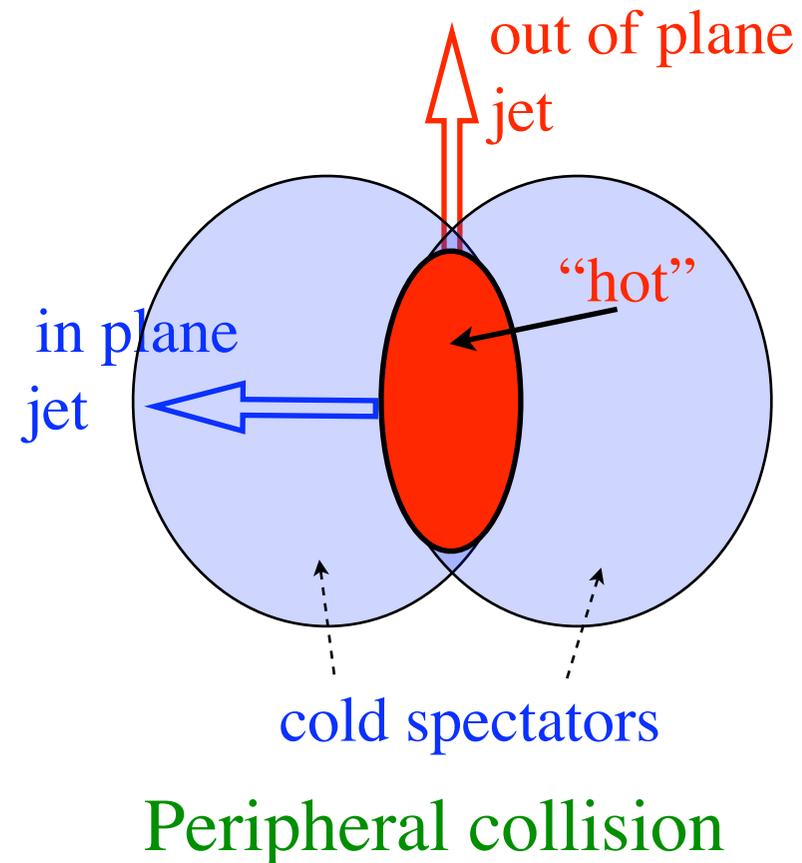
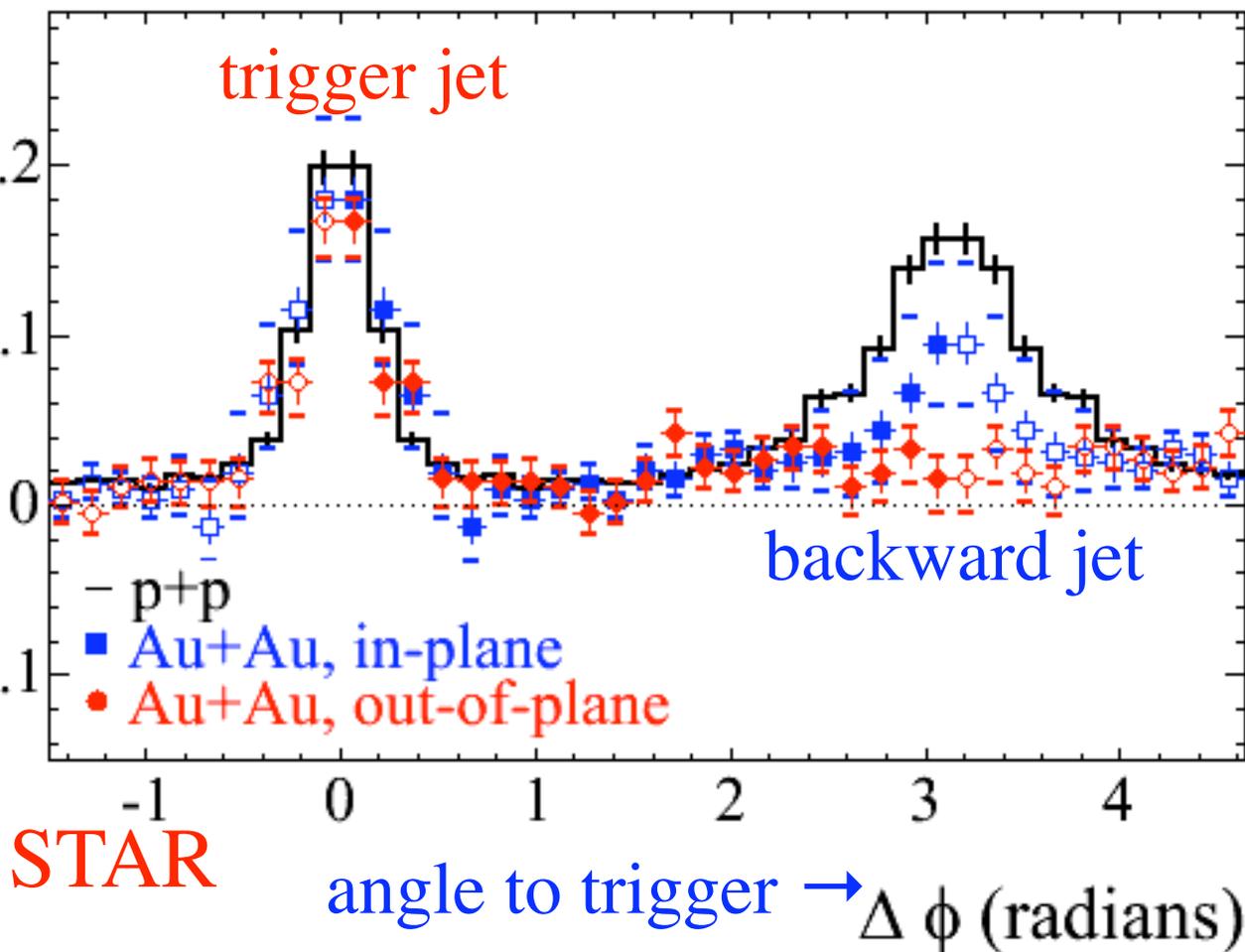
Adams *et al.*, Phys. Rev. Let. 91 (2003)



Peripheral coll.'s: more jets eaten *out* of plane

Peripheral collisions: “hot stuff” forms “almond”. In vs. out of reaction plane
Out: more “hot stuff”. *In*: less hot stuff, more cold nuclear matter

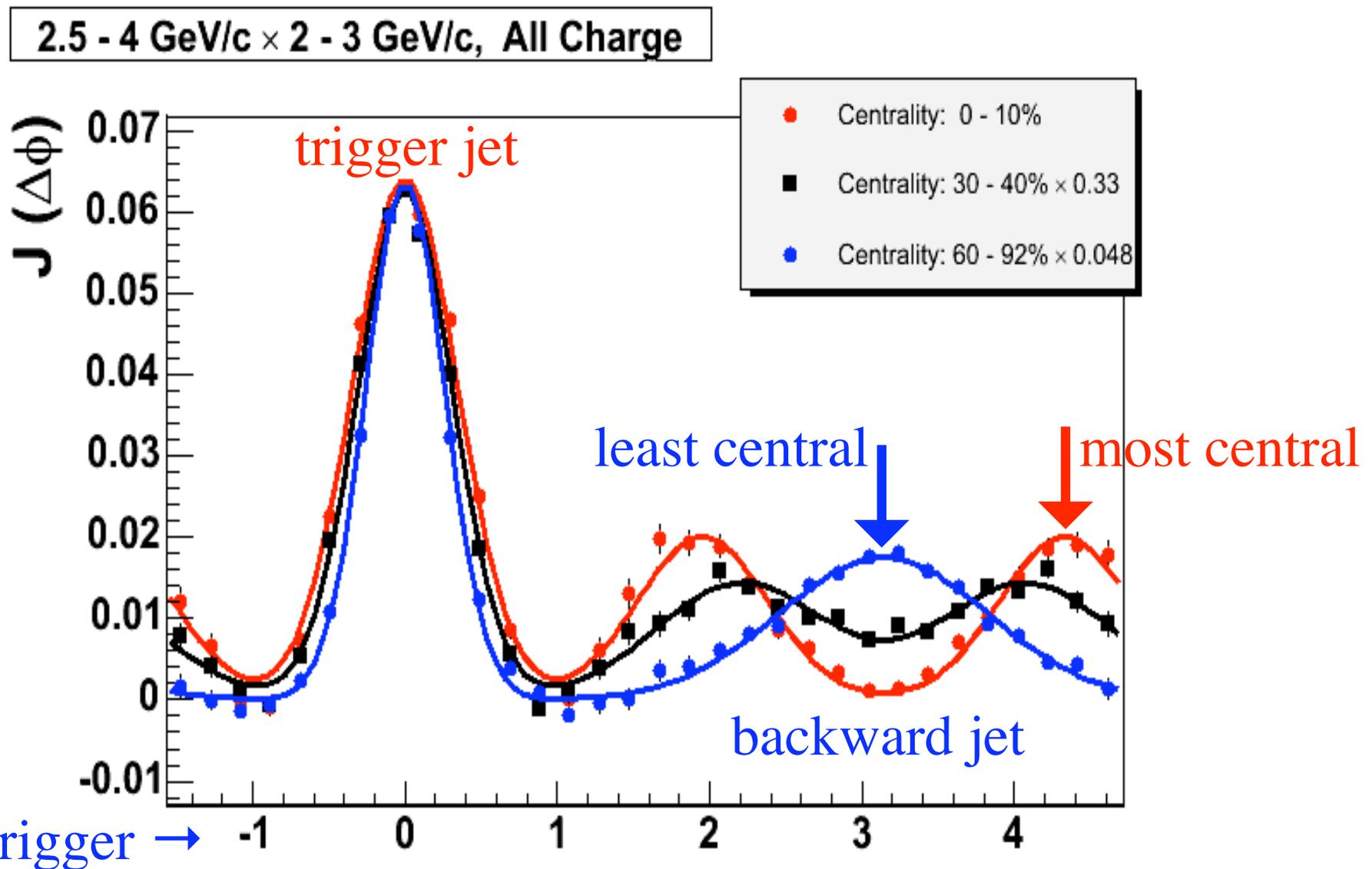
Exp.'y: backward jet more strongly suppressed *out* of plane than *in* plane =>
Geometrical test that central AA “eats” jets preferentially



AA collisions modifies jet shapes

PHENIX: shape of away side jet appears to be modified by “stuff”:

Mach cone or Cerenkov radiation?



The Body of the Unicorn: the sQGP

Particles peaked about zero transverse momentum

$T_c \sim 200$ MeV: expect thermal to $p_t \sim 2$ GeV.

Thousands of particles: use hydrodynamics? “Most Perfect Fluid on Earth”

me →



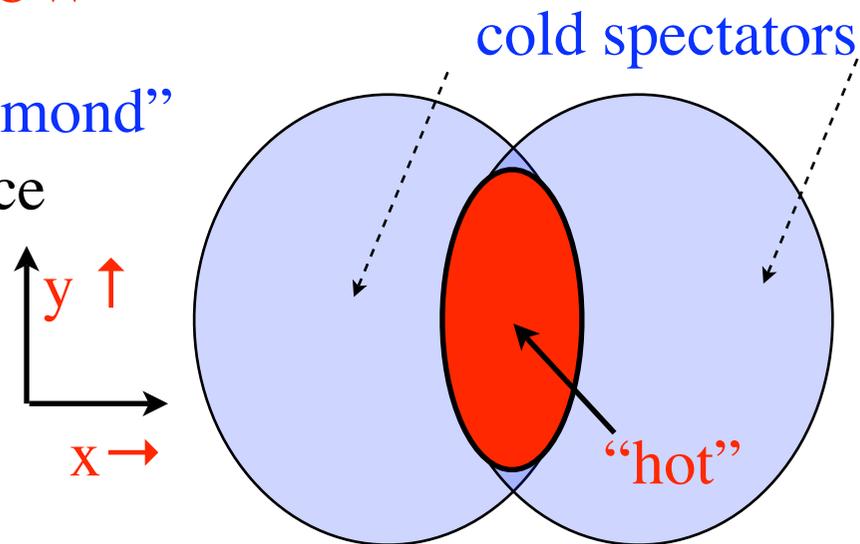
← Unicorn

“Elliptic Flow”

For peripheral collisions, overlap region is “almond” in coordinate space, sphere in momentum space

So start with spatial anisotropy,

$$\epsilon = \frac{\langle y^2 - x^2 \rangle}{\langle x^2 + y^2 \rangle}$$

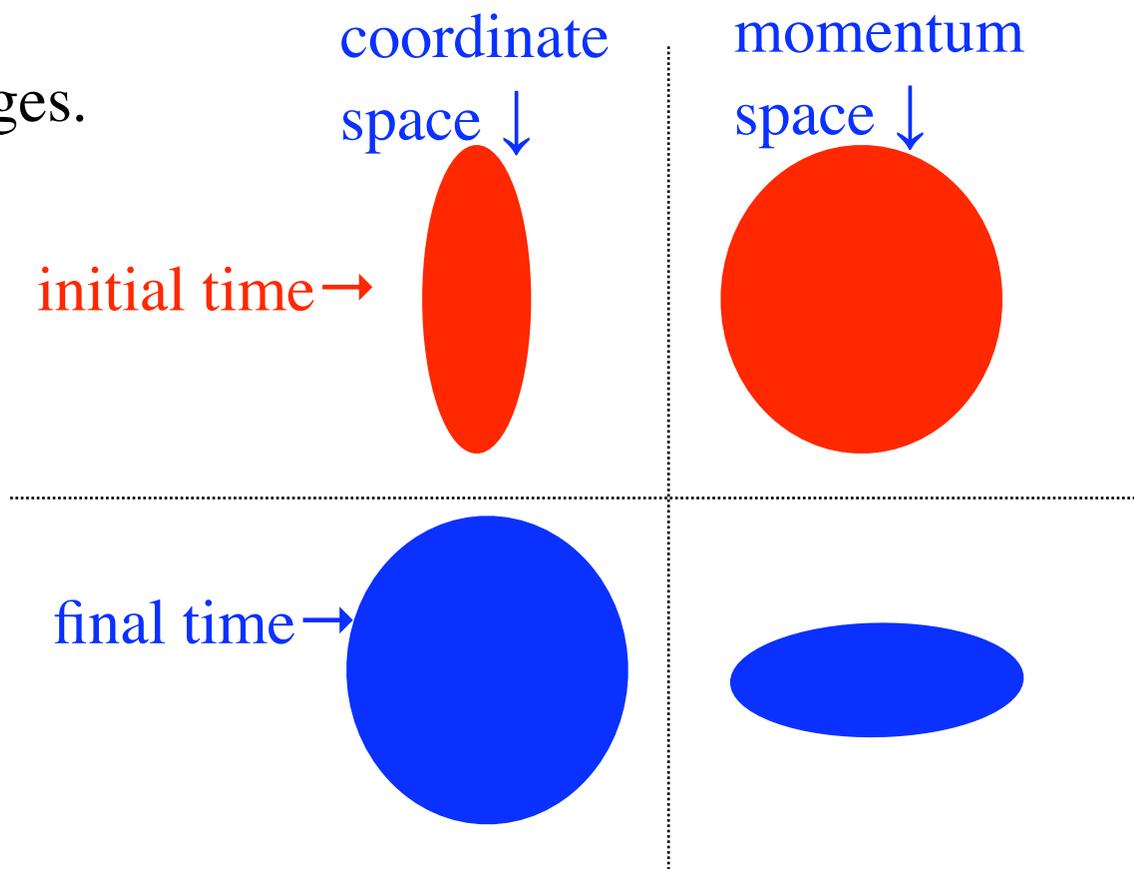


If particles free stream, nothing changes.

If collective effects present, end up with sphere in coordinate space, almond in momentum space:

“elliptic flow”

$$v_2 = \frac{\langle p_y^2 - p_x^2 \rangle}{\langle p_x^2 + p_y^2 \rangle}$$



“Most Perfect Fluid on Earth”

Large # particles: try *ideal* hydrodynamics:

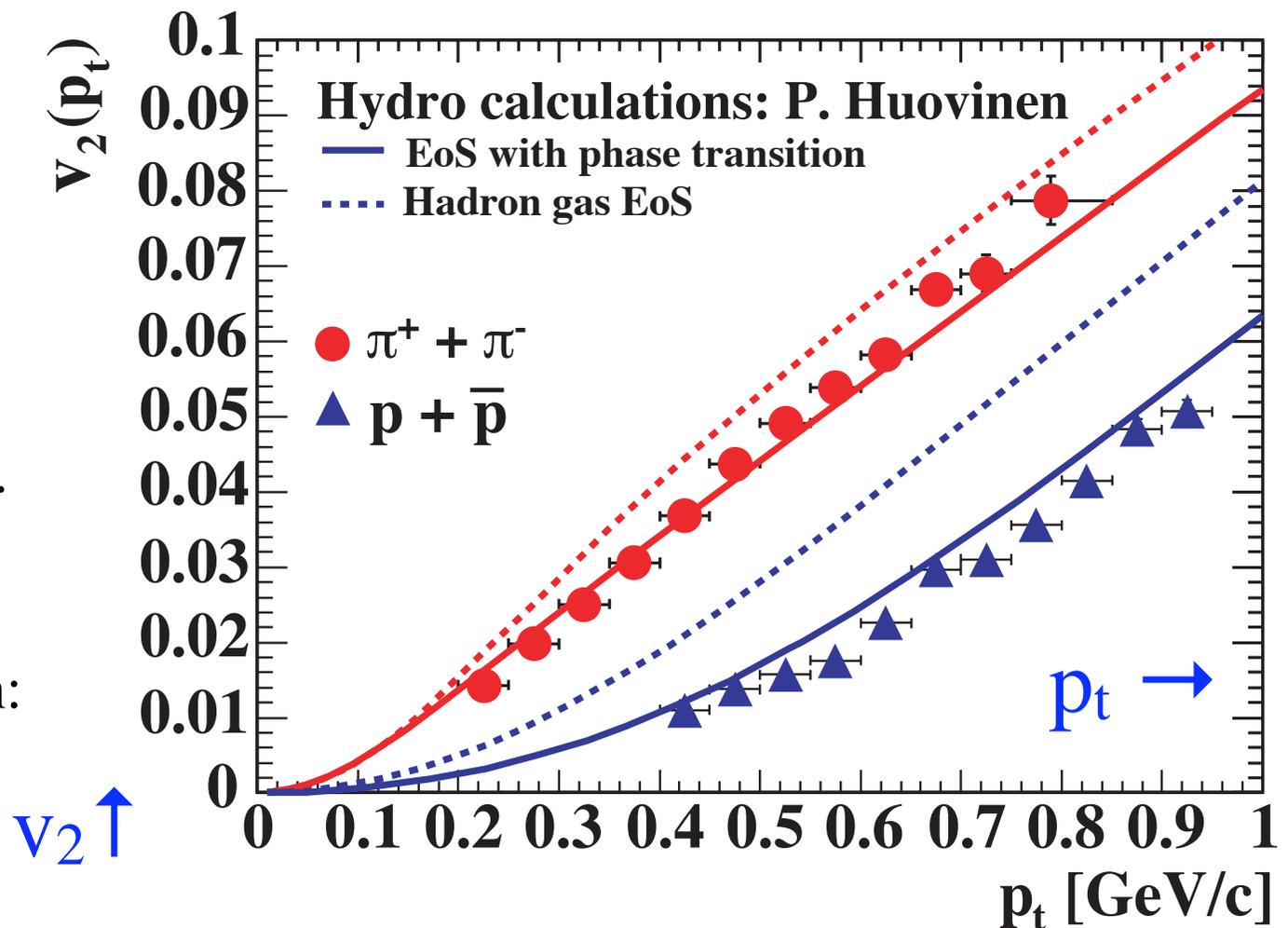
1. *Short* initial time (tune)
2. *MIT Bag* Equation of State (not correct)
3. *Zero* viscosity in QGP phase
4. Hadronic “afterburner”: \sim *large* hadronic viscosity

Good fit to π 's, K's, p's....
for *both* single particle
spectra *and* v_2 .

So far: *no* bound on
viscosity in QGP.

Viscous, relativistic hydro.
very hard.

Viscosity \sim 1/cross section:
small viscosity \Rightarrow
strong coupling
in sQGP?

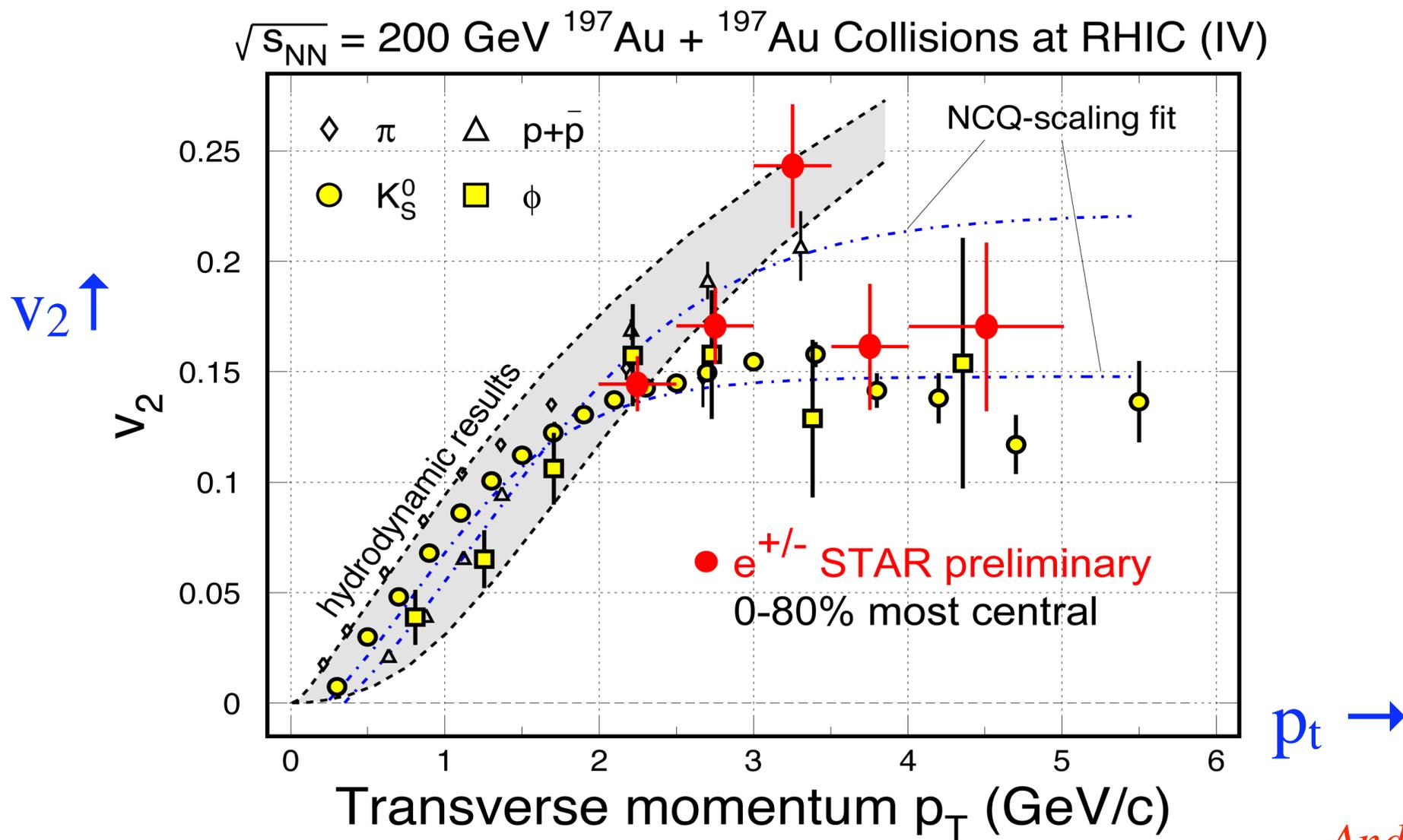


Even charm quarks flow

Look at charm quarks through single electrons.

See *large* elliptic flow: *no* suppression due to large mass.

Experimental definition of “sQGP”: heavy quarks affected *~same* as light quarks!



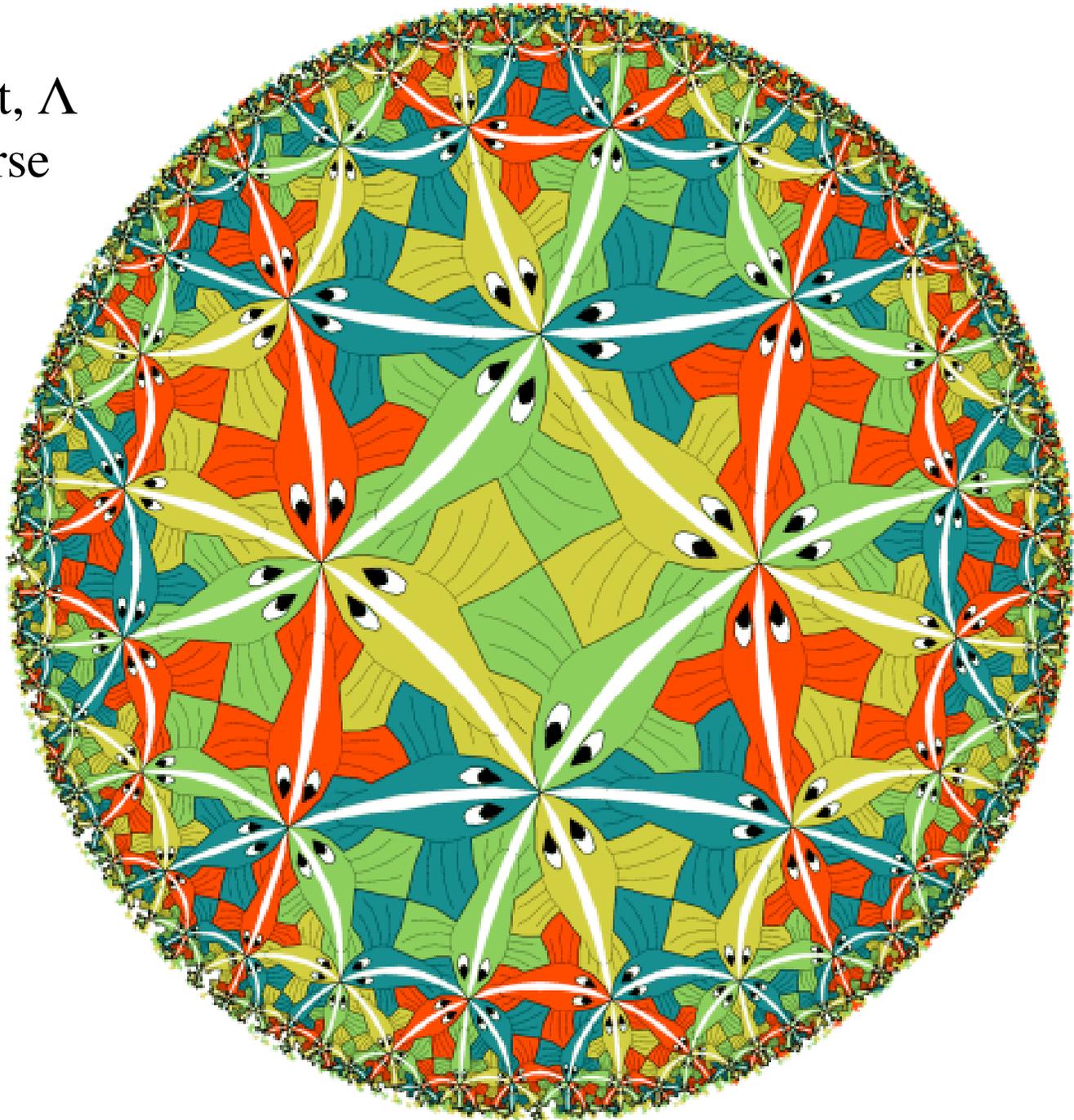
And now...

AdS: Anti de Sitter space

De Sitter space: Gravity with positive cosmological constant, Λ
Accelerated, expanding universe

Anti de Sitter: $\Lambda < 0$

Spatial cross section of AdS = hyperbolic space



M. K. Escher, courtesy of J. Maldacena

AdS/CFT Duality

Most supersymmetric QCD: “ $\mathcal{N}=4$ ” SUSY for SU(N) gluons: 4 supercharges.

Gluons (spin 1) + 4 spin 1/2 + 6 spin 0, all adjoint rep. No quarks.

One dimensionless coupling, α_s , but does *not* run! Extraordinary theory:

No mass scale, both classically and quantum mechanically!

Conformal Field Theory (CFT). Probably exactly soluble.

Maldacena’s Conjecture: $\mathcal{N}=4$ SU(∞) *dual* to string theory on AdS₅ x S⁵

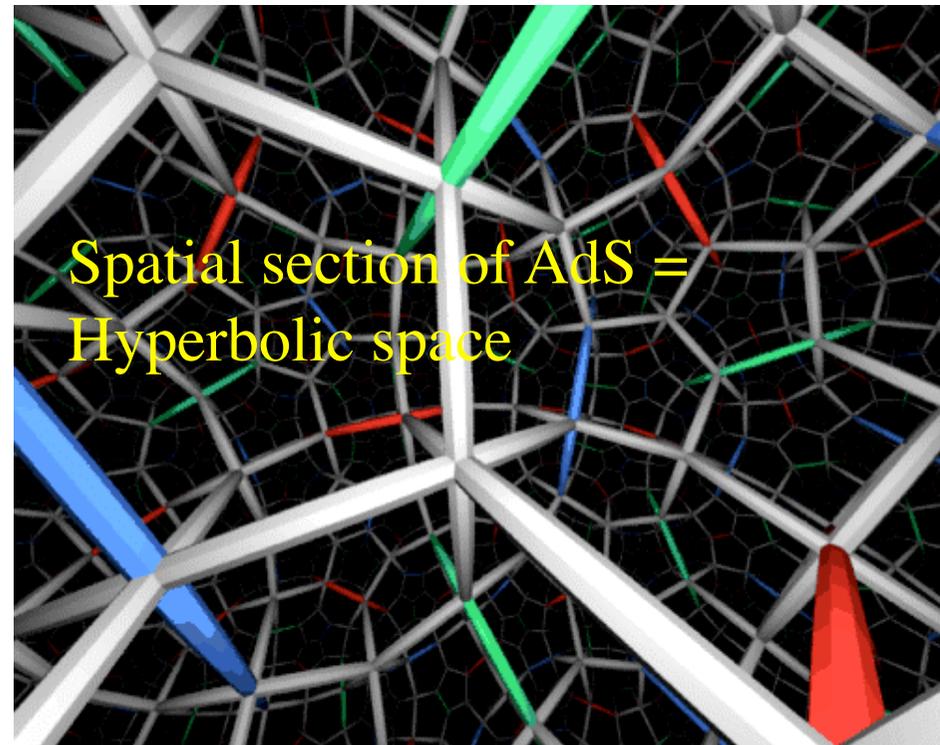
AdS₅ (AdS in 5 dimensions) + S⁵ (five sphere)

= Type IIB string in 10 dim.’s

AdS/CFT duality: Strong coupling in one theory is weak coupling in the other.

So *strong* coupling for $\mathcal{N}=4$ SU(∞) same as *weak* coupling on AdS₅ x S⁵.

Weak coupling string theory = *classical* supergravity!



Thermal AdS/CFT: $\alpha_s = \infty$

If one can compute with AdS/CFT,
often easier for $\alpha_s = \infty$ than $\alpha_s \approx 0$!

Results for $\mathcal{N} = 4$ SU(∞), infinite α_s :

Pressure 3/4 ideal! CFT \Rightarrow flat with T.

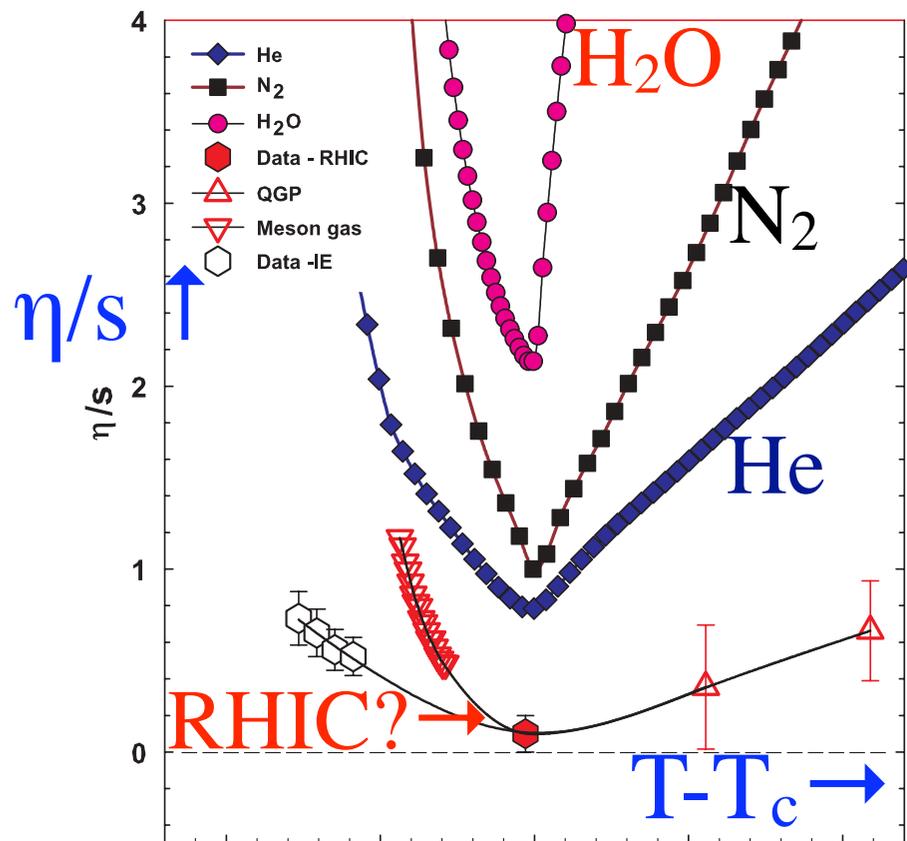
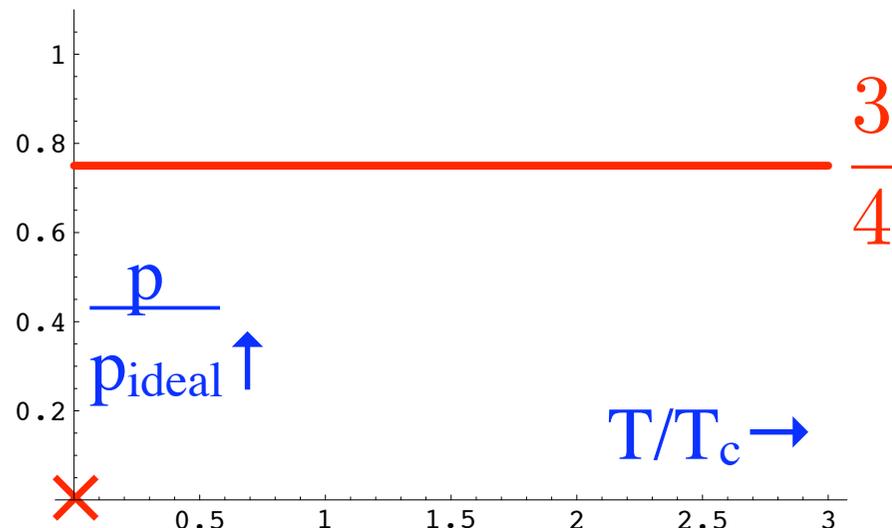
Not only equi., but also transport properties!

$\eta/s = \text{viscosity/entropy} = 1/(4\pi)$

Universal lower bound?

Perhaps. $\eta \sim 1/\alpha_s^2$ for $\alpha_s \approx 0$

In many systems, η/s has a minimum
near T_c . Maybe in the sQGP?



Shootout in the Unicorn Corral!

At RHIC, central $AA \neq A$ (pp) collisions

“Tail wags the Unicorn”

Clearest signal from “high” pt:

R_{AA} , jet suppression...

Body of the Unicorn: “sQGP”

Two packs of fighting dogs (aka theorists):

Strong coupling! AdS/CFT rules!

Weak coupling:

need lattice for equilibrium thermo.

need to extend to non-equi.: NaRPA!

Nothing better than a good dog fight...





"A possible eureka."

From the SPS: NA60 and a “thermal” ρ

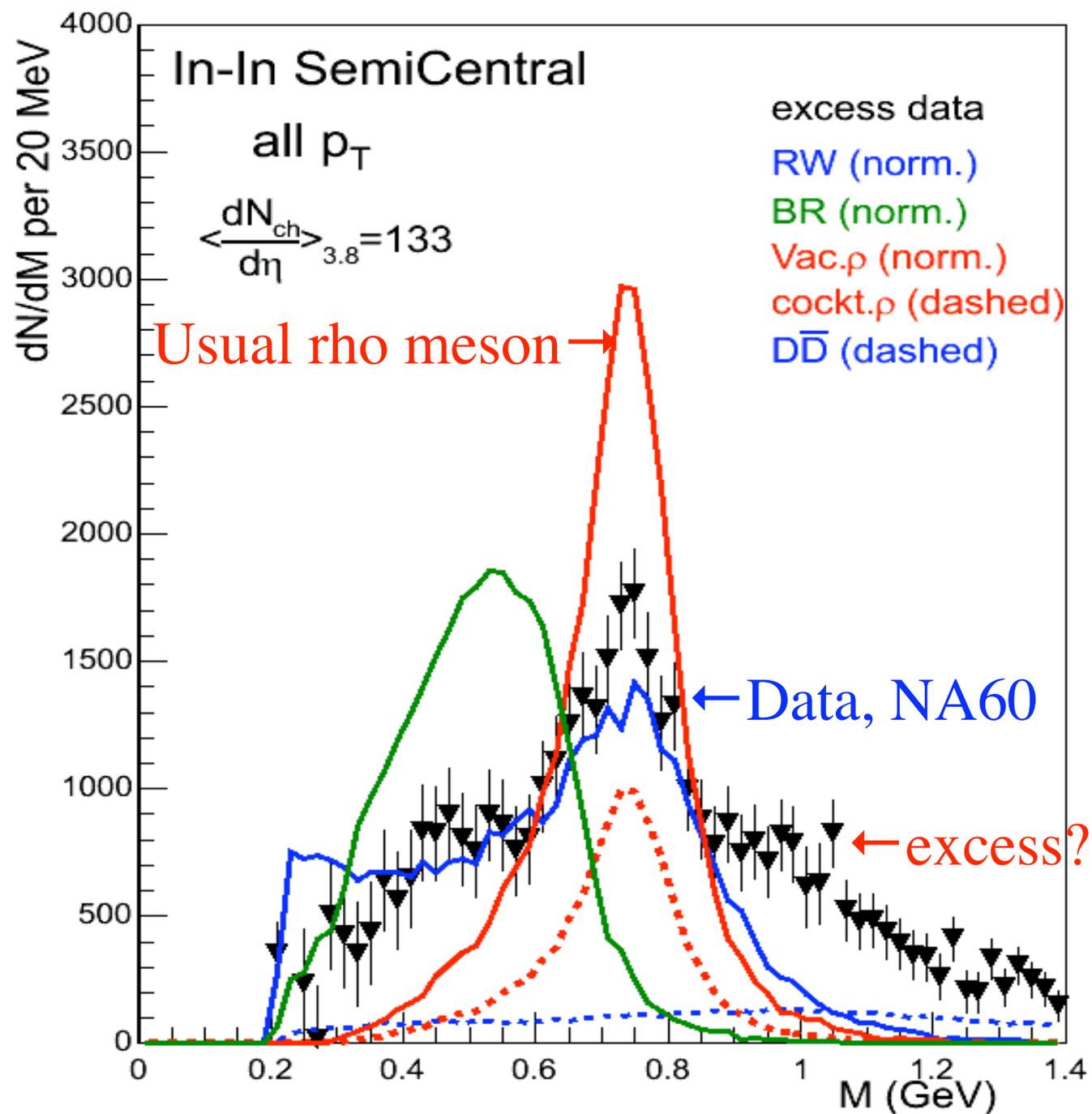
SPS = fixed target.
Kinematics more awkward, but can generate many more events.

Example: electromagnetic signals, such as e^+e^- pairs.

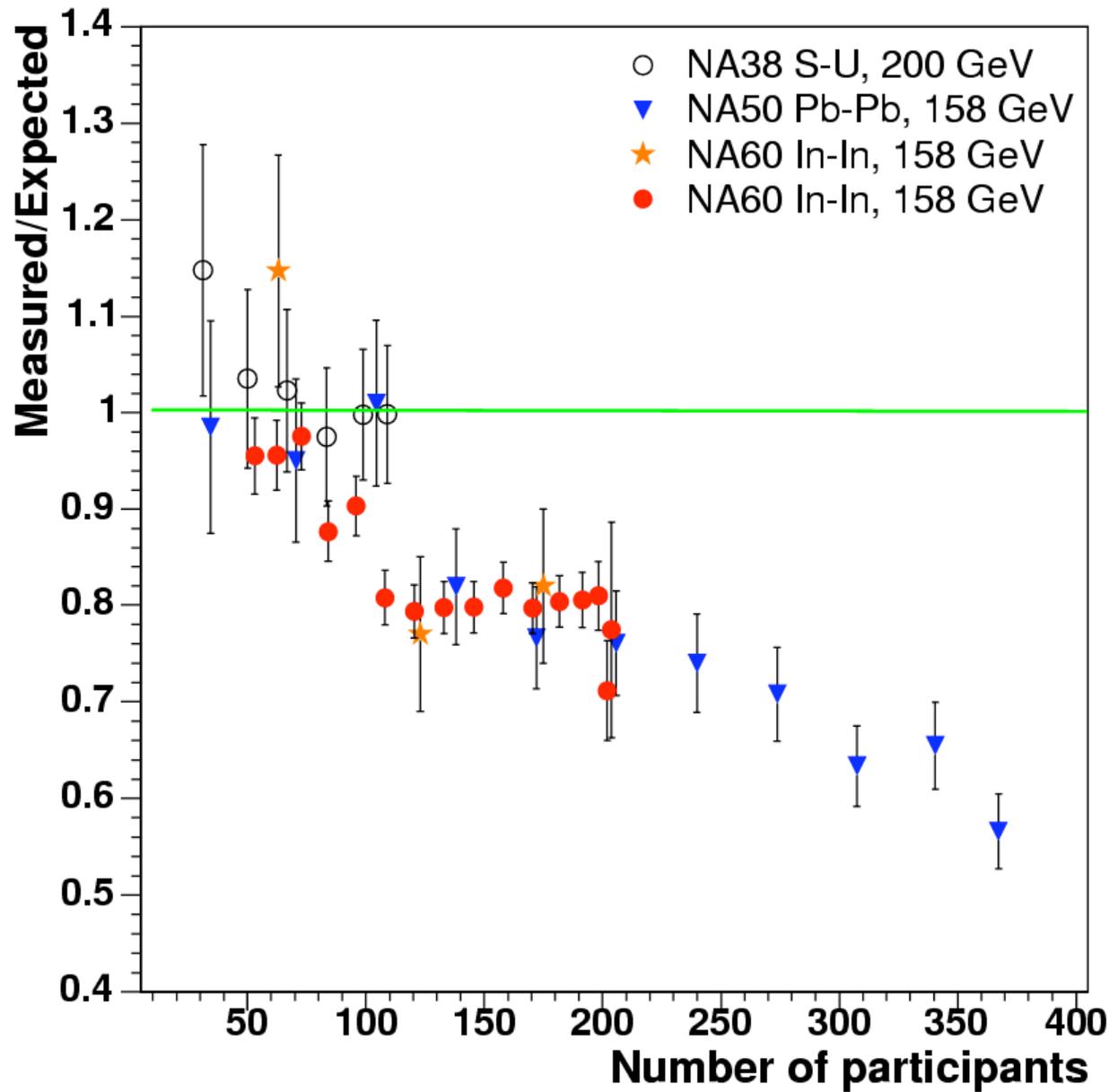
Look at “ ρ ” meson, mass ~ 770 MeV. Decays directly to e^+e^- .

Find “thermal” ρ : no shift in mass, thermal broadening.

Interesting excess above the ρ ?



SPS: NA50, NA60 J/Psi suppression



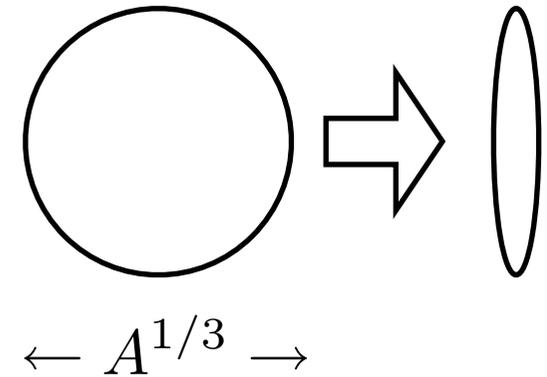
consistent, # participants good variable

Why do AA? “Saturation” as a Lorentz Boost

Incident nucleus Lorentz contracted at high energy

McLerran & Venugopalan: color charge bigger by $A^{1/3}$

$A \rightarrow \infty$: semi-classical methods, Color Glass



=> Logarithmic growth in multiplicity:

$$\frac{dN}{dy} \sim \frac{1}{g^2(\sqrt{s}/A)} \sim \log(\sqrt{s}/A)$$

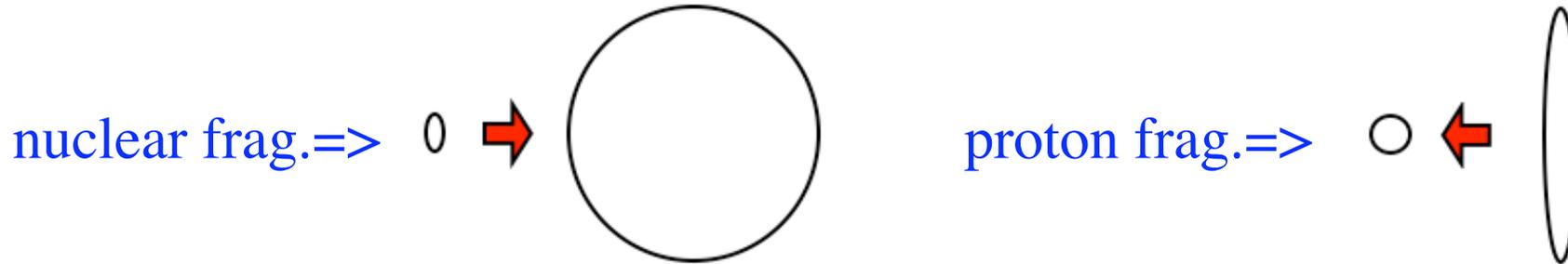
Expect at least same rise in $\langle p_t \rangle$.

Color Glass: “saturation momentum” function of energy, rapidity...

CG describes initial state. Final state?

Color Glass suppression: in dA, by the deuteron

Fragmentation region \sim rest frame. Incident projectile Lorentz contracted:



Nuclear fragmentation region: proton contracted. Study final state effects

Proton fragmentation region: study initial state effects

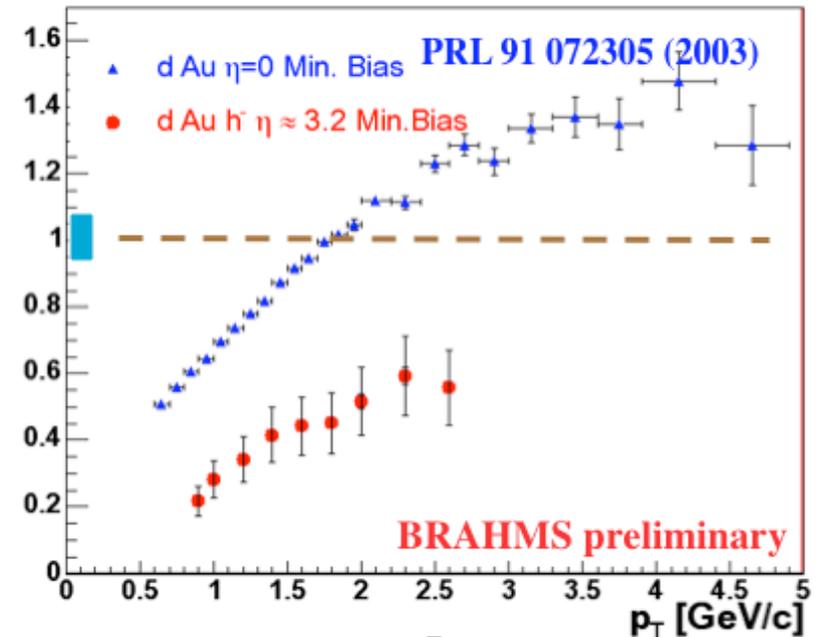
BRAHMS in dA:

enhancement @ zero rapidity
suppression @ proton frag. region.

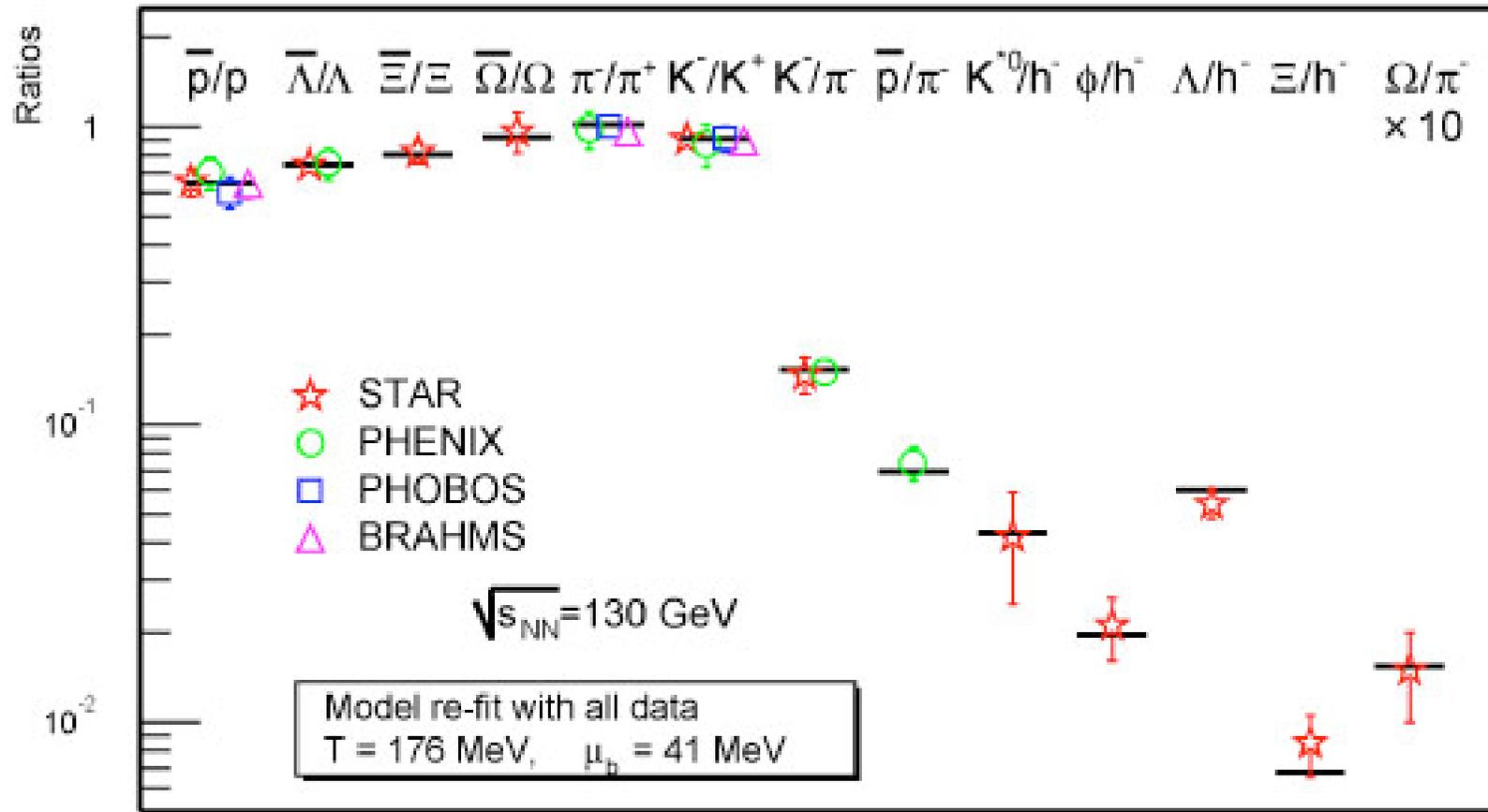
Supports color glass initial state.

Need to study all rapidities.

R_{dA} :



Total Chemical Ratios: $T_{\text{chemical}} \sim 160 \text{ MeV}$

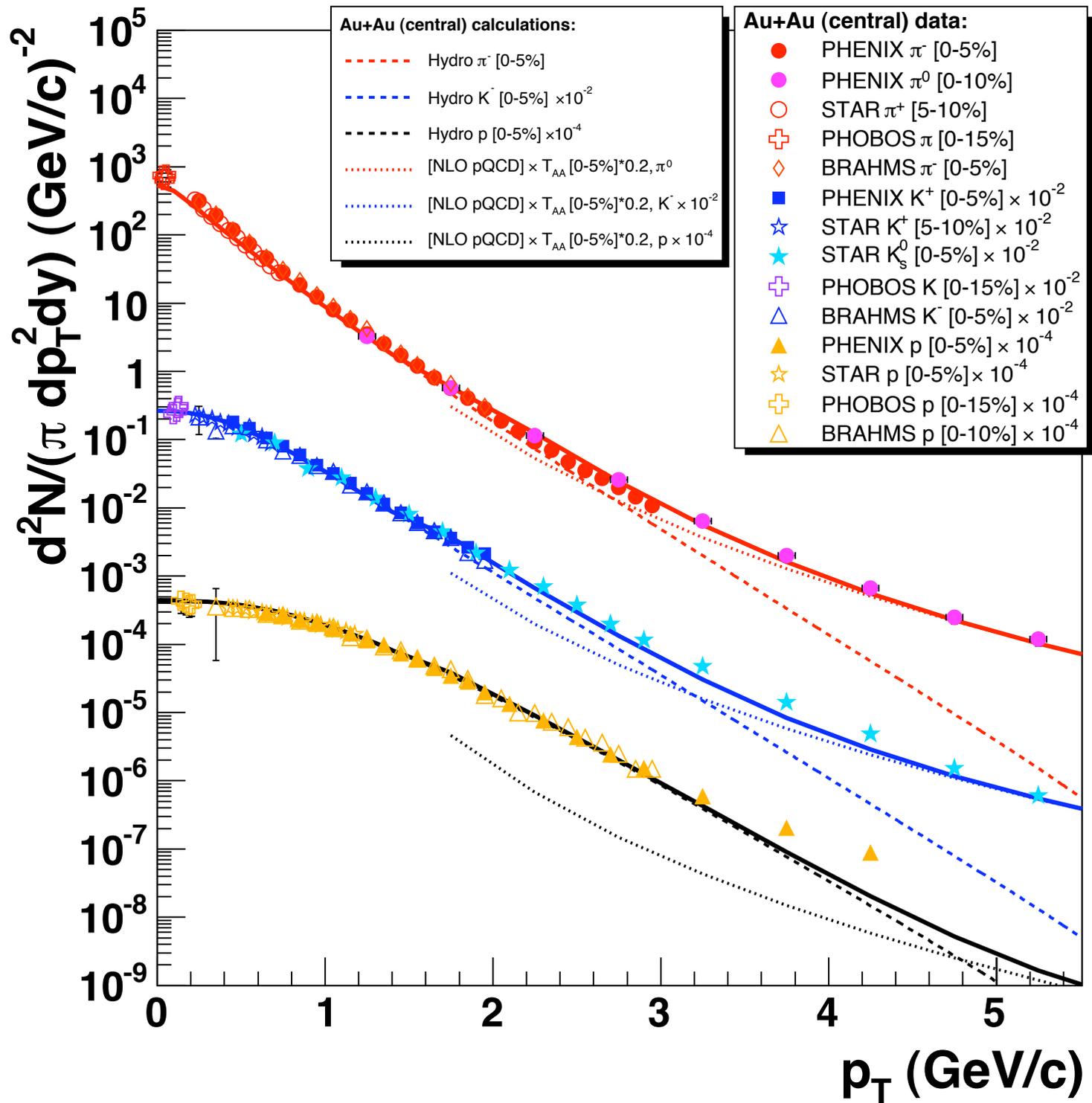


Braun-Munzinger et al., PLB 518 (2001) 41 D. Magestro (updated July 22, 2002)

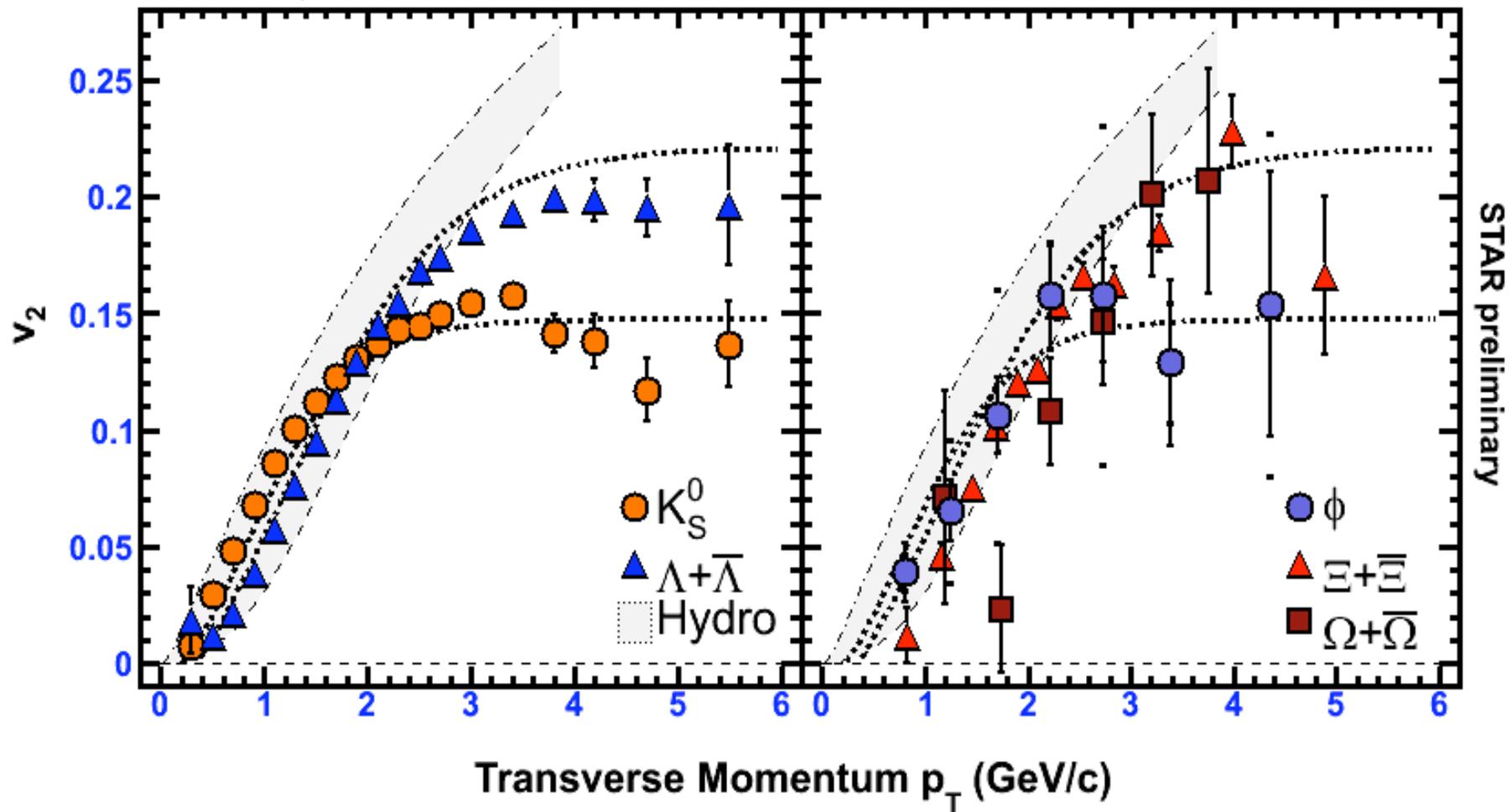
Overall chemical abundances for “long” lived states: well fit with

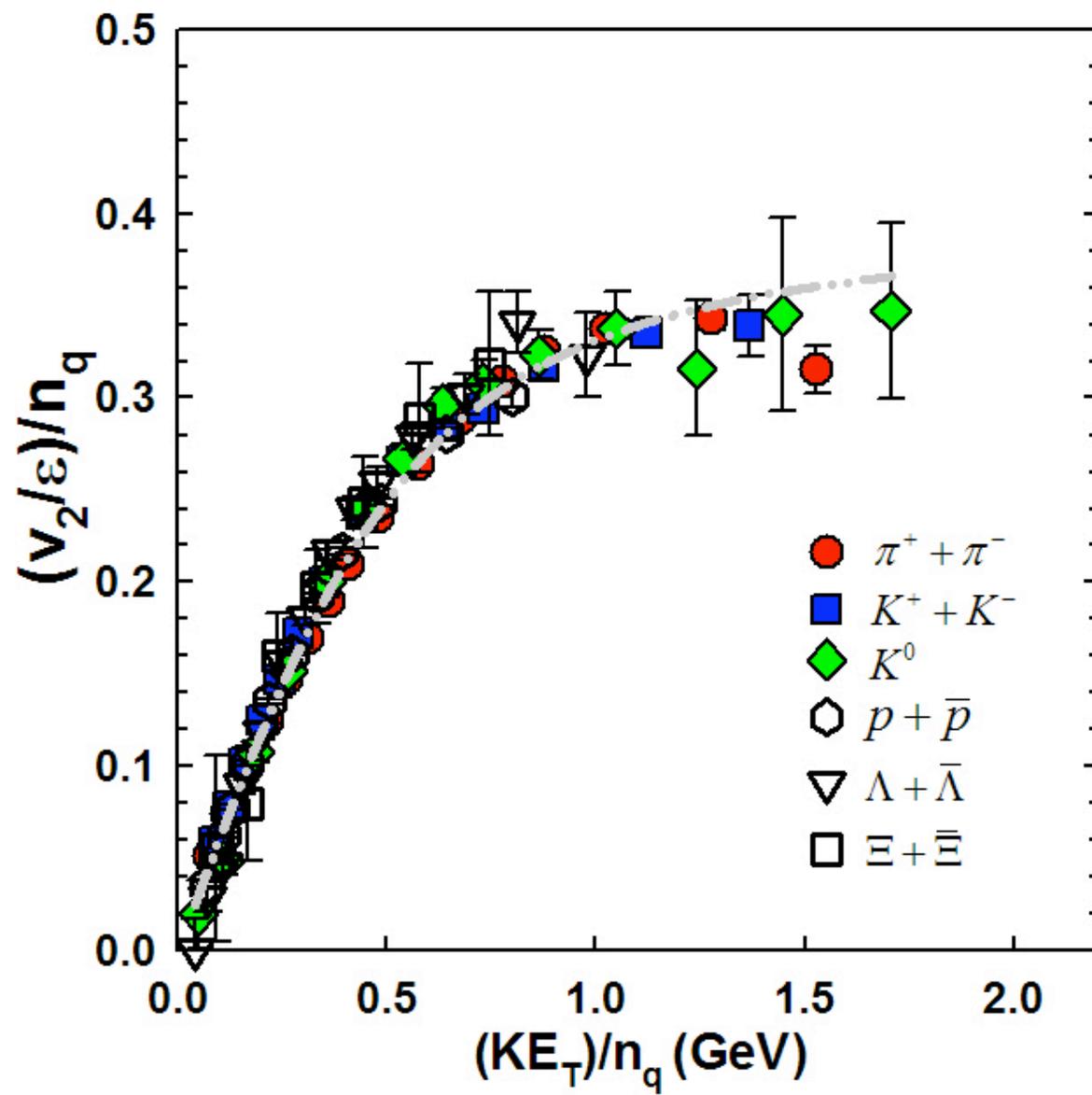
$$T_{\text{chemical}} = 160 \text{ MeV}, \mu_{\text{baryon}} = 24 \text{ MeV}$$

Not valid for “short” lived resonances: $\Delta, \phi, K^*, \Lambda^*$

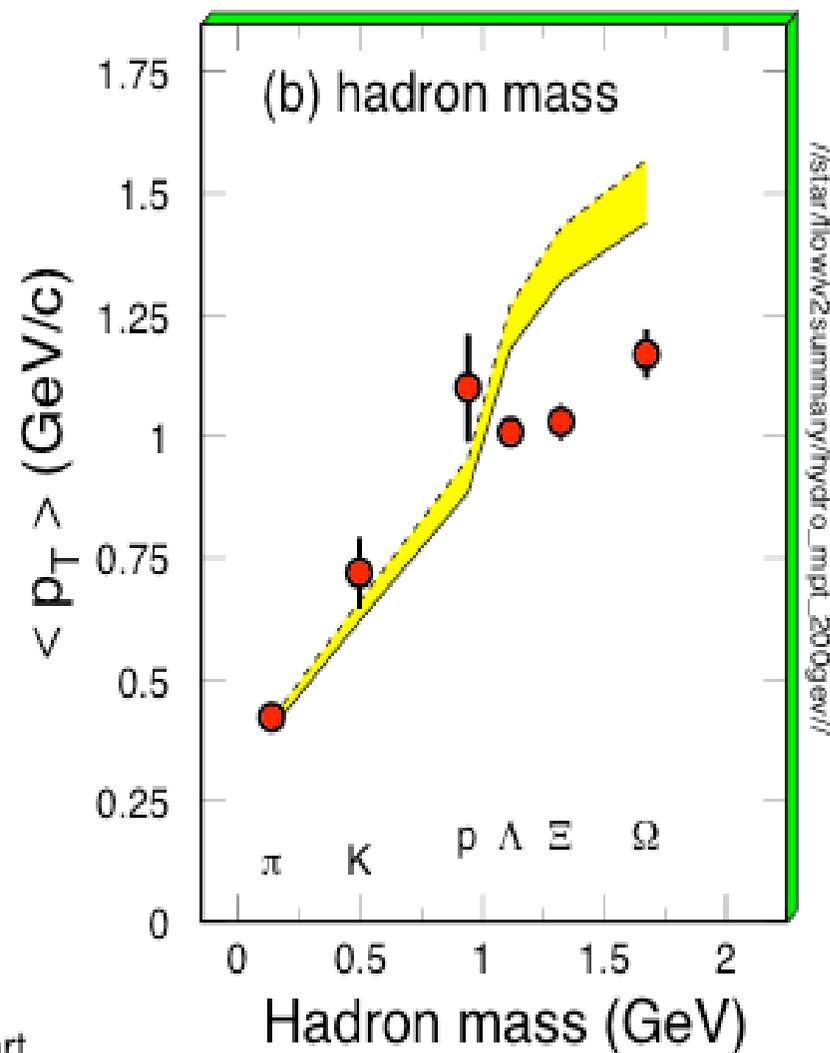
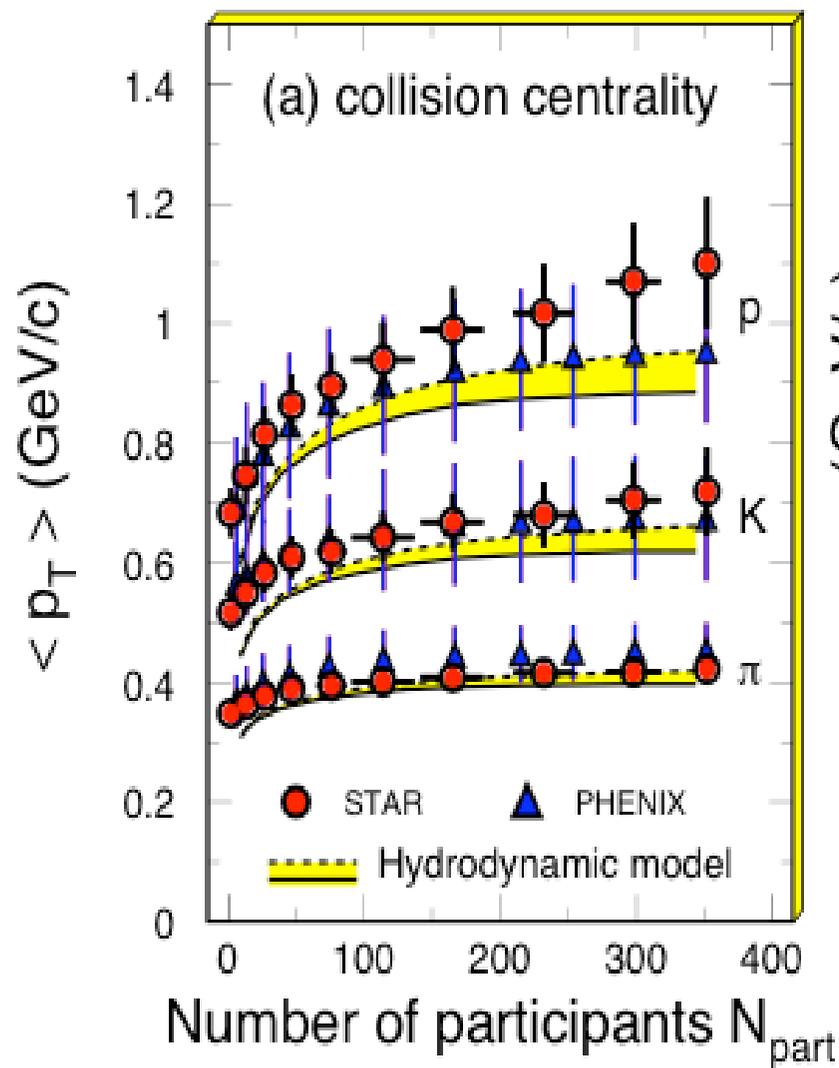


$\sqrt{s_{NN}} = 200 \text{ GeV } ^{197}\text{Au} + ^{197}\text{Au} \text{ Collisions at RHIC (run IV)}$





Au + Au Collisions at $\sqrt{s_{NN}} = 200$ GeV



/star/llw/v2/summary/hydro_mpl_200geV/

HBT radii: collisions “*explosive*”

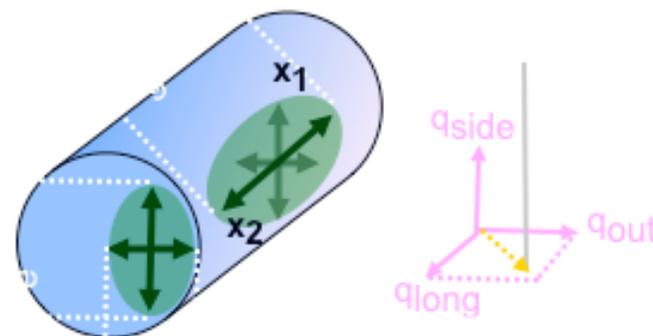
Hanbury-Brown-Twiss: two-particle correlations of identical particles

= *sizes at freezeout*. Three directions:

along beam R_{long} , along line of sight R_{out} , perpendicular R_{side} .

$$C(p_1, p_2) = N(p_1, p_2) / (N(p_1)N(p_2))$$

$$= 1 + \lambda \exp(-R_{HBT}^2 (p_1 - p_2)^2)$$



Hydro.: $R_{\text{out}}/R_{\text{side}} > 1$, *increases* with p_t
 (“burning log”)

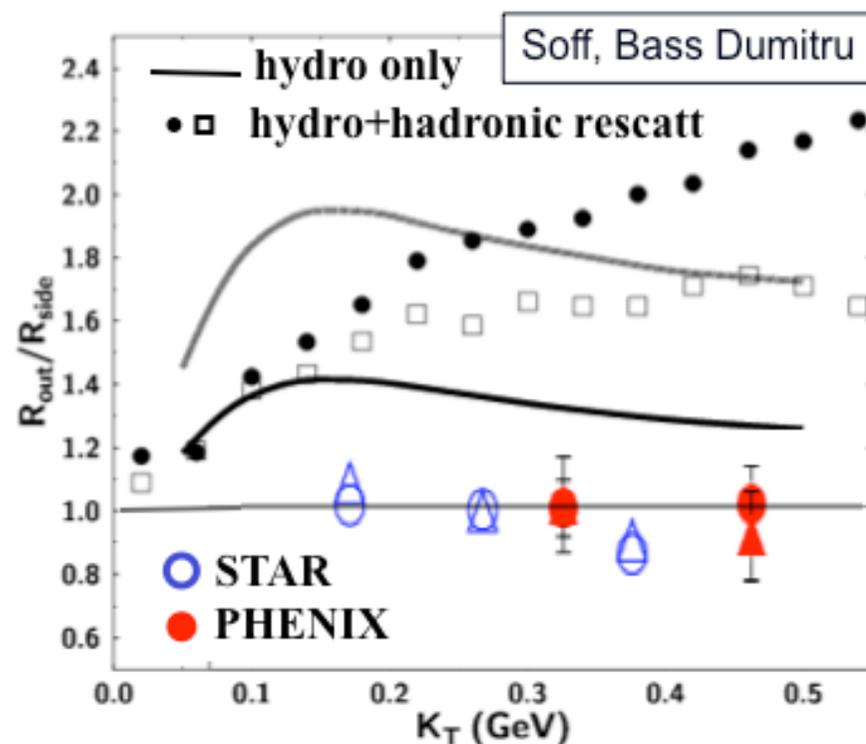
Exp.: $R_{\text{out}}/R_{\text{side}} \sim 1.0$, *flat* with p_t

Hydro. fails - *badly* - for HBT radii.

No big times from strong 1st order trans.!

HBT “*explosive*”: blast wave works:

Space-time history shell with
 lifetime $\sim 8\text{-}9$ fm/c, emission ~ 2 fm/c



HBT: p_t dependence same in pp, dA, AA!